

Special point and onset of deconfinement in the M-R diagram of neutron stars.

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Overview

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- 2 Properties of the special point
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 - Mass relation
 - Invariance w.r.t. the phase transition construction
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Introduction

What is the special point?

The TOV equation

The TOV equation

$$\frac{dP(r)}{dr} = \frac{(r)M(r) \left(1 + \frac{P(r)}{r}\right) \left(1 + \frac{4}{r} \frac{r^3 P(r)}{M(r)}\right)}{r^2 \left(1 - \frac{2M(r)}{r}\right)};$$

The TOV equation

$$\frac{dP(r)}{dr} = \frac{(r)M(r) \left(1 + \frac{P(r)}{r}\right) \left(1 + \frac{4}{r} \frac{r^3 P(r)}{M(r)}\right)}{r^2 \left(1 - \frac{2M(r)}{r}\right)}; \quad \frac{dM(r)}{dr} = 4\pi r^2 \rho(r)$$

The TOV equation

$$\frac{dP(r)}{dr} = \frac{(r)M(r)}{r^2} \frac{1 + \frac{P(r)}{r}}{1 - \frac{2M(r)}{r}} \left(1 + \frac{4}{r} \frac{r^3 P(r)}{M(r)} \right); \quad \frac{dM(r)}{dr} = 4\pi r^2 \rho(r)$$

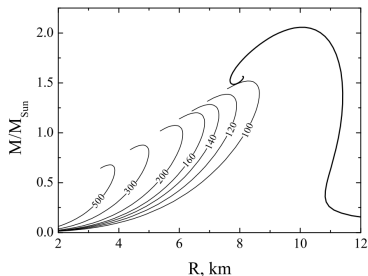


Fig. 1. Mass-radius diagram for a star made of ordinary matter (thick line) and purely quark stars (thin lines). The numbers at the lines indicate the parameter B .

¹Yudin et al., *Astron. Lett.* **40** (2014), 201

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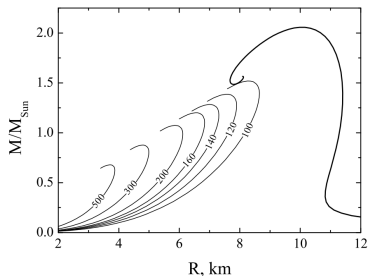


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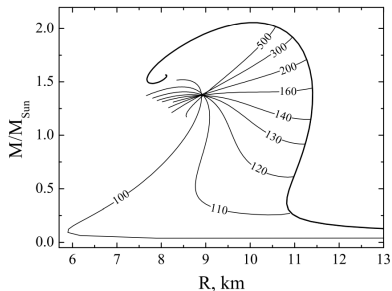


Fig. 2. Mass-radius diagram of hybrid stars for various values of the parameter B

¹Yudin et al., *Astron. Lett.* **40** (2014), 201

Properties of the special point

Hadronic EoS invariance

The constant-speed-of-sound (CSS) model:

²Alford, Han, Prakash, Phys. Rev. **D 88** (2013) no.8, 083013

The constant-speed-of-sound (CSS) model:

– dimensionless baryochemical potential

$$\hat{\mu}_B = \frac{\mu_B}{\text{scale}} = \frac{\mu + B}{A}^{1+(1+\alpha)};$$

– pressure

$$p(\mu_B) = A \hat{\mu}_B^{1+\alpha};$$

– baryon density

$$n_B(\mu_B) = (1 + \alpha) \frac{A}{\text{scale}} \hat{\mu}_B^\alpha;$$

– energy density

$$= B + A \hat{\mu}_B^{1+\alpha};$$

– $p(\mu_B)$ relation: $= p + (1 + \alpha) B$:

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– dimensionless baryochemical potential

$$\hat{\mu}_B = \frac{\mu_B}{\text{scale}} = \frac{\mu_B + B}{A} \quad \mu_B = (1 + \frac{B}{A}) \mu_B;$$

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$$p(\mu_B) = A \hat{\mu}_B^{1+} \quad \mu_B;$$

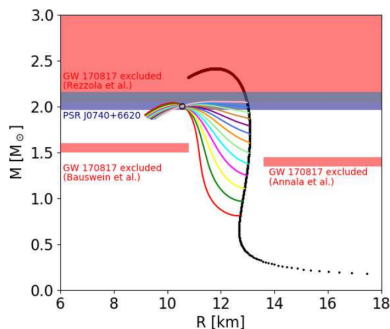
– baryon density

$$n_B(\mu_B) = (1 + \frac{B}{A}) \frac{A}{\text{scale}} \hat{\mu}_B;$$

– energy density

$$= \mu_B + A \hat{\mu}_B^{1+} \quad \mu_B;$$

– $p(\mu_B)$ relation: $= p + (1 + \frac{B}{A}) \mu_B$:



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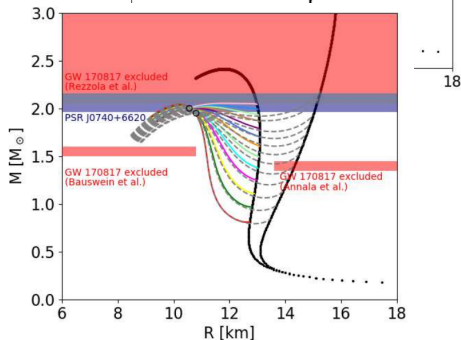
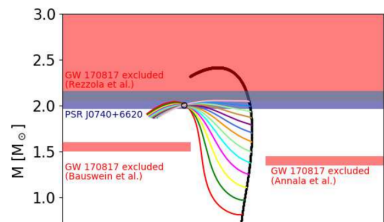
– baryon density

$$n_B(\mu_B) = (1 + \dots) \frac{A}{\text{scale}} \hat{\mu}_B^i;$$

– energy density

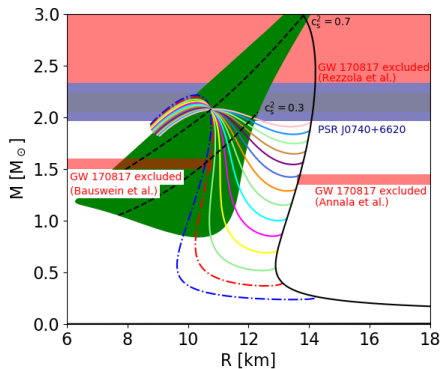
$$= B + A \hat{\mu}_B^{1+} \quad \mu_B;$$

– $p(\mu_B)$ relation: $= p + (1 + \dots) B:$



Properties of the special point

Mass relation



³Cierniak, Blaschke, Eur.Phys.J.ST **229** (2020) no.22-23, 3663-3673

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Properties of the special point

Invariance w.r.t. the phase transition construction

The mixed phase parabolic ansatz:

⁴Abgaryan, et al., Universe (2018), 94

The mixed phase parabolic ansatz:

$$P_M(\mu) = \frac{1}{2}(\mu - \mu_c)^2 + \mu_1(\mu - \mu_c) + (1 + \mu_1)P_c;$$

Gibbs condition for phase equilibrium:

$$\begin{aligned} P_H(\mu_H) &= P_M(\mu_H); \\ P_Q(\mu_Q) &= P_M(\mu_Q); \\ \frac{\partial^k}{\partial \mu^k} P_H(\mu_H) &= \frac{\partial^k}{\partial \mu^k} P_M(\mu_H); \\ \frac{\partial^k}{\partial \mu^k} P_Q(\mu_Q) &= \frac{\partial^k}{\partial \mu^k} P_M(\mu_Q); \end{aligned}$$

Derived parameters ($k = 1$):

$$\mu_1 = \frac{2}{2} \frac{1 + \frac{2(\mu_c - \mu_H)}{Q}}{Q(\frac{\mu_c - \mu_H}{H} - \frac{\mu_c - \mu_Q}{Q})};$$

$$\mu_2 = \frac{2}{2} \frac{1 + \frac{2(\mu_c - \mu_Q)}{H}}{H(\frac{\mu_c - \mu_Q}{H} - \frac{\mu_c - \mu_H}{Q})};$$

$$\mu_1 = n_Q(\mu_c - \mu_Q) - n_H(\mu_c - \mu_H) + P_Q - P_H;$$

$$\mu_2 = n_Q - n_H;$$

⁴Abgaryan, et al., Universe 4 (2018), 94

The mixed phase parabolic ansatz:

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Derived parameters ($k = 1$):

$$\mu_1 = \frac{2}{2} \frac{1 + \frac{2(\mu_c - \mu_H)}{Q - H}}{Q - H};$$

$$\mu_2 = \frac{2}{2} \frac{1 + \frac{2(\mu_c - \mu_Q)}{H - Q}}{H - Q};$$

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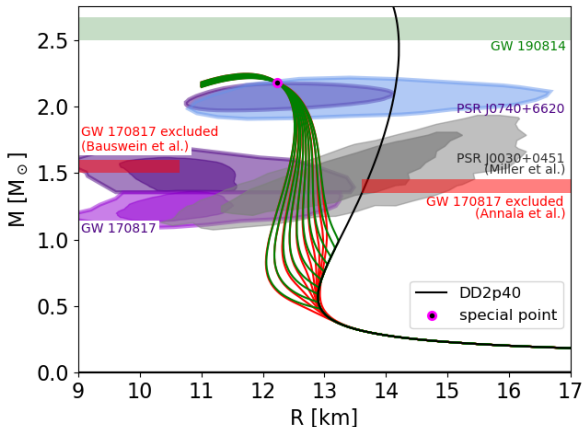
$$\mu_2 = \frac{2}{2(\mu_c - \mu_H) + (\mu_H - \mu_Q)};$$

$$\mu_1 = \eta_Q(\mu_c - \mu_Q) - \eta_H(\mu_c - \mu_H) + P_Q - P_H;$$

$$\mu_2 = \eta_Q - \eta_H;$$

⁴Abgaryan, et al., Universe (2018), 94

Invariance w.r.t. Maxwell > mixed phase construction (pasta phases)



⁵Cierniak, Blaschke, Astron. Nachr. (2021), arXiv: 2106.06986

Invariance w.r.t. Maxwell $>$ interpolation construction (soft - stiff transition)

⁵Cierniak, Blaschke, Astron. Nachr. (2021), arXiv: 2106.06986

Special point locations for constant sound speed

⁵Cierniak, Blaschke, Astron. Nachr. (2021), arXiv: 2106.06986

| c_s^2 | M_{SP} [M] | R_{min} [km] | R_{max} [km] |
|---------|---------------------|-------------------|-------------------|
| 0.35 | 1.82 | - | - |
| 0.40 | 2.07 | 12.18 | 12.29 |
| 0.45 | 2.30 | 11.84 | 13.41 |
| 0.50 | 2.50 | 11.56 | 13.91 |
| 0.55 | 2.68 | 11.30 | 14.20 |
| 0.60 | 2.86 | 11.05 | 14.45 |
| 0.70 | 3.22 | 10.67 | 14.67 |
| 1.00 | 4.00 | 9.95 | 14.84 |

The values of c_s^2 , largest possible M_{SP} and the radii range ($R_{min} - R_{max}$) of a $2 M$ hybrid star.

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The values of c_s^2 , largest possible M_{SP} and the radii range ($R_{min} - R_{max}$) of a $2 M$ hybrid star. **Bold red rows correspond to the nINJL t from [6].**

⁶Antic, ShahrbaF, Blaschke, Grunfeld, arXiv: 2105.00029

Conclusions:

- A special point is a feature of neutron star mass radius relations unique to hybrid stars.
- The special point does not depend on the choice of the hadronic equation of state, phase transition onset density or the type of transition, thus
- The novel NICER PSR J0740+6620 measurement suggests a stiff high density equation of state, in tension with hyperon hadron models. The existence of hybrid stars solves this problem and hints at a possible hybrid neutron star nature of the $2.6 M_{\odot}$ constituent of GW190814.

Outlook:

- Special point of rotating neutron stars?
- Relation between the special point and tidal deformability.
- The special point at a finite temperature.