

Constraints on realistic EoS

IST EoS

Modelling of neutron stars

Conclusions

The unified equation of state with induced surface tension consistent with astrophysical, gravitational, high- and low- energy nuclear physics data

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Strongly Interacting Matter Phase Diagram

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Constraints on the EoS



General Requirements

- causality
- thermodynamic consistency
- multicomponent character (n, p, e, ...)

- electric neutrality
- β-equilibrium
- realistic interaction between the constituents



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EoS with hard core repulsion

• Hard core reduces volume available for motion of particles by $V_{excl} = Nb$

 $V \rightarrow V - V_{excl} \Rightarrow \overbrace{\rho = nT}^{ideal \ gas} \stackrel{EoS}{=} \frac{NT}{V} \rightarrow \frac{NT}{V - V_{excl}} = \frac{NT}{V - Nb} = \underbrace{\frac{nT}{nT}}_{1 - nb}$

- Van der Waals EoS (b = const) in the Grand Canonical Ensemble
 - $\begin{cases} p = p(T, \mu) \\ n = \frac{\partial p}{\partial \mu} \end{cases} \Rightarrow p = T \int_{\vec{k}} \exp\left(\frac{\mu pb \sqrt{m^2 + k^2}}{T}\right) = p_{id}(T, \mu pb)$
 - Excluded volume (per particle) depends on density ($b \neq const$)



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The Induced Surface Tension (IST) EoS

Quantitites of the Boltzmann ideal gas

$$p = nT$$
, $n = \sum_{i} n_i^{id}$, $n_i^{id} = \frac{p_i^{id}}{T}$

Virial expansion for one particle species

$$\frac{p}{T} = n + a_2 n^2 + a_3 n^3 + a_4 n^3 + \dots$$

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- Ideal gas: $a_2 = 0, a_3 = 0, a_4 = 0, ...$ Hard spheres: $a_2 = 4v, a_3 = 10v^2, a_4 = 18.365v^3, ...$ R.K.Pathria, Statistical Mechanics, Pergamon Press, Oxford, 1972
- Excluded volume at low density $(v_i = \frac{4\pi}{3}R_i^3 \text{ and } s_i = 4\pi R_i^2)$

$$a_2^{ij} = \frac{1}{2} \cdot \frac{4\pi}{3} (R_i + R_j)^3 = \frac{1}{2} \cdot (v_i + s_i R_j + R_i s_j + v_j)$$



Virial expansion for many particle species

$$\frac{p}{T} = \overbrace{\sum_{i} n_{i}^{id}}^{\simeq n} - \overbrace{\sum_{i,j} a_{2}^{ij} n_{i}^{id} n_{j}^{id}}^{\simeq a_{2}n^{2}} + \dots = \sum_{i} \frac{p_{i}^{id}}{T} \left(1 - v_{i} \sum_{j} \frac{p_{j}^{id}}{T} - s_{i} \sum_{j} \frac{p_{j}^{id}}{T} R_{j}\right) + \dots$$



EoS with induced surface tension (IST)

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lpha accounts for not uniqueness of extrapolation to high densities

High density extrapolation (gives exponentials)

$$\begin{cases} p = \sum_{i} p_{i}^{id} \exp\left(-\frac{pv_{i}+\Sigma s_{i}}{T}\right) \\ \Sigma = \sum_{i} p_{i}^{id} \exp\left(-\frac{pv_{i}+\alpha\Sigma s_{i}}{T}\right) R_{i} \end{cases} \rightarrow \begin{cases} p = \sum_{i} p_{i}^{id} (\mu_{i}-pv_{i}-\Sigma s_{i}) \\ \Sigma = \sum_{i} p_{i}^{id} (\mu_{i}-pv_{i}-\alpha\Sigma s_{i}) R_{i} \end{cases}$$

VS, A. Ivanytskyi, K. Bugaev, I. Mishustin, NPA 924, 24 (2014)



Physical Origin of the Induced Surface Tension

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Vacuum

attraction of constituents \Rightarrow eigen surface tension







- Hard core repulsion only in part is accounted by eigen volume
- The rest corresponds to surface tension and curvature tension Curvature tension can be accounted explicitly or implicitly
- Physical clusters tend to have spherical (in average) shape



Determination of $\boldsymbol{\alpha}$

• One component EoS with IST and $\alpha > 1$

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$$\begin{cases} p = p^{id} \exp\left(-\frac{pv + \Sigma s}{T}\right) \\ \Sigma = p^{id} \exp\left(-\frac{pv + \alpha \Sigma s}{T}\right) R \end{cases} \Rightarrow \begin{cases} p = p^{id} \exp\left(-\frac{pb}{T}\right) R \\ \Sigma = pR \exp\left(\frac{(1-\alpha)\Sigma s}{T}\right) \end{cases}$$

VVS, A. I. Ivanytskyi, K. A. Bugaev, I. N. Mishustin, Nucl. Phys. A, 924, 24 (2014) Excluded volume: $\frac{b}{v} = 1 + 3e^{\frac{(1-\alpha)\Sigma s}{T}} \rightarrow \begin{cases} 4, & \Sigma \rightarrow 0\\ 1, & \Sigma \rightarrow \infty \end{cases}$

$\alpha > 1$ switches different regimes of excluded volume

■ Virial expansion of one component EoS with IST Second virial coefficient: $a_2 = 4V$ is reproduced always Third virial coefficient: $a_3 = 10V^2 \Rightarrow \alpha = \frac{4}{3}$ a_4 is not reproduce

Fourth virial coefficient: $a_4 \simeq 18.365 V^3 \Rightarrow \alpha \simeq 1.245$ a_3 - reproduced with 16% accuracy

 $\alpha>1.245$ reproduces two (3rd and 4th) virial coefficients



Effect of the IST

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Thermodynamic parameters with and without IST





Nuclear Matter Properties Near the (3)CEP

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FIG. 3. Values of incompressibility constant K_0 and critical temperature T_C , which obey the proton flow constraint are located between the lines ABC and FED. The lines ABC and FED are, respectively, generated by the lower and upper branches of the proton flow constraint. The vertical lines AF, BE, and CD correspond to K_0 values 200 MeV, 250 MeV, and 315 MeV, respectively.

VVS, et al., Nucl. Phys. A, 924, 24 (2014)

A. Ivanytskyi et al., PRC 97, 064905 (2018)



Hadron Resonance Gas Model

- Hadrons with masses \leq 2.5 GeV (widths, strong decays, zero strangeness)
- 111 independent particle ratios measured at 14 energies (from 2.7 GeV to 200 GeV)
- 14 × 4 local parameters $(T, \mu_B, \mu_{13}, \gamma_s)$ + 5 global parameters (hard core radii)



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 $R_b = 0.365 \ fm, \ R_m = 0.42 \ fm, \ R_\pi = 0.15 \ fm, \ R_K = 0.395 \ fm, \ R_\Lambda = 0.085 \ fm$ Overall $\chi^2/dof \simeq 1.038$

> K.A. Bugaev, et al., NPA 970, p. 133-155, (2018) VVS, Ukr. J. Phys. 59, 755 (2014)



Hadron Resonance Gas at ALICE Energies

- 11 independent particle yields, 6 parameters (temperature + 5 hard core radii)
- Overal $\chi^2/dof \simeq 0.89$
- Freeze out temperature $T_{FO} = 148 \pm 7 MeV$



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 $p = \sum_{i}^{all \ particles} [p_{id}(T, \mu_i - pV_i - \Sigma S_i + U_{at} \pm U_{sym}) + p_{id}(\mu_e) - p_{at} + p_{sym}]$ $\Sigma = \sum_{i}^{all \ particles} p_{id}(T, \mu_i - pV_i - \alpha \Sigma S_i + U_0)R_i$

 p_{id} – pressure of the ideal gas for quantum statistics Σ – induced surface tension U_0, α – model parameters

Thermodynamic consistency of the model : $\frac{\partial p_{int}}{\partial n_{id}} = n_{id} \frac{\partial U(n_{id})}{\partial n_{id}}$ Parametrization of the mean field potential : $U_{at} = -C_d^2 n_{id}^{\kappa}$

> VVS, et al., Nucl. Phys. A, 924, 24 (2014) A. Ivanytskyi et al., PRC 97, 064905 (2018)



Mean field interaction for nuclear matter

Thermodynamic consistency provides identity $\frac{\partial p}{\partial u} = n$

x - any quantity (density, asymmetry parameter, ...)

K. A. Bugaev and M. I. Gorenstein, Z. Phys. C 43, 261 (1989)

D.H. Rischke, et al. Z. Phys. C 51, 485 (1991)

$$p(\mu) = p^{id}(\mu - U(x)) + p_{int}(x), \quad p_{int}(x) = \int_0^x dx' x' \frac{\partial U(x')}{\partial x'}$$

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Typical in-medium potential Long range attraction (negative contribution to pressure) $p_{at}(x) = -\frac{\kappa}{m+1}C_d^2 x^{1+\kappa}, \quad x = n_n^{id} + n_p^{id}$

 $\kappa < 1, \ C_d^2$ – fitted to flow constraint and properties of ground state A. Ivanytskyi et al., PRC 97, 064905 (2018)

Repulsion due to symmetry energy (positive contribution to pressure)

binding energy of stable nuclei: $E_{sym} = a_{sym} \frac{(N-Z)^2}{A}$, $a_{sym} = 30 \pm 4 MeV$ $p_{sym}(x) = \frac{A^{sym}x^2}{[1+(B^{sym}x)^2]^2}, \quad x = n_n^{id} - n_p^{id}$

A^{sym}, B^{sym} – fitted to a_{sym} and slope of symmetry energy at ground state



EoS for NS interiors

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- Degrees of freedom: neutrons, protons, electrons
- Quantum statistics is accounted by construction of p_{id} and n_{id}
- Realistic interaction: HC repulsion, MF attraction, symmetry energy repulsion
- Virial coefficients of classic hard spheres are reproduced
- Causal behaviour up to densities where QGP is expected and Flow constraint



VS, I. Lopes, A. Ivanytskyi, APJ, 871, 157 (2019) VS, G. Panotopoulos, I. Lopes, PRD 101, 063025 (2020)



$\ensuremath{\mathsf{M}}\xspace{-}\ensuremath{\mathsf{R}}\xspace$ relation and compactness of $\ensuremath{\mathsf{NS}}\xspace$



VS, I. Lopes, A. Ivanytskyi, APJ 871, 157 (2019) VS, G. Panotopoulos, I. Lopes, PRD 101, 063025 (2020)



New formulation of the IST EoS with density dependent mass

Total pressure and surface tension

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$$p = \sum_{i}^{all \ particles} [p_{id}(T, \mu_i - pV_i - \Sigma S_i \pm U_{sym}) + p_{id}(\mu_e) - p_{at} + p_{sym}]$$

$$\Sigma = \sum_{i}^{all \ particles} p_{id}(T, \mu_i - pV_i - \alpha \Sigma S_i \pm U_{sym} + U_0)R_i$$

Effective masses

$$\left\{ \begin{array}{l} m_{i}^{p}=m_{nucl}-U_{at} \\ m_{i}^{\Sigma}=m_{nucl}-U_{at} \end{array} \right.$$

VS, O. Ivanytskyi, A. Maselli, C. Providência, In preparation 2021



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Summary

IST approach was successfully applied to the description of compressible nuclear matter properties near the (3)CEP

VVS, et al., Nucl. Phys. A, 924, 24 (2014)

• heavy-ion collision data between $\sqrt{S_{NN}} = 2.7$ GeV – 2.76 TeV

NS properties at T=0

VVS, I. Lopes, A. Ivanytskyi, APJ, 871, 157 (2019)

K.A. Bugaev, et al., NPA 970, p. 133-155, (2018)

Advantages of the IST EoS

- can be easily generalized to any number of particle species
- can be formulated to finite temperatures \Rightarrow proto-neutron stars
- provide a unified description of hadron and nuclear matter
- extract the surface tension of nuclear matter \Rightarrow pasta phase
- IST approach can be applied to the description of the deconfinement phase transition

Future prospects

- will be available soon on CompOSE
- add hyperons and heavy leptons
- formulate a hybrid star model with QGP core



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