



Holographic approach of the spinodal instability to criticality

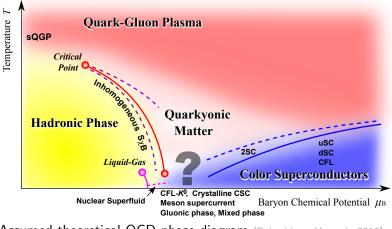
Maximilian Attems (CERN)

arXiv:1905.12544 JHEP (2020), arXiv:2012.15687 JHEP (2021)

A Virtual Tribute to Quark Confinement and the Hadron Spectrum 2021

Motivation - QCD phase diagram

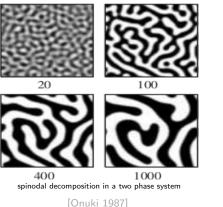
The critical point (CP) is the endpoint of the 1st order phase transition between the Quark-Gluon Plasma and the hadrons:



Assumed theoretical QCD phase diagram [Fukushima, Hatsudo 2010]

Holographic setup

Spinodal instability is key signature of 1st-order phase transition: spinodal instability: Gregory-Laflamme instability:

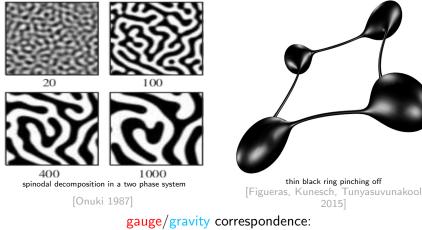


thin black ring pinching off [Figueras, Kunesch, Tunyasuvunakool 2015]

gauge/gravity correspondence:

bridge between physical phenomena in gauge theories and gravity.

Spinodal instability is key signature of 1st-order phase transition: Gregory-Laflamme instability: spinodal instability:



bridge between physical phenomena in gauge theories and gravity.

thin black ring pinching off

Holographic bottom-up model: setup

Dual field theory: 'mimics'a deformation of N=4 SYM with a dimension 3 operator ${\it O}$ and source Λ as 'mass'

$$S_{
m Gauge Theory} = S_{
m conformal} + \int d^4 x \Lambda O$$

Einstein-Hilbert action coupled to a scalar with non-trivial potential in five-dimensional bottom-up model:

$$S = \frac{2}{\kappa_5^2} \int d^5 x \sqrt{-g} \left[\frac{1}{4} \mathcal{R} - \frac{1}{2} \left(\nabla \phi \right)^2 - V(\phi) \right]$$

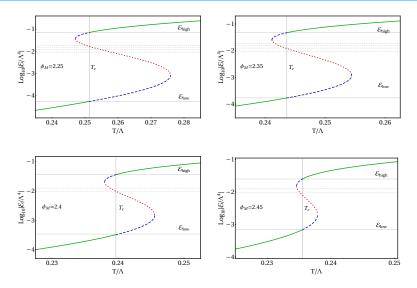
We derive the potential from the superpotential [Bianchi, Freedman, Skenderis 2002]: $V(\phi) = -\frac{4}{3}W(\phi)^2 + \frac{1}{2}W'(\phi)^2$

$$\ell^2 W(\phi) = -rac{3}{2} - rac{\phi^2}{2} - rac{\phi^4}{4\phi_M^2}$$

using a single parameter with critical value $\phi_M = 2.521$

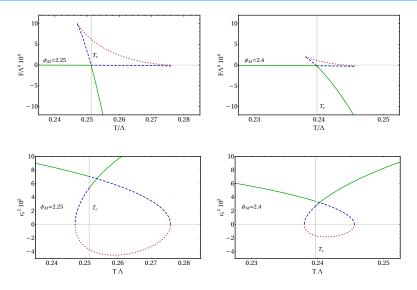
$$\ell^2 V(\phi) = -3 - rac{3\phi^2}{2} - rac{\phi^4}{3} - rac{\phi^6}{3\phi_M^2} + rac{\phi^6}{2\phi_M^4} - rac{\phi^8}{12\phi_M^4}$$

Holographic bottom-up model: tuning criticality



Equation of states with stronger or softer first-order phase transition with $\phi M = \{2.25, 2.3, 2.35, 2.4, 2.45\}$.

Holographic bottom-up model: tuning criticality II



free energy F and the speed of sound squared cs^2 over temperature for the theories with different criticality

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four generic stages:

- 1 linear stage
- 2 reshaping
- 3 merger
- 4 final:

static + phase-separated

endstate:

phase-separated configuration

conjecture:

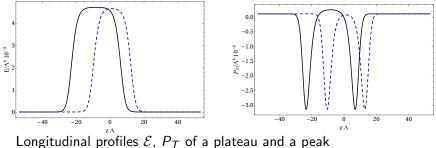
all static, non-phase separated configurations are dynamically unstable

Spinodal instability - criterium

inhomogeneous structures: 1 phase separated, 2 phase mixed maxima: minima:

I plateauxI valley2 peak2 gorge

New transverse pressure $P_T \ge 0.8P_c$ criterium for distinction:



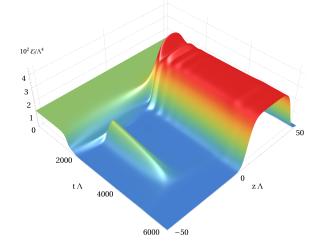
Formation time of the spinodal instability with varying criticality:

$oldsymbol{\phi}_{\mathrm{M}}$	2.25	2.35	2.4	2.45
$t_{ m formation}\Lambda$	880	1380	1725	2660

The time until \mathcal{E}_{high} is reached heavily depends on criticality.

Spinodal instability dissipation

Dissipation of a peak into a plateau:



First simulations where the rigid interface disappear for different criticality.

Main lessons from holographic studies of criticality

- Fast spinodal instability away from CP
- distinction of the inhomogeneous structures of a spinodal instability via new criterium: peaks versus domains; valleys versus gorges
- classification of the dynamics of a spinodal instability: quasi-linear regime, reshaping, merger, preferred final state
- new dissipation with spinodal instabilites
- More studies on the way..