|Vub| and |Vcb| determinations in the presence of New Physics

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Based on:

R. Fleischer, R. Jaarsma and GTX: 2104.04023 Eur.Phys.J.C 81 (2021) 7, 658

R. Fleischer, R. Jaarsma and GTX: 1809.09051 Eur.Phys.J.C 78 (2018) 11, 911

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Generalities

Effective Hamiltonian

 $\mathcal{H}_{\text{eff}}^{q} = \frac{4G_{\text{F}}}{\sqrt{2}} V_{qb} \left[\tilde{C}_{V_{L}}^{q,\ell} \mathcal{O}_{V_{L}}^{q,\ell} + \tilde{C}_{V_{R}}^{q,\ell} \mathcal{O}_{V_{R}}^{q,\ell} + \tilde{C}_{S}^{q,\ell} \mathcal{O}_{S}^{q,\ell} + \tilde{C}_{P}^{q,\ell} \mathcal{O}_{P}^{q,\ell} + \tilde{C}_{T}^{q,\ell} \mathcal{O}_{T}^{q,\ell} \right] + \text{h.c.},$

Operator Basis

$$\mathcal{O}_{V_L}^{q,\ell} = (\bar{q}\gamma^{\mu}P_Lb)(\bar{\ell}\gamma_{\mu}P_L\nu_{\ell}), \quad \mathcal{O}_{V_R}^{q,\ell} = (\bar{q}\gamma^{\mu}P_Rb)(\bar{\ell}\gamma_{\mu}P_L\nu_{\ell}), \quad \mathcal{O}_S^{q,\ell} = \frac{1}{2}(\bar{q}b)(\bar{\ell}P_L\nu_{\ell}),$$
$$\mathcal{O}_P^{q,\ell} = \frac{1}{2}(\bar{q}\gamma_5b)(\bar{\ell}P_L\nu_{\ell}), \quad \mathcal{O}_T^{q,\ell} = (\bar{q}\sigma^{\mu\nu}P_Lb)(\bar{\ell}\sigma_{\mu\nu}P_L\nu_{\ell}),$$

Introduction of New Physics

$$\tilde{C}_a^{q,\ell} = C_a^{(\mathrm{SM})q,\ell} + C_a^{q,\ell},$$

Generalities

Establish bounds on the NP Wilson Coefficients through ratios of leptonic and semileptonic processes

$$\mathcal{B}(B_q^- \to \ell^- \bar{\nu}_\ell) \qquad \mathcal{B}(\bar{B} \to P \ell^- \bar{\nu}_\ell) \qquad \mathcal{B}(\bar{B} \to \tilde{V} \ell^- \bar{\nu}_\ell)$$

Work under the assumption of NP in all the fermionic generations.

Correlate NP in light generations according to

Scenario 1: $C_X^e = 0.1 C_X^{\mu}$, Scenario 2: $C_X^e = C_X^{\mu}$,

Scenario 3: $C_X^e = 10C_X^{\mu}$.

where X=S, P, V, T



Strategy

1) Determine the NP regions using ratios of branching fractions where the CKM contributions cancel out.

2) Extract the CKM element by substituting the NP contributions in a CKM-NP dependent observable.

 $b \rightarrow u$ transitions: $|V_{ub}|$ $b \rightarrow c$ transitions: $|V_{cb}|$

Determine NP regions through |V_{cb}| |independent quantities

$$\mathcal{R}(D) = \frac{\mathcal{B}(\bar{B} \to D\tau^- \bar{\nu}_{\tau})}{\mathcal{B}(\bar{B} \to D\ell'^- \bar{\nu}_{\ell'})}, \qquad \mathcal{R}(D^*) = \frac{\mathcal{B}(\bar{B} \to D^* \tau^- \bar{\nu}_{\tau})}{\mathcal{B}(\bar{B} \to D^* \ell'^- \bar{\nu}_{\ell'})}$$

$$F_L(D^*) = \frac{\Gamma(B \to D_L^* \tau \bar{\nu}_\tau)}{\Gamma(B \to D^* \tau \bar{\nu}_\tau)} \qquad R_\mu^e(D^*) = \frac{\mathcal{B}(B^0 \to D^{*-} e^+ \nu_e)}{\mathcal{B}(B^0 \to D^{*-} \mu^+ \nu_\mu)}$$

 $b \rightarrow c$ transitions Experimental values 1909.12524 $\mathcal{R}(D) = 0.340 \pm 0.030$ $\mathcal{R}(D^*) = 0.295 \pm 0.01$ Belle: 1903.03102 Belle: 1809.03290 $R^e_{\mu}(D^*) = 1.01 \pm 0.01 \pm 0.03$ $F_L(D^*) = 0.60 \pm 0.08 \pm 0.04$ Theory values: 2104.04023 1703.05330 $\mathcal{R}(D^*)|_{\rm SM} = 0.253 \pm 0.005$ $\mathcal{R}(D)|_{\rm SM} = 0.300 \pm 0.006,$ $R^{e}_{\mu}(D^{*})|_{\rm SM} = 1.0045(1)$ $F_L(D^*)|_{\rm SM} = 0.458 \pm 0.004$



Results after applying all the constraints

Scenario	$C_{P}^{c,\mu}(1$	TeV)	$C_P^{c, au}($	1 TeV)	
	$\mathcal{R}(D^*)$		$\mathcal{R}(D^*)$	$\mathcal{R}(D^*), F_L(D^*)$	
	and	$R^e_\mu(D^*)$	and	and	
	$F_L(D^*)$		$F_L(D^*)$	$R^e_\mu(D^*)$	
	[-0.27, 0.27]		[-3.78, -2.32]	[-2.76, -2.32]	
$C_P^{c,e} = 10C_P^{c,\mu}$	([-0.05, 0.05]			
	[-0.27, 0.27]		[0.48, 1.94]	[0.48, 0.90]	
	[-2.0, 1.90]		[-3.78, -2.32]	[-3.52, -2.32]	1
$C_P^{c,e} = C_P^{c,\mu}$	([-1.46, 1.04]) (
	[-2.00, 1.89]		[0.48, 1.94]	[0.48, 1.66]	
	[-2.86, 2.65]		[-3.78, -2.32]	[-2.73, -2.32]	1
$C_P^{c,e} = 0.1C_P^{c,\mu}$		[-0.55, 0.34])
	[-2.84, 2.64]		[0.48, 1.94]	[0.48, 0.89]	

Determination of $|V_{cb}|$ To extract the CKM element we consider $\langle \mathcal{B}(\bar{B} \to D^* \tau^- \bar{\nu}_{\tau}) \rangle$ In particular we use the experimental result

$$\langle \mathcal{B}(\bar{B} \to D^* \tau^- \bar{\nu}_\tau) \rangle = (1.62 \pm 0.08) \times 10^{-2}$$

together with the constraint

 $\mathcal{B}(B_c \to \tau^- \bar{\nu}_\tau) < 0.60$





Exclusive value from HFLAV Collaboration:1909.12524 |Vcb|=0.03958+/-0.00117

Predictions on leptonic decays

$$\mathcal{B}(B_q^- \to \ell^- \bar{\nu}_\ell)|_{\mathrm{SM}} = \frac{G_{\mathrm{F}}^2}{8\pi} |V_{qb}|^2 M_{B_q} m_\ell^2 \left(1 - \frac{m_\ell^2}{M_{B_q}^2}\right)^2 f_{B_q}^2 \tau_{B_q}$$
$$\mathcal{B}(B_q^- \to \ell^- \bar{\nu}_\ell) = \mathcal{B}(B_q^- \to \ell^- \bar{\nu}_\ell)|_{\mathrm{SM}} \left|1 + C_{V_L}^{q,\ell} - C_{V_R}^q + \frac{M_{B_q}^2}{m_\ell} m_b + m_q\right) C_P^{q,\ell} \Big|^2$$



$b \rightarrow c$ transitions: vector(left handed)

Scenario	$C_{V_L}^{c,\mu}(1 \text{ TeV})$	$C_{V_L}^{c, au}(1 \text{ TeV})$
	$R^e_\mu(D^*)$	$\mathcal{R}(D^*), \mathcal{R}(D) ext{and} R^e_\mu(D^*)$
	[-0.183, -0.181]	$[-1.913, -1.857] \cup [-0.143, -0.087]$
$C_{V_L}^{c,e} = 10C_{V_L}^{c,\mu}$		
	[-0.001, 0.002]	$[-2.120, -2.045] \cup [0.045, 0.120]$
	[-1.828, -1.805]	$[-1.911, -1.855] \cup [-0.144, -0.089]$
$C_{V_L}^{c,e} = 0.1 C_{V_L}^{c,\mu}$		
	[-0.020, 0.015]	$[-2.116, -2.042] \cup [0.042, 0.116]$

CKM |*V*_{*cb*}| *results:* 0.038<|*V*_{*cb*}|<0.043

	Predictions		
Scenario	$\mathcal{B}(B_c \to e^- \bar{\nu}_e)$	$\mathcal{B}(B_c \to \mu^- \bar{\nu}_\mu)$	$\mathcal{B}(B_c \to \tau^- \bar{\nu}_{\tau})$
$C_{V_L}^{c,e} = 10C_{V_L}^{c,\mu}$	$[1.91 \times 10^{-9}, 2.36 \times 10^{-9}]$	$[8.39 \times 10^{-5}, 1.10 \times 10^{-4}]$	[0.026, 0.029]
$C_{V_L}^{c,e} = 0.1 C_{V_L}^{c,\mu}$	$[1.96 \times 10^{-9}, 2.36 \times 10^{-9}]$	$[8.08 \times 10^{-5}, 1.10 \times 10^{-4}]$	[0.026, 0.029]

Different possibilities to construct |Vub| independent quantities

$$\mathcal{R}^{\ell}_{\langle e,\mu\rangle;\rho\ [q^2\leq 12]\ \mathrm{GeV}^2} \equiv \mathcal{B}(B^- \to \ell^- \bar{\nu}_{\ell}) / \left\langle \mathcal{B}(\bar{B} \to \rho \ell'^- \bar{\nu}_{\ell'}) \right\rangle_{[\ell'=\ e,\mu],\ q^2\leq 12\ \mathrm{GeV}^2},$$

$$\mathcal{R}^{\ell}_{\langle e,\mu\rangle;\pi} \equiv \mathcal{B}(B^{-} \to \ell^{-} \bar{\nu}_{\ell}) / \left\langle \mathcal{B}(\bar{B} \to \pi \ell'^{-} \bar{\nu}_{\ell'}) \right\rangle_{[\ell'=e,\mu]},$$

 $\mathcal{R}^{\langle e,\mu\rangle;\rho}_{\langle e,\mu\rangle;\pi} \stackrel{[q^2 \le 12] \text{ GeV}^2}{=} \left\langle \mathcal{B}(\bar{B} \to \rho \ell'^- \bar{\nu}_{\ell'}) \right\rangle_{[\ell'=\ e,\mu],\ q^2 \le 12 \text{ GeV}^2} / \left\langle \mathcal{B}(\bar{B} \to \pi \ell'^- \bar{\nu}_{\ell'}) \right\rangle_{[\ell'=e,\mu]}$

Experimental results leptonic

$$\mathcal{B}(B^- \to \mu^- \bar{\nu}_{\mu}) = (5.3 \pm 2.2) \times 10^{-7}$$
$$\mathcal{B}(B^- \to \tau^- \bar{\nu}_{\tau}) = (1.09 \pm 0.24) \times 10^{-4}$$
$$\mathcal{B}(B^- \to e^- \bar{\nu}_e) < 9.8 \times 10^{-7} (90\% \text{ C.L.}).$$

Belle: 1911.03186 PDG Belle:0611045

Experimental results semileptonic

$$\left< \mathcal{B}(\bar{B} \to \pi \ell'^- \bar{\nu}_{\ell'}) \right>_{[\ell'=e,\mu]} = (1.53 \pm 0.04) \times 10^{-4} \quad PDG$$

 $\mathcal{B}(\bar{B}^0 \to \pi^+ \tau^- \bar{\nu}_{\tau}) < 2.5 \times 10^{-4} \,(90\% \text{ C.L.})$ Belle: 1509.06521

$$\left\langle \mathcal{B}(\bar{B} \to \rho \ell'^- \bar{\nu}_{\ell'}) \right\rangle_{[\ell'=e,\mu], q^2 \le 12 \text{ GeV}^2} = (1.98 \pm 0.12) \times 10^{-4}.$$

Belle: 1306.2781

Example Pseudoscalar Operator



Example Pseudoscalar Operator

Scenario: $C_P^{u,e} = C_P^{u,\mu}$			
Region	$C_P^{u,\mu}(1 \text{ TeV})$	$C_P^{u, au}(1 \text{ TeV})$	$ V_{ub} $
1	[-0.0001, 0.0055]	[-0.05, 0.13]	$\left[0.0030, 0.0037 ight]$
2	[-0.0001, 0.0055]	[-0.44, -0.27]	$\left[0.0030, 0.0037 ight]$
3*	[-0.024, -0.018]	[-0.44, -0.27]	$\left[0.0030, 0.0037 ight]$
4*	[-0.024, -0.018]	[-0.05, 0.13]	$\left[0.0030, 0.0037 ight]$

Exclusive value from HFLAV Collaboration:1909.12524

|V_{ub}|=0.00367+/-0.00015

Predictions

Scenario: $C_P^{u,e} = C_P^{u,\mu}$		
Region	$\mathcal{B}(B^- \to e^- \bar{\nu}_e)$	$\mathcal{B}(\bar{B} \to \rho \tau^- \bar{\nu}_\tau)_{q^2 \le 12 \text{ GeV}^2}$
1	$[0, 1.4 \times 10^{-7}]$	$[5.2 \times 10^{-5}, 1.2 \times 10^{-4}]$
2	$[0, 1.4 \times 10^{-7}]$	$[4.3 \times 10^{-5}, 9.3 \times 10^{-5}]$



Final Remarks

- We have presented a simple strategy to extract the CKM elements $|V_{ub}|$ and $|V_{cb}|$ in the presence of NP contributions.
- To illustrate our approach we have assumed the presence of NP in all the generations of fermions.
- Our studies took into account different Lorenz structures finding values for $|V_{cb}|$ consistent with the exclusive determinations.
- The presence of pseudoscalar mediators leads to potential enhancements in the branching fractions of not yet measured leptonic decays.