

Improving the cold-QM pressure via soft interactions at N3LO

Tyler Gorda
TU Darmstadt

vConf21, 5 August 2021

Work in Collaboration with:

A. Kurkela, R. Paatelainen, S. Säppi,
and A. Vuorinen

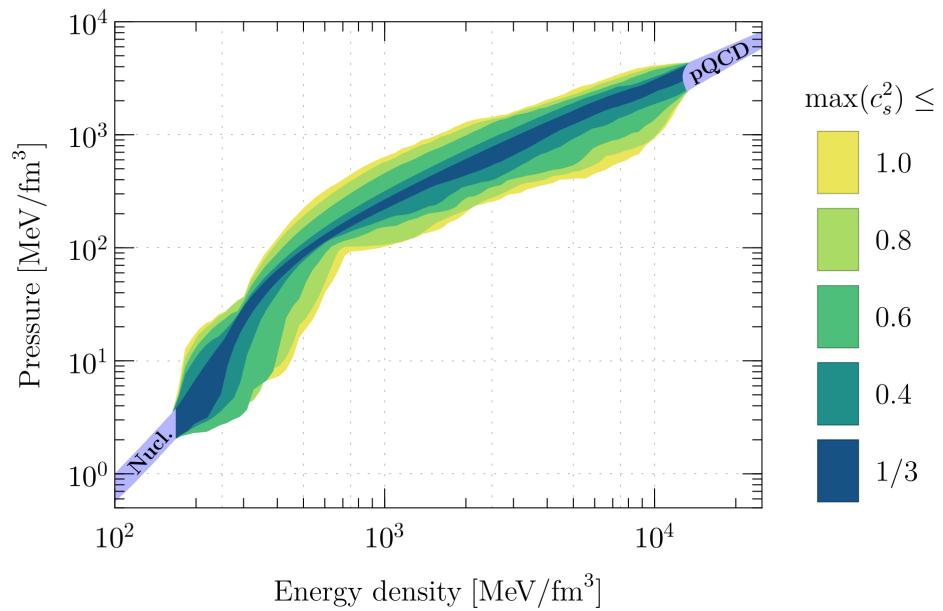
arXiv 2103.05658, 2103.07427



TECHNISCHE
UNIVERSITÄT
DARMSTADT

Motivation

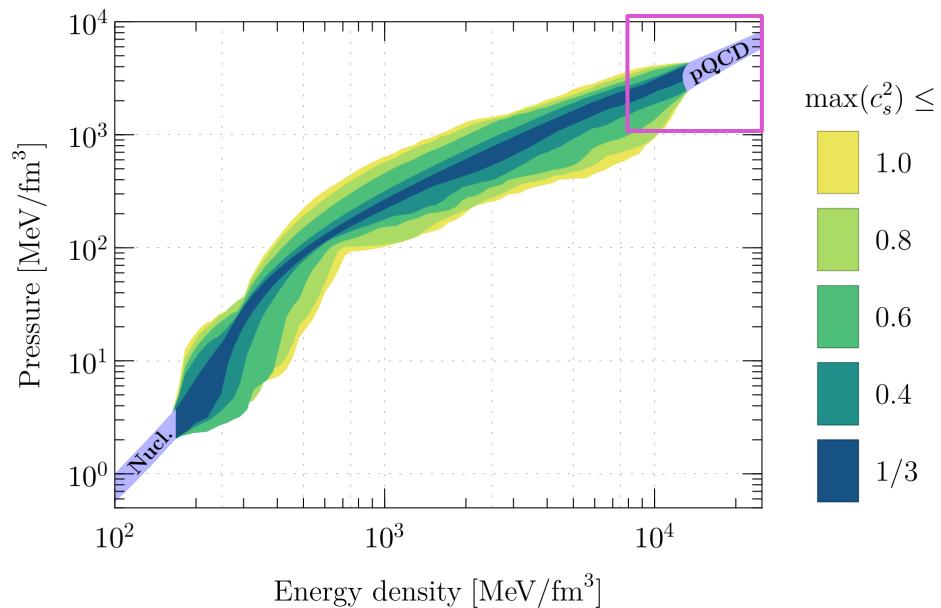
- Cold QM pQCD EOS relevant for constraining NS EOS:



Annala, TG, Kurkela, Näyttälä,
Vuorinen, Nature Physics (2020)

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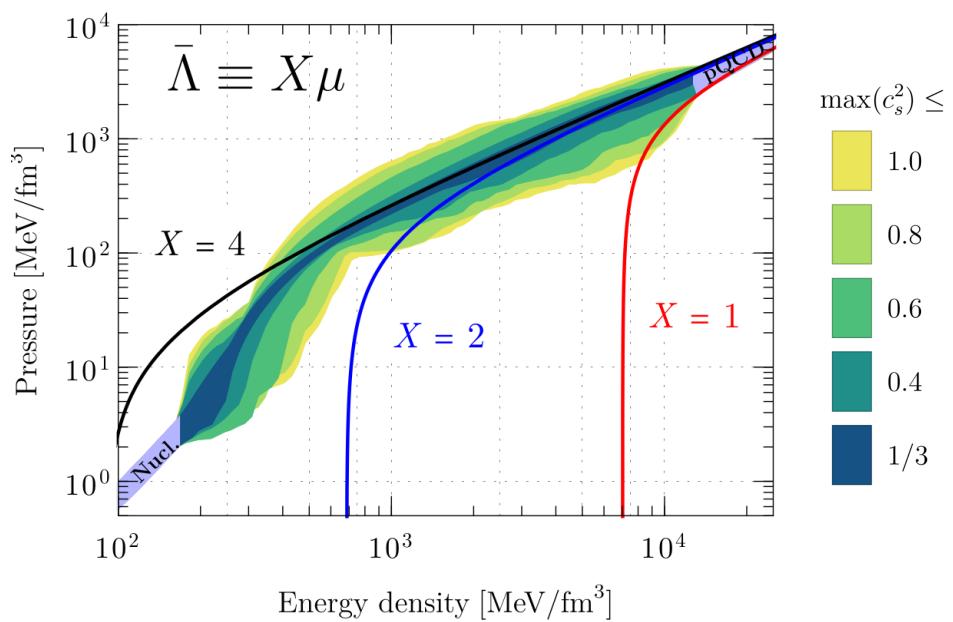


Even though pQCD
EOS only valid for
 $n_B \gtrsim 40n_0$,
CONSTRAINS NS
EOS at lower
densities

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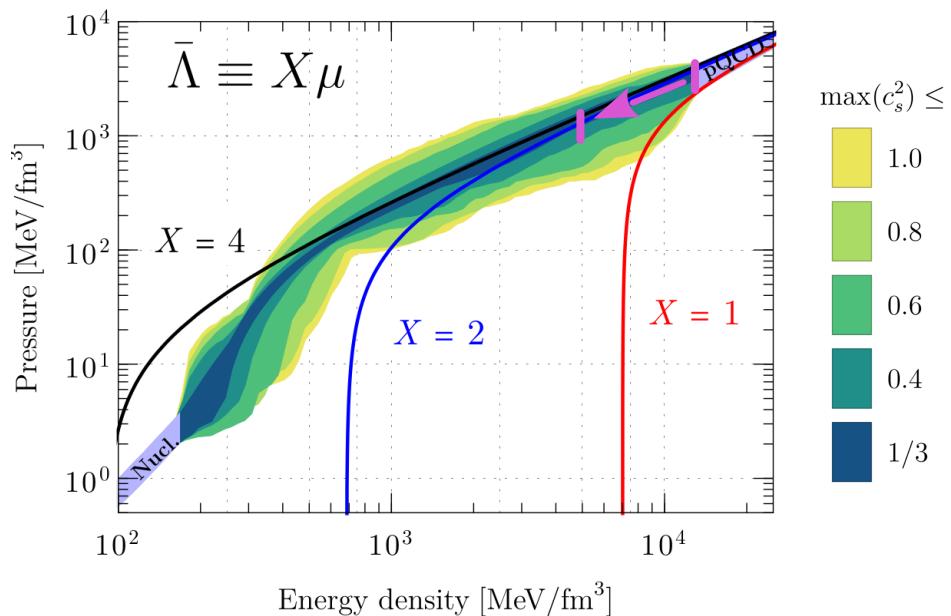
- Cold QM pQCD EOS has poor (renorm.) scale sensitivity, which prevents us from using it to lower densities:



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- Cold QM pQCD EOS has poor (renorm.) scale sensitivity, which prevents us from using it to lower densities:



If we could improve the dependence on the renorm. scale, could have dramatic consequences

Annala, TG, Kurkela, Näyttälä,
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Calculating the pressure; a tale of EFTs

Calculating the pressure; a tale of EFTs

Hot QGP

Three scales:

- 1) $P \sim T$: Full-theory (hard) diagrams
- 2) $P \sim \alpha_s^{1/2} T$: EFT for (massive) chromo-electric fields
- 3) $P \sim \alpha_s T$: Lattice EFT for (massless) chromo-magnetic fields

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Cold QM

Two scales:

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No softer scale b/c gluons not thermally occupied at $T = 0$: Great!

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Not great

Calculating the pressure; a tale of EFTs

Effective in-medium mass scale:

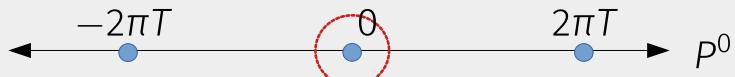
$$m_E \sim \alpha_s^{1/2} \mu$$

$$m_E^2 = \lim_{|P| \rightarrow 0} \Pi^{\mu\nu}(P)$$

(coeff in front; also angular function)

Hot QGP

sum-integrals: $T \sum_{n=-\infty}^{\infty} \int \frac{d^3 P}{(2\pi)^3}$



- Only zero-mode requires special treatment for $m_E \ll T$
- **3d EFT of massive zero mode “dimensional reduction”**

Cold QM

4d-Euclid.
Integrals:

$$\int \frac{d^4 P}{(2\pi)^4}$$



- **No simple separation**
- **No simple EFT to deal with IR problems**

Pressure contributions in cold QM

$$p = p_{\text{FD}} + \boxed{p_1^h \alpha_s + p_2^h \alpha_s^2 + p_3^h \alpha_s^3} \quad P \gtrsim \mu$$
$$\quad \quad \quad + \boxed{p_2^s \alpha_s^2 + p_3^s \alpha_s^3} \quad P \lesssim m_E$$
$$\quad \quad \quad + \boxed{p_3^m \alpha_s^3}, \quad \text{Both scales}$$

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Ambiguity in soft/hard split ($m_E \ll K \ll \mu$) gives logarithmic sensitivity to a **factorization mass scale Λ_h , which cancels out of sum over all kinematic regions (**columns!**)

Pressure contributions in cold QM

Freedman,
McLerran
(1977)

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This WORK

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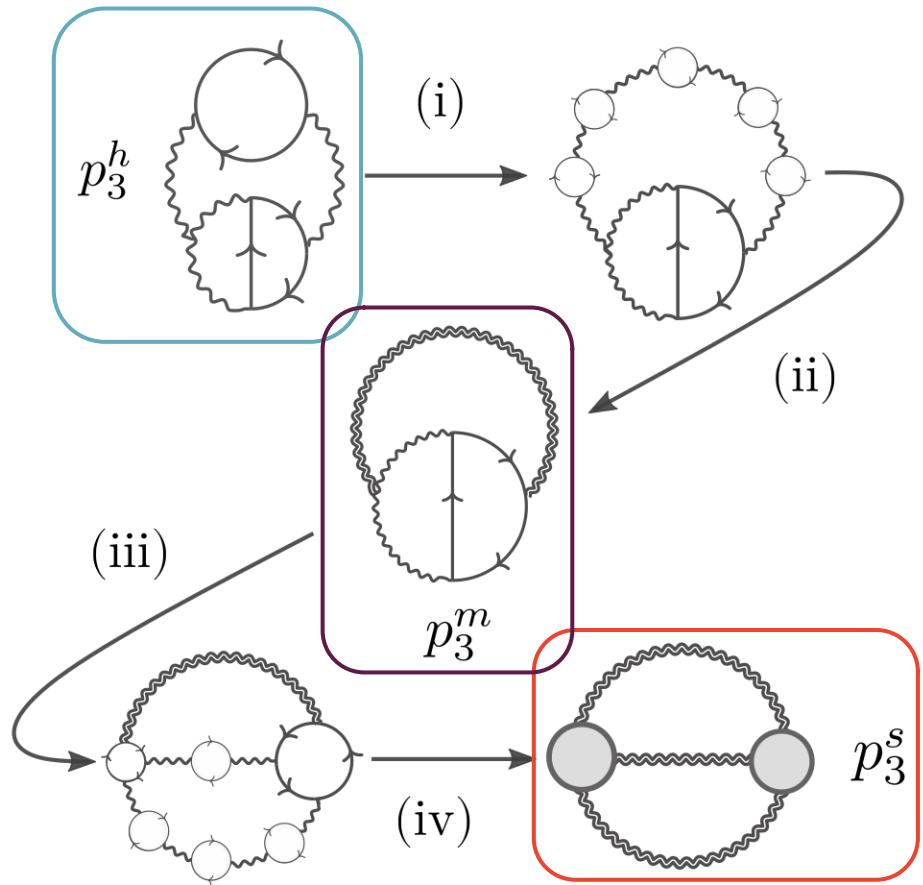
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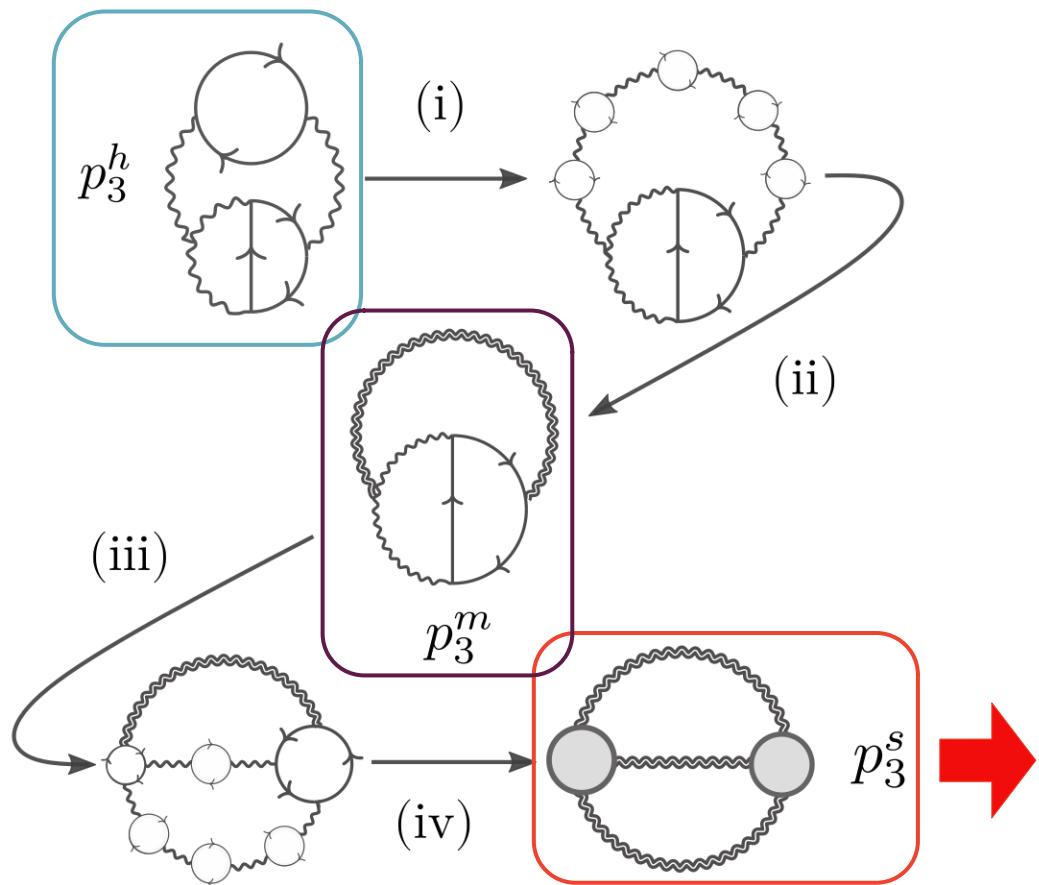
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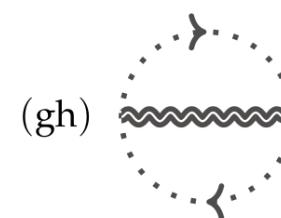
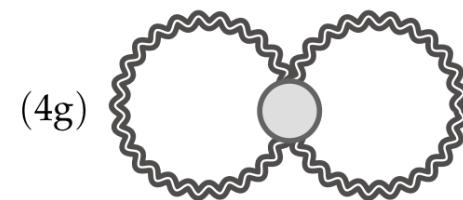
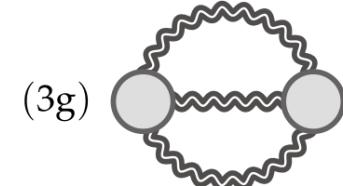
Cold QM approach



Cold QM approach



p_3^s contributions

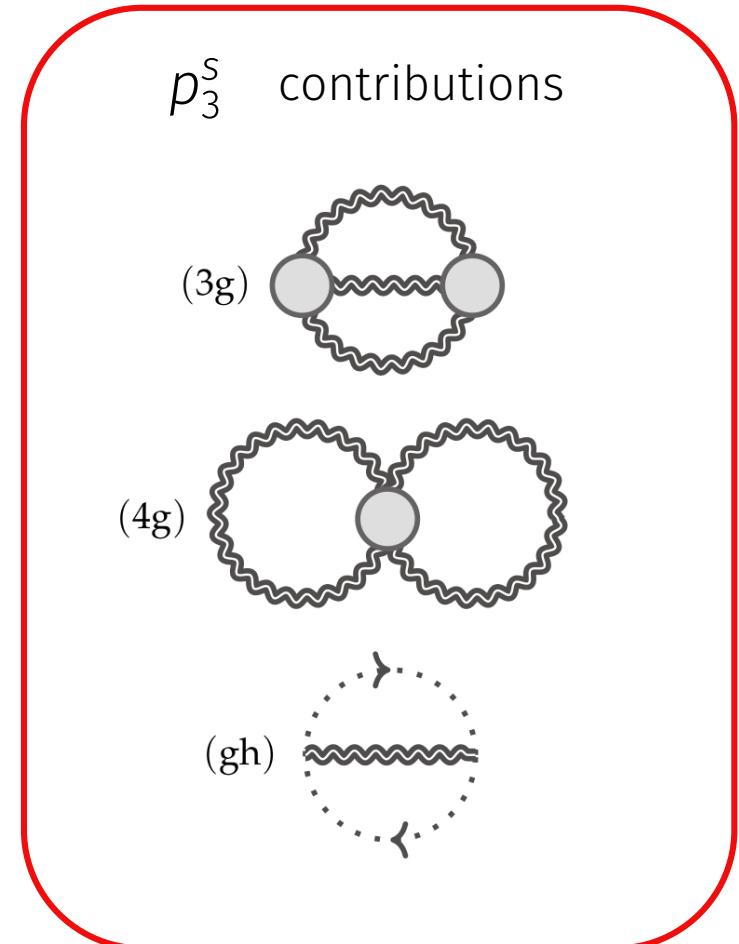


Cold QM approach

Remarks:

- Examined in [Andersen et al. Phys.Rev.D 66 (2002)], but we need the full unexpanded computation
- Here, don't need full quark kinematics, just $\Gamma_{\text{HTL}}^{\mu\nu}(\hat{P})$
- Hence, radial integrals over $|K|, |P|$ in 2-loop HTL diagrams are simpler

$$\alpha_s^3 p_3^s = \frac{\alpha_s N_c d_A m_E^4}{(8\pi)^2} \left(\frac{m_E}{\Lambda_h} \right)^{-4\epsilon} \left(\frac{p_{-2}}{(2\epsilon)^2} + \frac{p_{-1}}{2\epsilon} + p_0 \right)$$

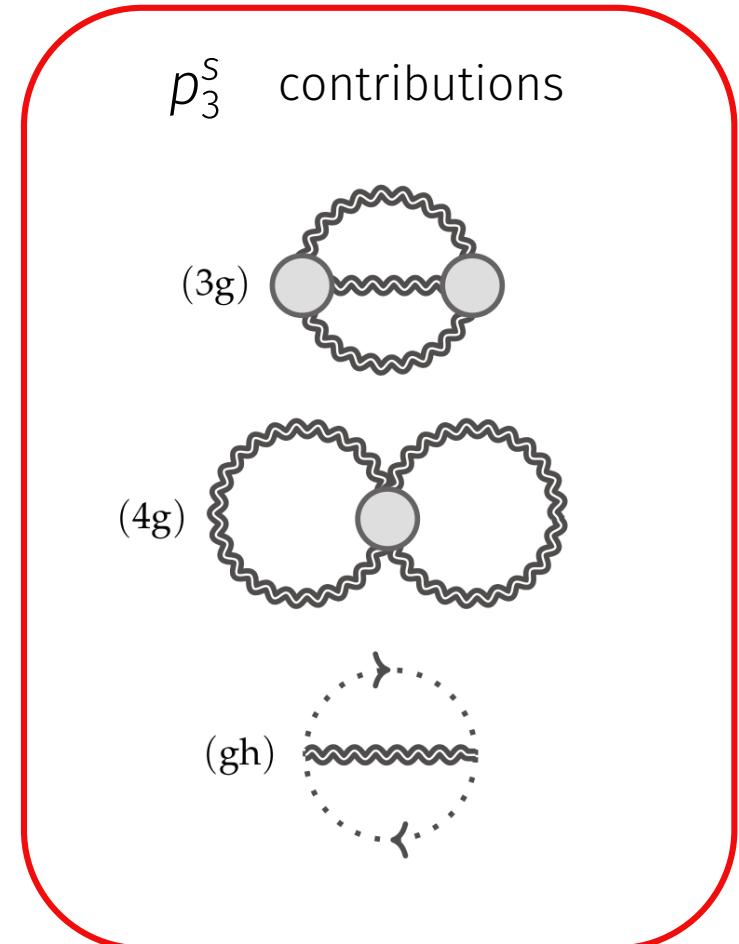


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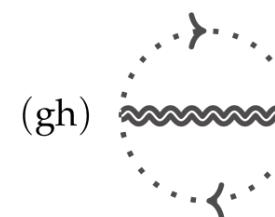
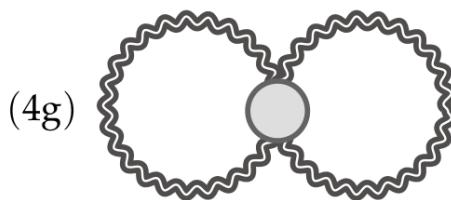
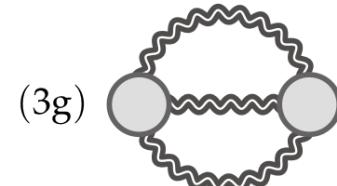
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TG, Kurkela, Romatschke,
Säppi, Vuorinen PRL (2018)

p_3^s contributions



Results

Structure of the soft results

Freedman,
McLerran
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Expanded, gives:

UV div.s cancel with IR div.s from other regions

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Finite results can be plotted and compared to
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$$+ \left. p_0 - 2p_1 \log(m_E/\Lambda_h) + 2p_2 \log^2(m_E/\Lambda_h) \right)$$

Can set Λ_h using
PMS prescription!

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p_{-1}^{NNLO}	1
p_0^{NNLO}	1.17201
p_{-2}	$11/(6\pi)$
p_{-1}	$1.50731(19)$
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Pure-glue beta
function! Running of
 m_E from N2LO

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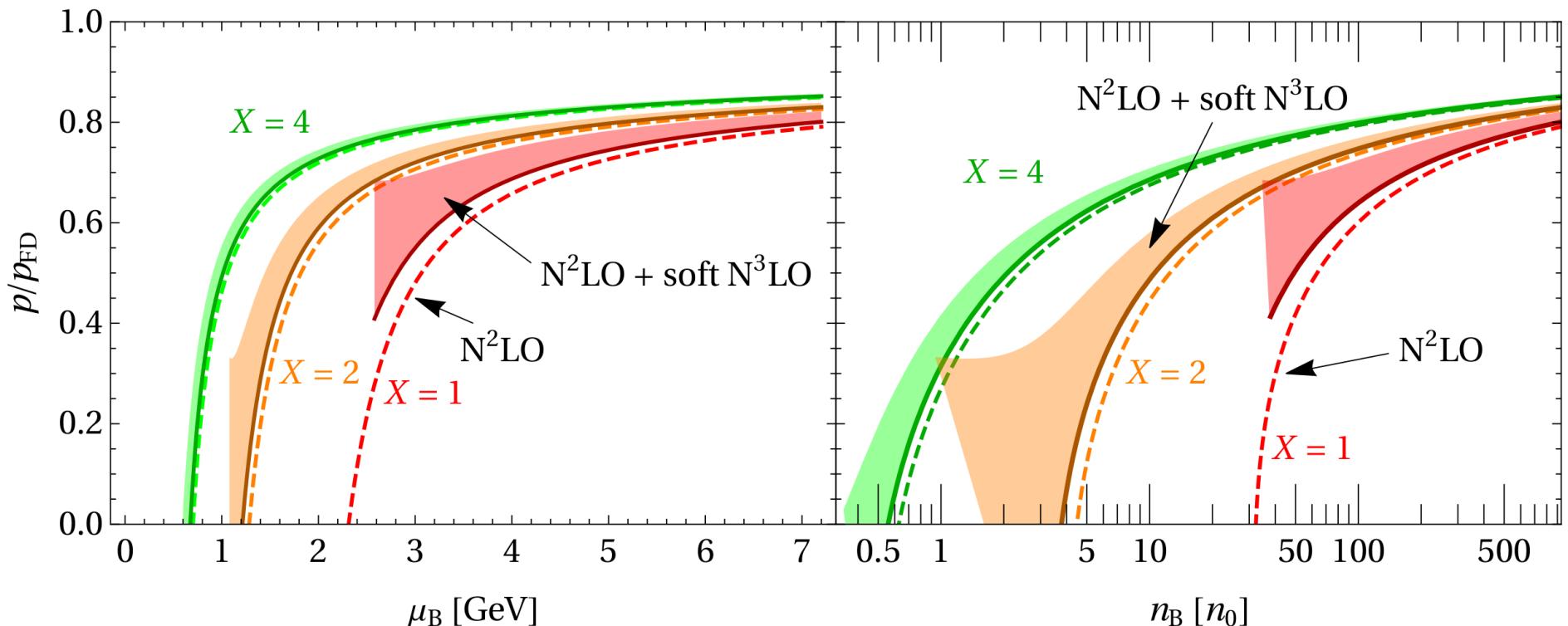
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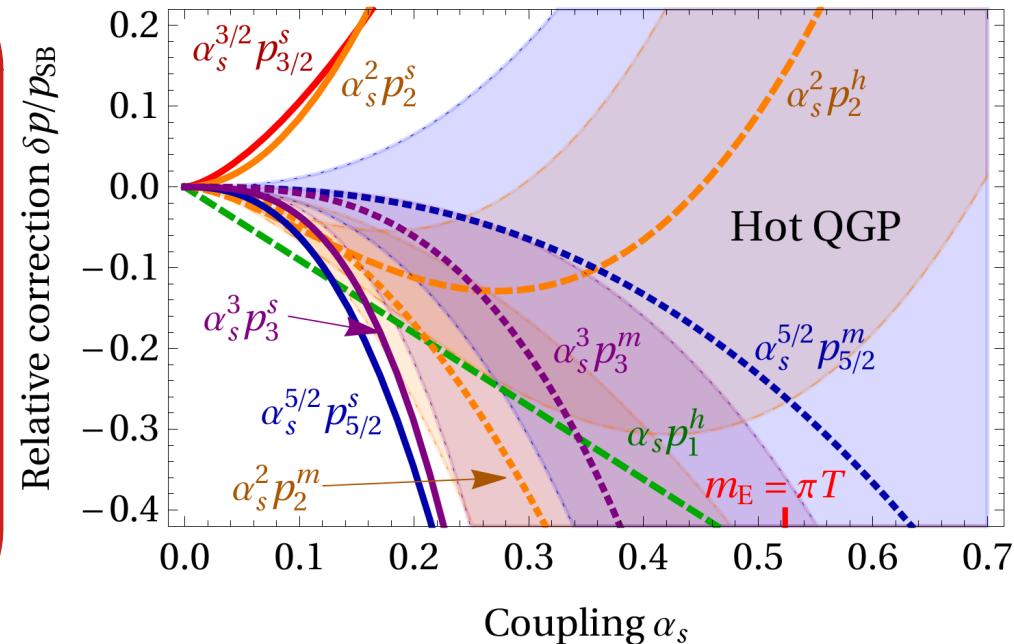
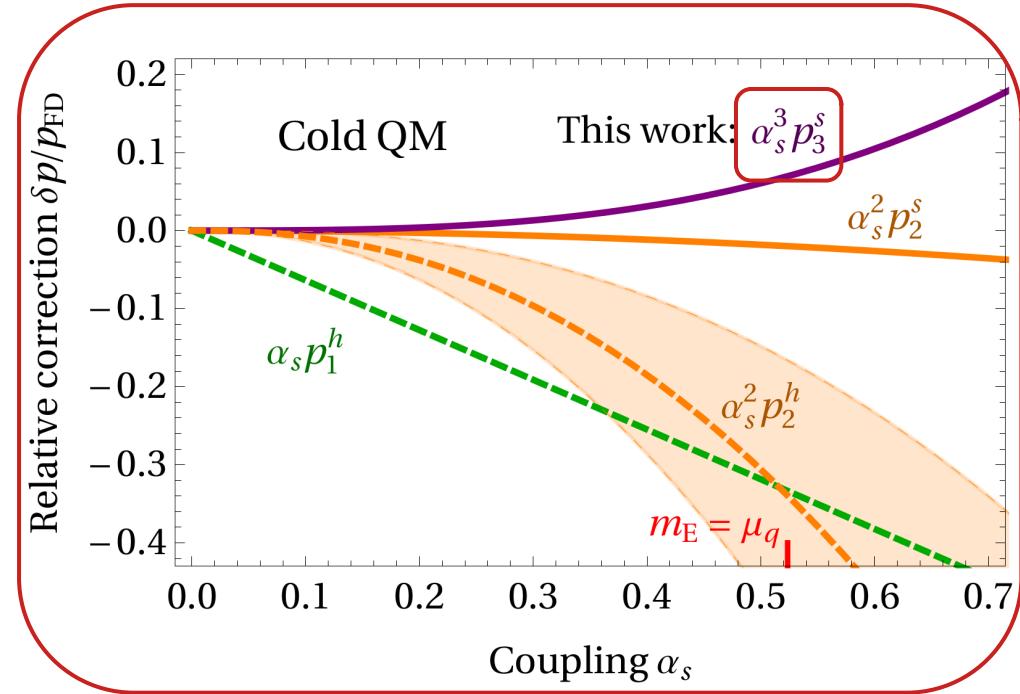
* Makes sense b/c
 p_{-2} comes from UV
limit of soft, pure-
glue, theory

Plots of soft results: 1/2



- * Shifts EOS to slightly higher pressures;
- * Slightly decreases scale dependence;
may improve NS applications: $40n_0 \rightarrow 35n_0$?

Plots of soft results: 2/2



Cold QM perturbatively well behaved – better than hot QGP

Conclusions + outlook

- Have improved the EOS of perturbative cold QM – **new state-of-the-art result**
- Result shows fine convergence of the soft sector; contrary to high T behavior. **Perturbation theory may be more powerful in cold QM.**
- Appearance of β_0 : **Leading term in soft sector simple to all orders?**
- Mixed “ m ” in progress, will improve EOS further. (PMS on renormalization scale too?)
- **Constitutes first full calculation of the interactions between soft, screened gluons in QCD thermodynamics**

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Thank you for your attention!

Backup Slides

Extracting UV divergences: steps

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1. Isolate UV divergences by looking for: $m_E^4 \times (\text{dimensionless int.})$; power-count to extract them

$$m_E^4 \int_{\hat{K}, \hat{P}} \int_{|K|, |P|} f(|K|, |P|, \Pi(\hat{K}), \Pi(\hat{P}))$$

Extracting UV divergences: steps

1. Isolate UV divergences by looking for: $m_E^4 \times (\text{dimensionless int.})$; power-count to extract them
2. $(|P|, |K|)$ to Euclid. polar coordinates: (X, χ) . Radial integral gives one divergence, sometimes angle gives another. Often both tractable.

$$\begin{aligned} m_E^4 \int_{\hat{K}, \hat{P}} \int_{|K|, |P|} f(|K|, |P|, \Pi(\hat{K}), \Pi(\hat{P})) &\mapsto m_E^4 \int_{\hat{K}, \hat{P}} \int_{X, \chi} f(X, \chi, \hat{K}, \hat{P}, \Pi(\hat{K}), \Pi(\hat{P})) \\ &\mapsto \frac{m_E^4}{\varepsilon^2} \int_{\hat{K}, \hat{P}} g(\hat{K}, \hat{P}, \Pi(\hat{K}), \Pi(\hat{P})) + O\left(\frac{1}{\varepsilon}\right) \end{aligned}$$

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2. $(|P|, |K|)$ to Euclid. polar coordinates: (X, χ) . Radial integral gives one divergence, sometimes angle gives another. Often both tractable.
3. Do tensor reduction to simplify remaining contractions between $\hat{P} \cdot \hat{K}$, since we are averaging over remaining angles

$$\begin{aligned} m_E^4 \int_{\hat{K}, \hat{P}} \int_{|K|, |P|} f(|K|, |P|, \Pi(\hat{K}), \Pi(\hat{P})) &\mapsto m_E^4 \int_{\hat{K}, \hat{P}} \int_{X, \chi} f(X, \chi, \hat{K}, \hat{P}, \Pi(\hat{K}), \Pi(\hat{P})) \\ &\mapsto \frac{m_E^4}{\varepsilon^2} \int_{\hat{K}, \hat{P}} g(\hat{K}, \hat{P}, \Pi(\hat{K}), \Pi(\hat{P})) + O\left(\frac{1}{\varepsilon}\right) \\ &\mapsto \frac{m_E^4}{\varepsilon^2} \int_{\hat{K}, \hat{P}} h(\Pi(\hat{K}), \Pi(\hat{P})) + O\left(\frac{1}{\varepsilon}\right) \end{aligned}$$

Non-UV parts

Also need to deal with nasty HTL vertex corrections:

1. Introduce Feynman parametrization
2. Use extensive Ward identities
3. Analytically perform some angular integrals
4. Numerically compute final 6d numerical integral using Monte-Carlo integration

Quick aside about logarithms

Ambiguity in soft/hard split ($m_E \ll K \ll \mu$) gives logarithmic sensitivity to a **factorization mass scale Λ_h** ;

When summing over all kinematic regions, gives logs of coupling:

$$\begin{aligned} p_2^h + p_2^s &= p_2^{\text{LL}} \log \alpha_s + p_2^{\text{const}}, \\ p_3^h + p_3^s + p_3^m &= p_3^{\text{LL}} (\log \alpha_s)^2 + p_3^{\text{NLL}} \log \alpha_s + p_3^{\text{const}}. \end{aligned}$$

Each individually depend on new mass scale Λ_h ; cancels out of sum

TG, Kurkela, Romatschke, Säppi, Vuorinen PRL (2018)