

University of Michigan & RIKEN (Theoretical Quantum Physics + iTHEMS)

ENRICO RINALDI

**COMPOSITE DM &
GW SIGNATURES**



Lattice Strong Dynamics collaboration

Argonne: Jin, Osborn

Liverpool: Schaich

Bern: Gasbarro

LLNL: Vranas, Howarth

Boston: Brower, Rebbi

UC Davis: Kiskis

Nvidia: Weinberg

Yale: Appelquist,
Fleming, Cushman

Colorado: Neil,
Hasenfratz

Oregon: Kribs

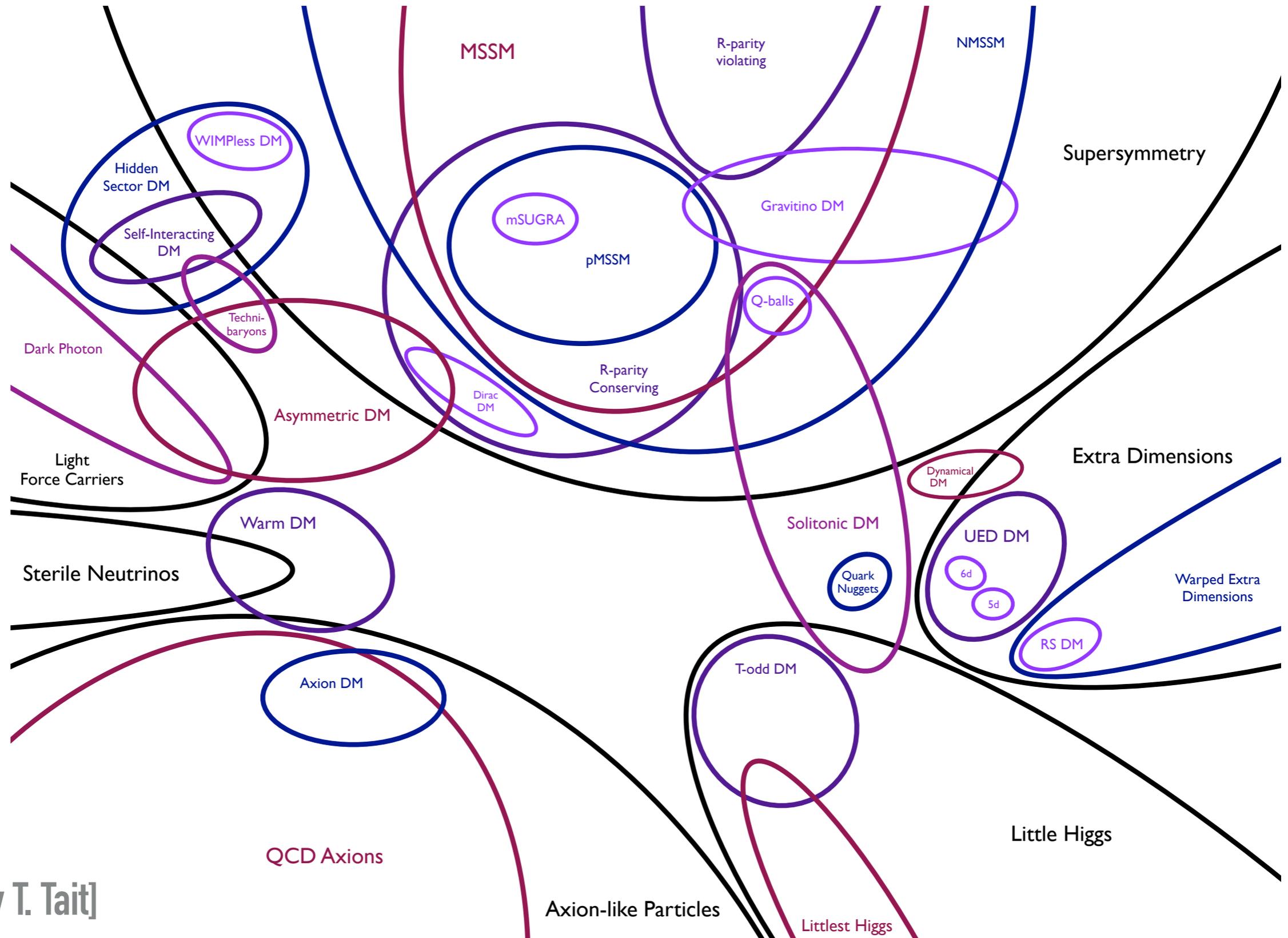
Siegen: Witzel

RIKEN: ER



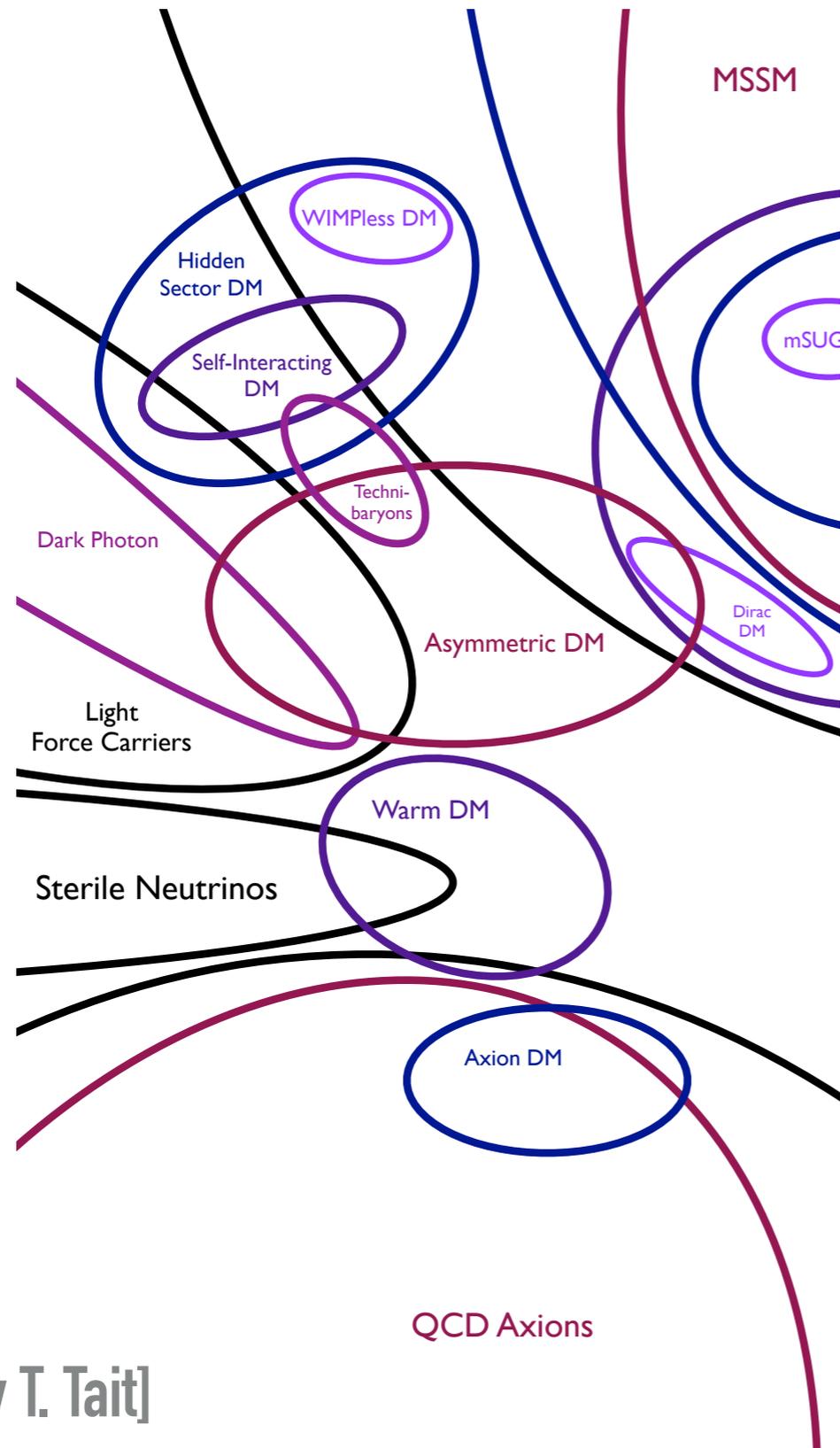
dark
matter ?

THE DARK MATTER THEORY LANDSCAPE



[picture by T. Tait]

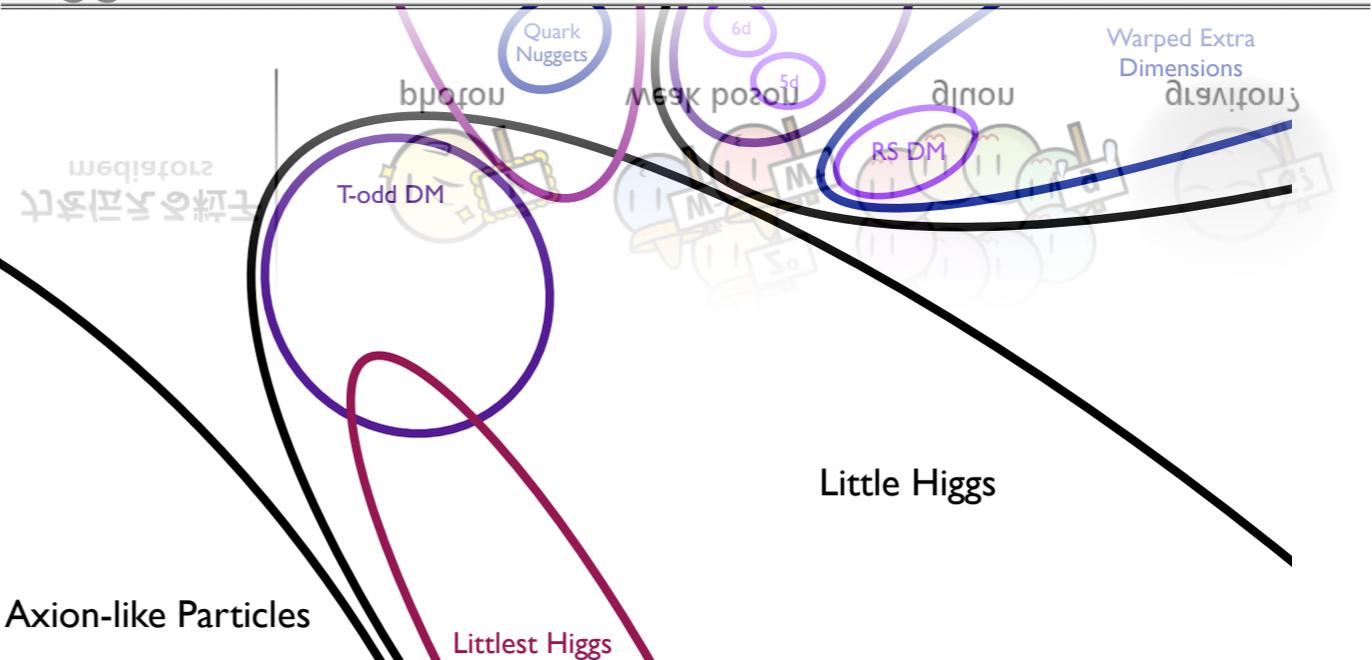
THE DARK MATTER THEORY LANDSCAPE



基本相互作用
Fundamental Interaction

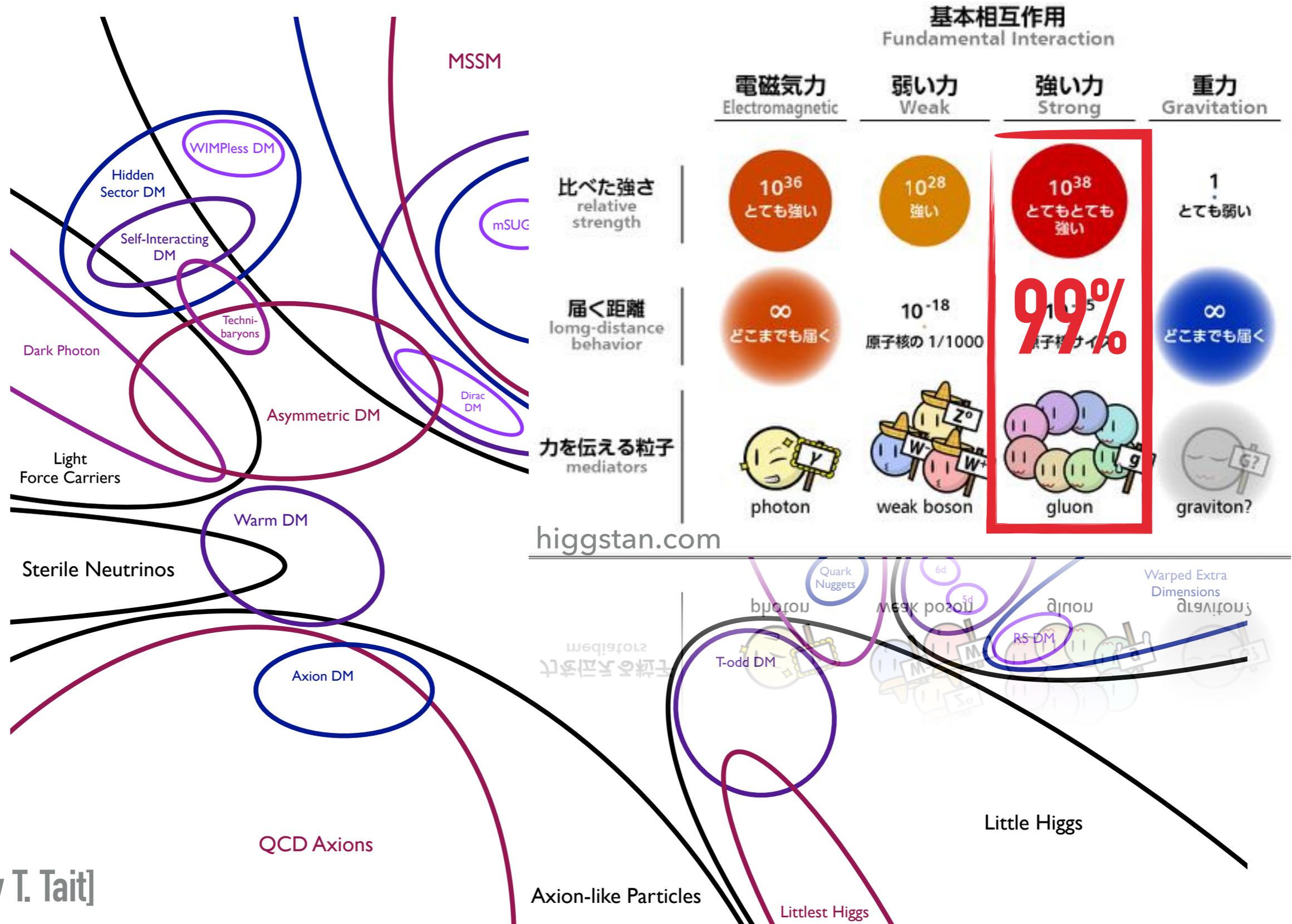
	電磁気力 Electromagnetic	弱い力 Weak	強い力 Strong	重力 Gravitation
比べた強さ relative strength	10^{36} とても強い	10^{28} 強い	10^{38} とてもとても強い	1 とても弱い
届く距離 long-distance behavior	∞ どこまでも届く	10^{-18} 原子核の 1/1000	10^{-15} 原子核サイズ	∞ どこまでも届く
力を伝える粒子 mediators	photon	weak boson	gluon	graviton?

higgstan.com



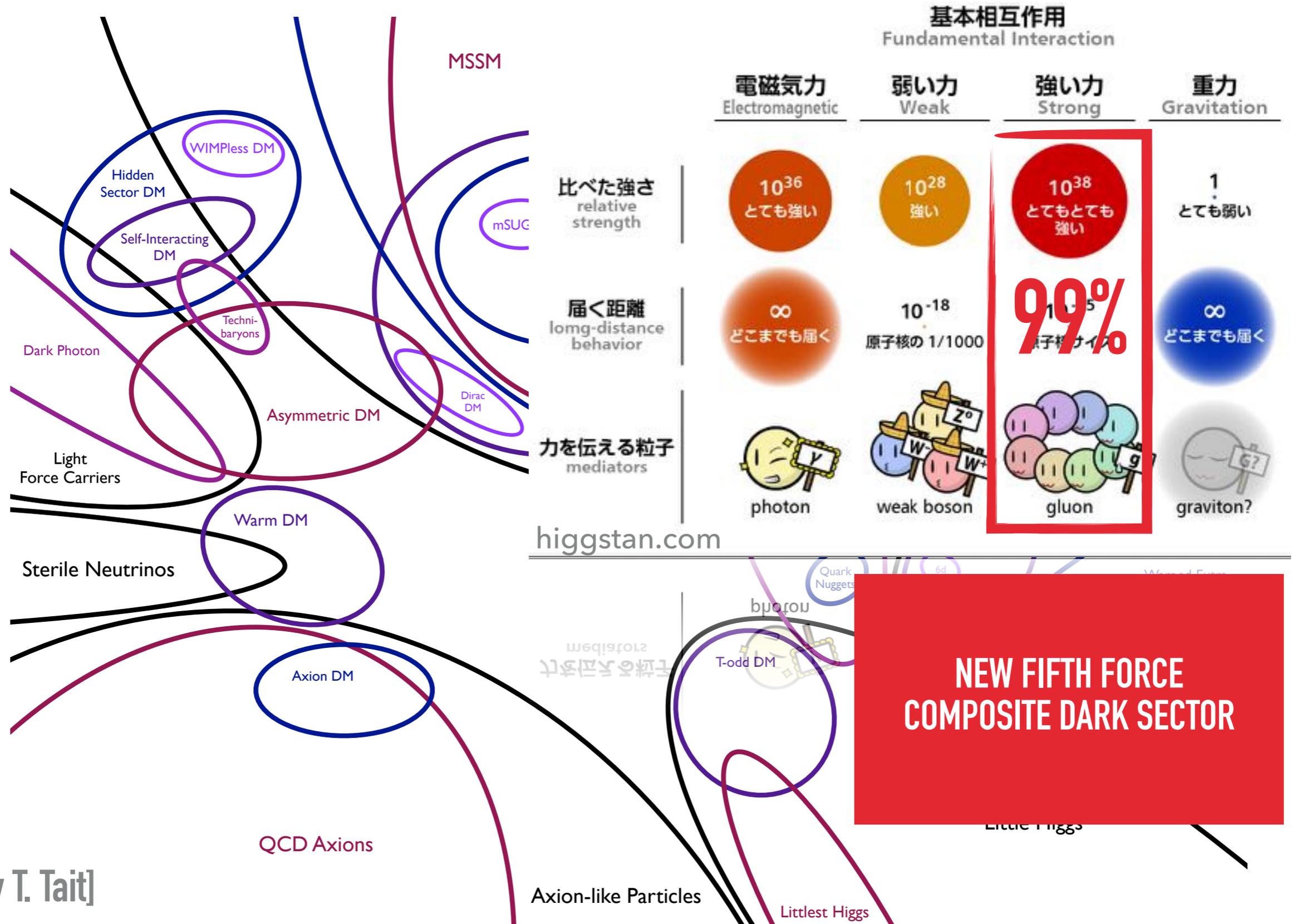
[picture by T. Tait]

THE DARK MATTER THEORY LANDSCAPE



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THE DARK MATTER THEORY LANDSCAPE



**NEW FIFTH FORCE
COMPOSITE DARK SECTOR**

[picture by T. Tait]

DETECTING COMPOSITE DARK SECTORS AND NEW PHYSICS: THE TRIAD

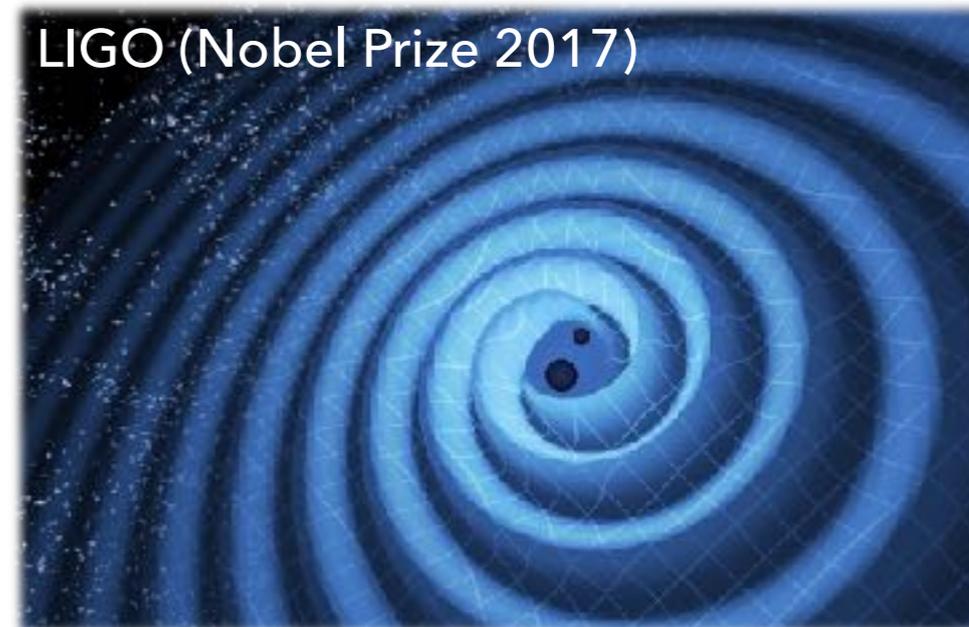
**UNDERSTAND FIFTH FORCE TO GUIDE EXPERIMENTAL
DISCOVERY**

UNDERSTAND FIFTH FORCE TO GUIDE EXPERIMENTAL
DISCOVERY



DARK PHASE TRANSITION GENERATING
GRAVITATIONAL WAVES

LIGO (Nobel Prize 2017)

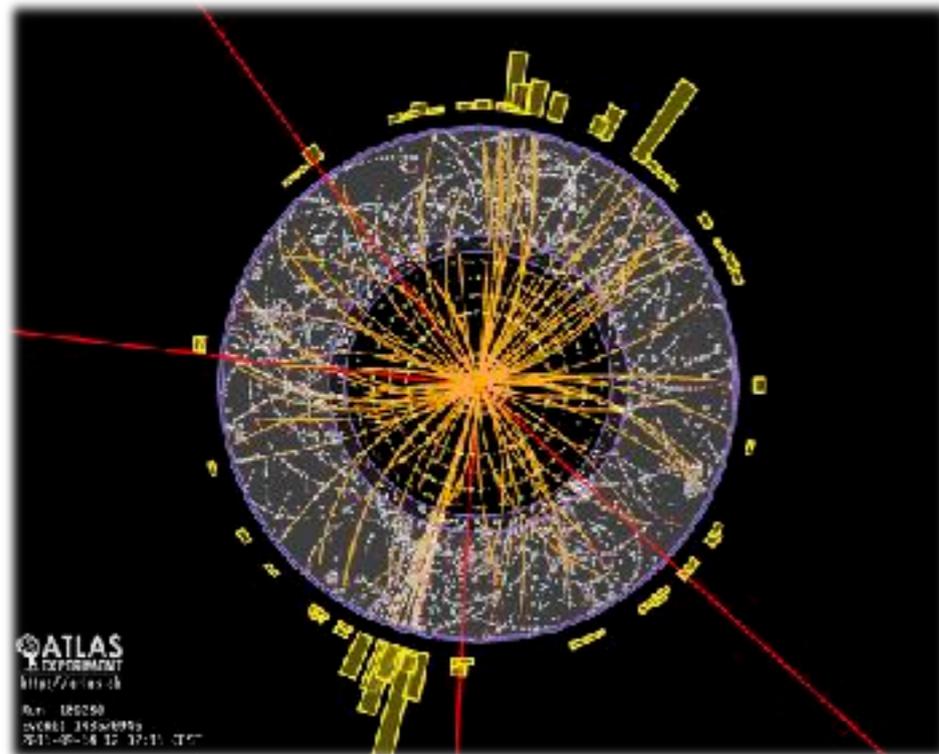


DETECTING COMPOSITE DARK SECTORS AND NEW PHYSICS: THE TRIAD

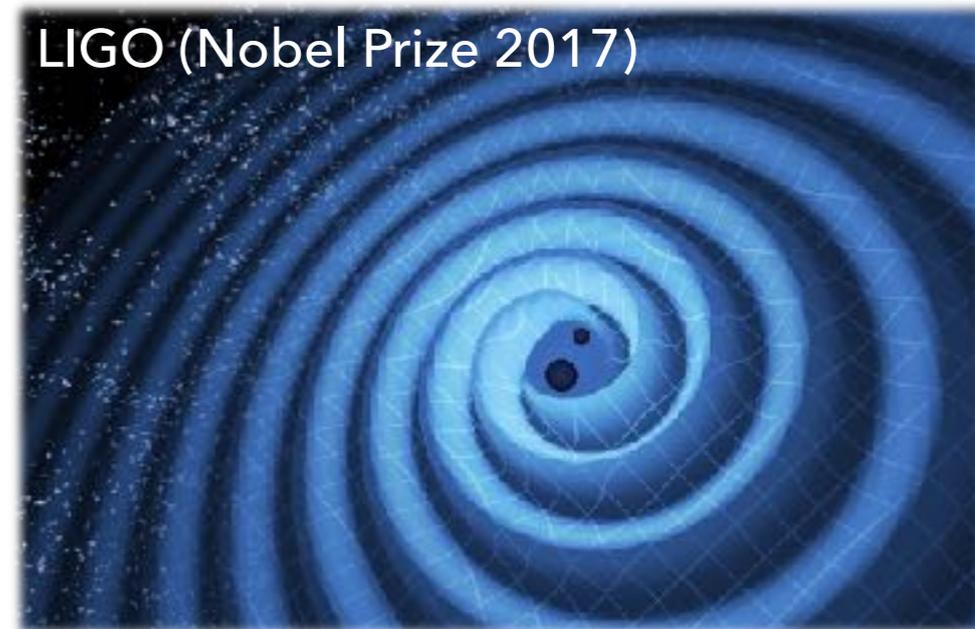
UNDERSTAND FIFTH FORCE TO GUIDE EXPERIMENTAL DISCOVERY

DARK SECTOR PARTICLES PRODUCED AT HIGH-ENERGY COLLIDERS

DARK PHASE TRANSITION GENERATING GRAVITATIONAL WAVES



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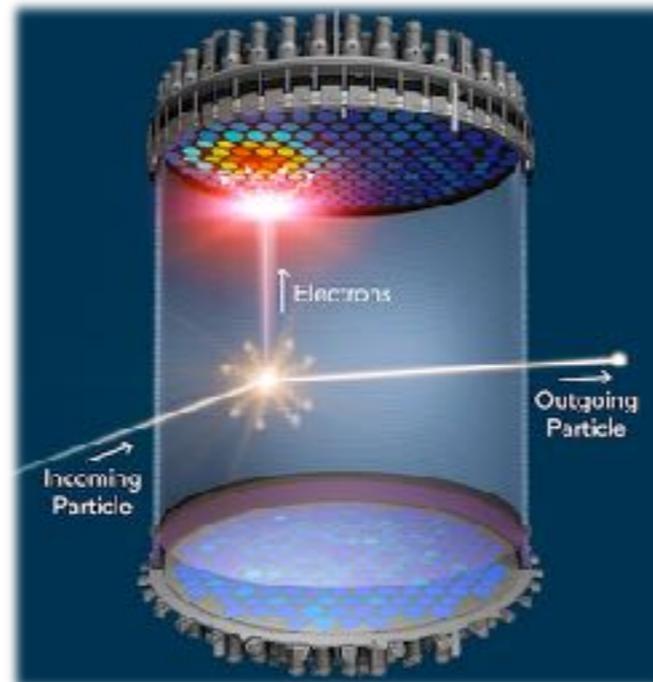
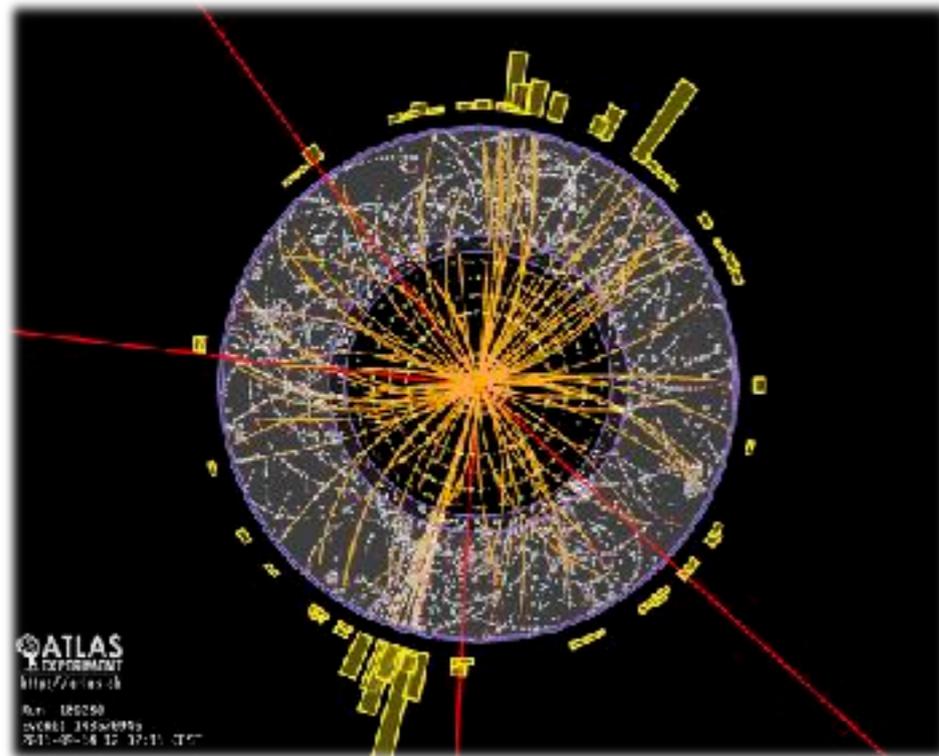
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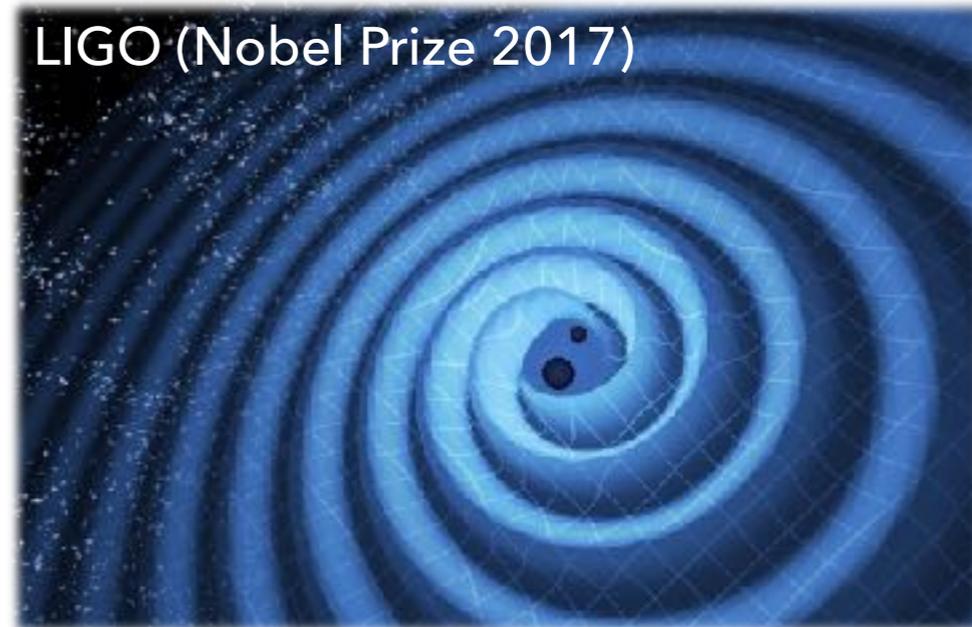
DARK SECTOR PARTICLES PRODUCED AT HIGH-ENERGY COLLIDERS

DIRECT DETECTION THROUGH DARK AND NUCLEAR FORM FACTORS

DARK PHASE TRANSITION GENERATING GRAVITATIONAL WAVES



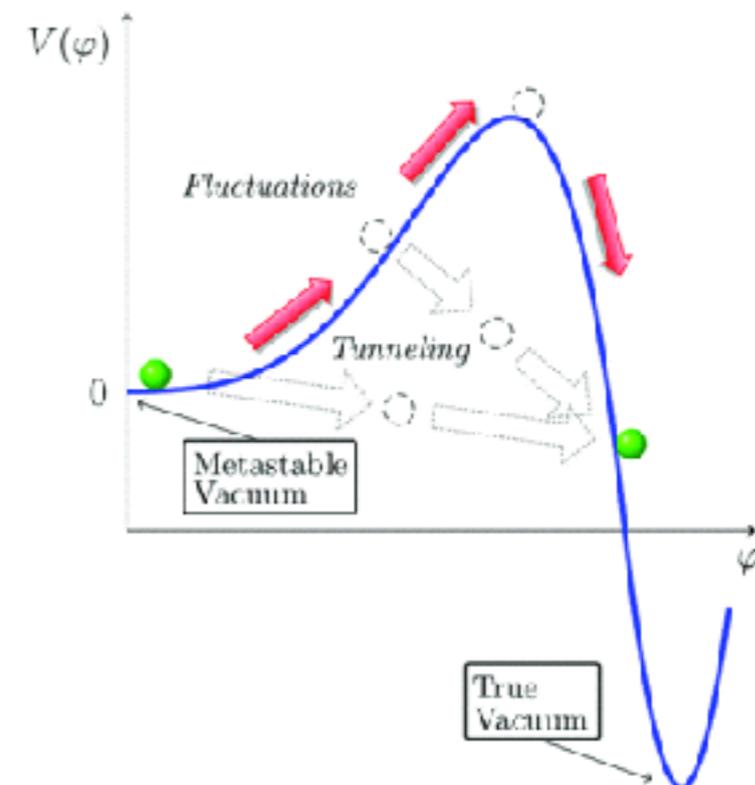
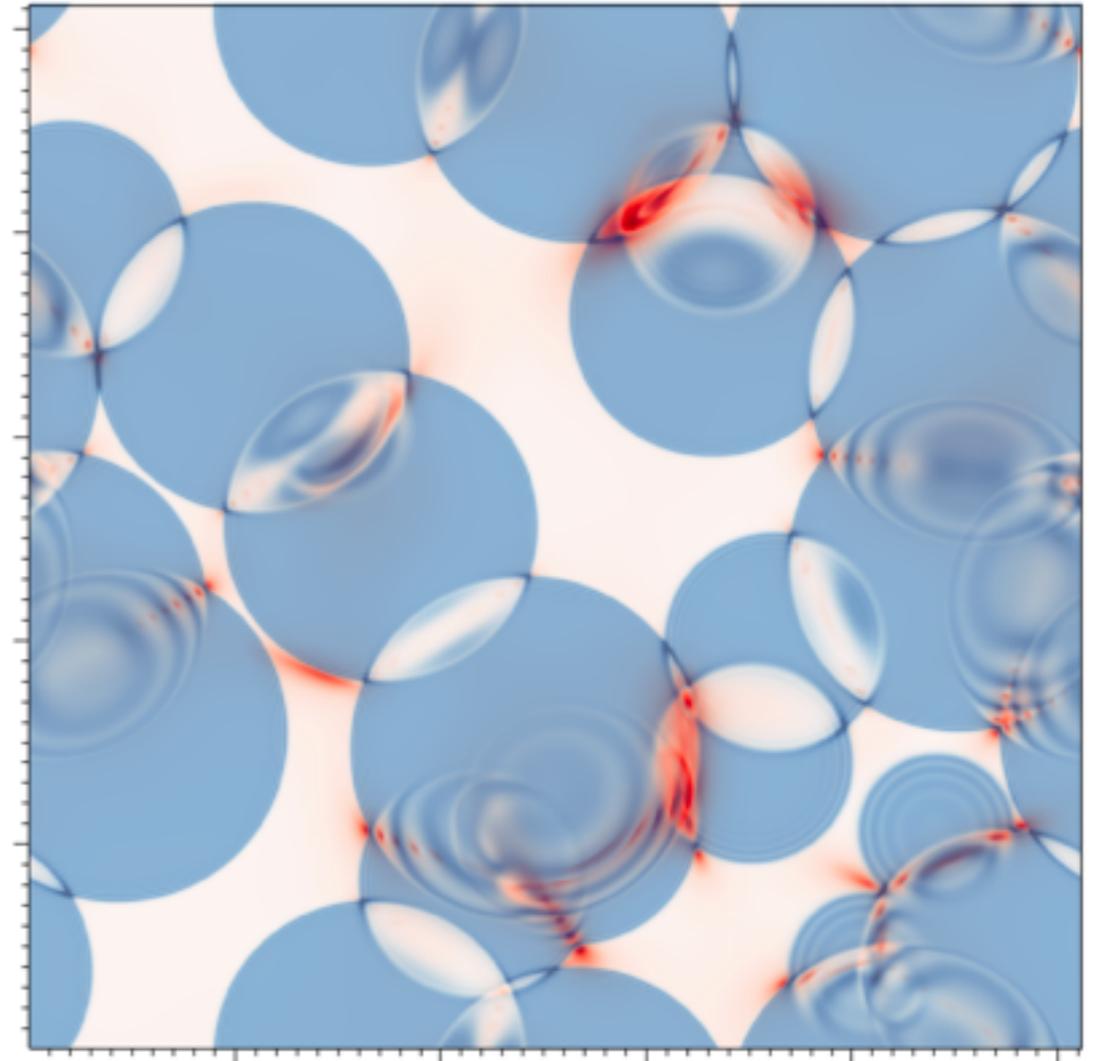
LIGO (Nobel Prize 2017)



Phase Transitions (PT) are everywhere in nature!

They can be *cosmological*

Gravitational waves from vacuum first-order phase transitions

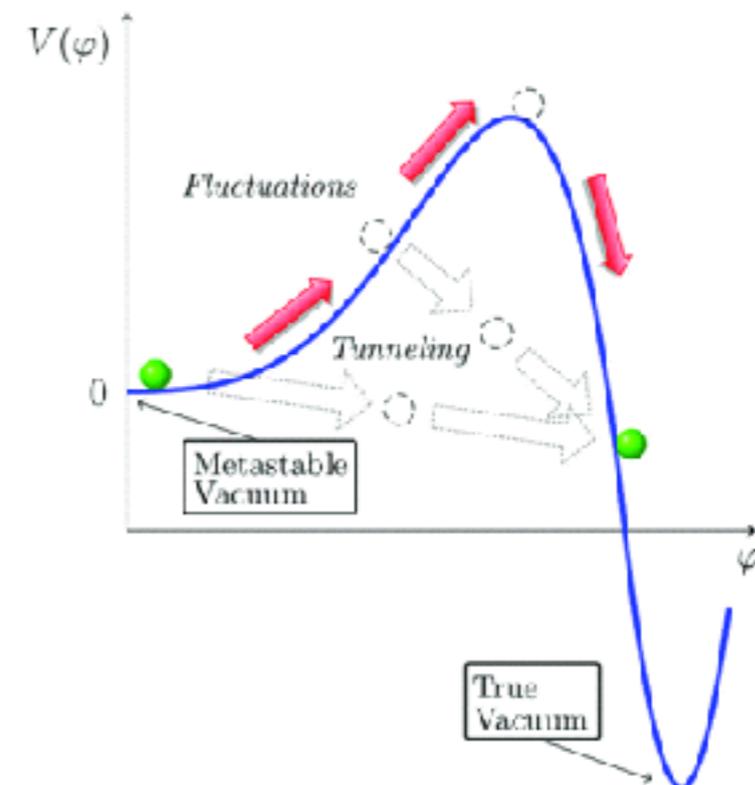
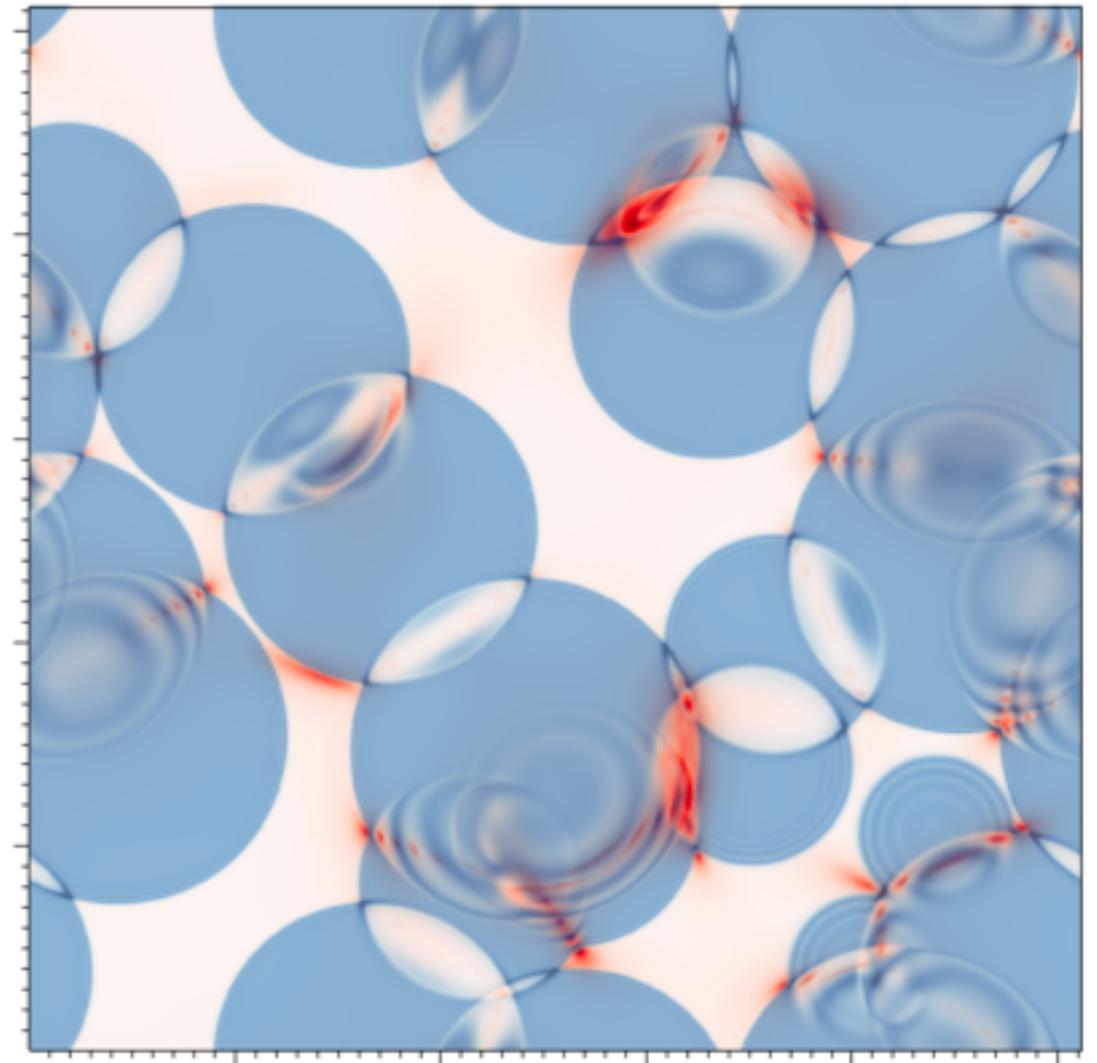


Cosmological Aspects of Higgs Vacuum Metastability

Phase Transitions (PT) are everywhere in nature!

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Focus on 1st order PTs: the universe changes from a metastable high energy (symmetric) phase to a stable lower energy (broken) phase.

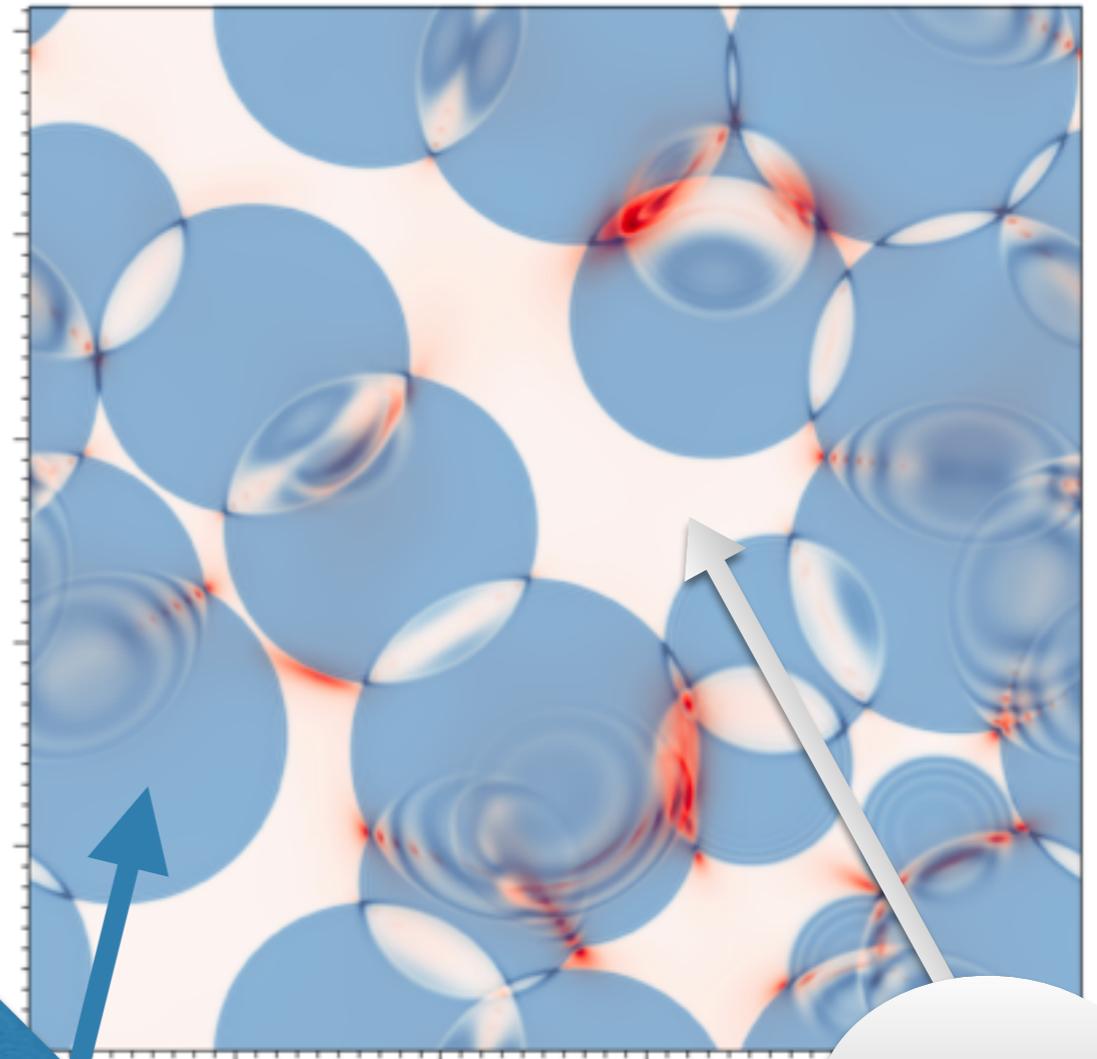


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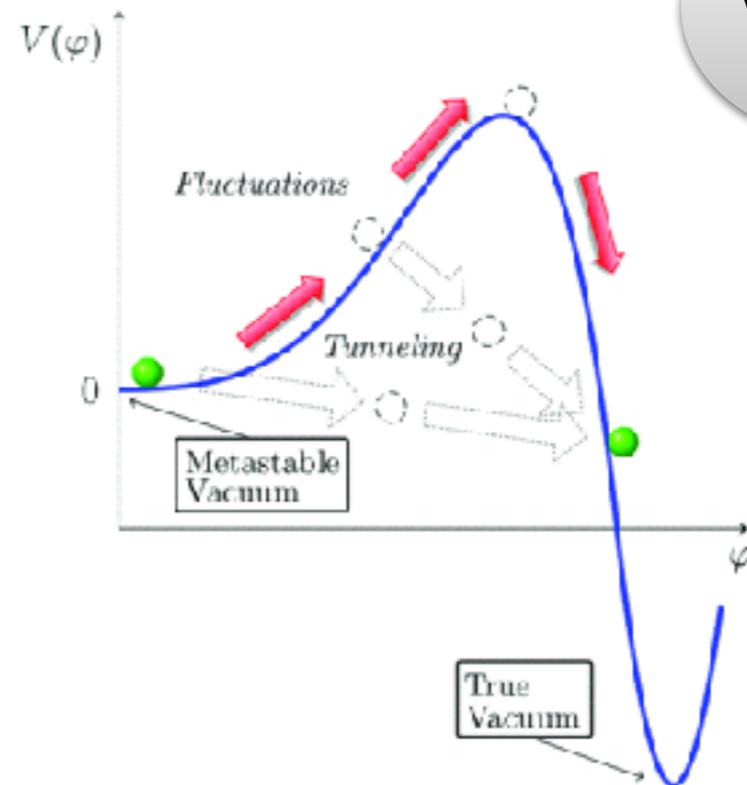
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VEV

Vacuum



EARLY UNIVERSE TRANSITIONS

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QCD confinement-
deconfinement

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crossover...maybe 1st order at finite density...

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Dark sector transition

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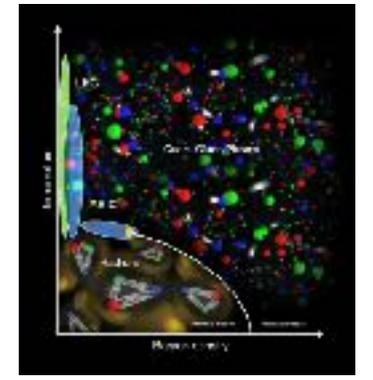
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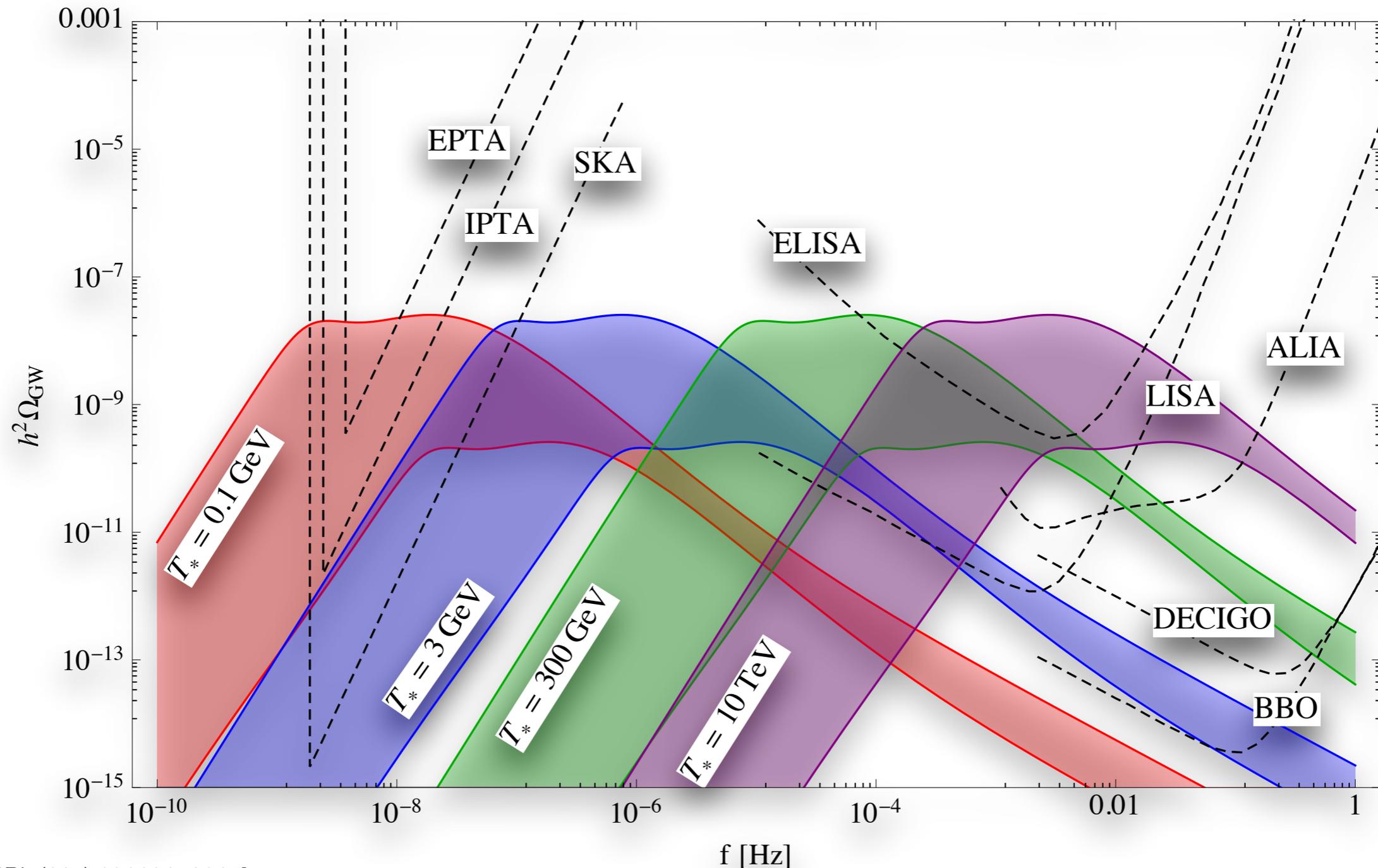
Dark sector transition

maybe strong 1st order if N is large...



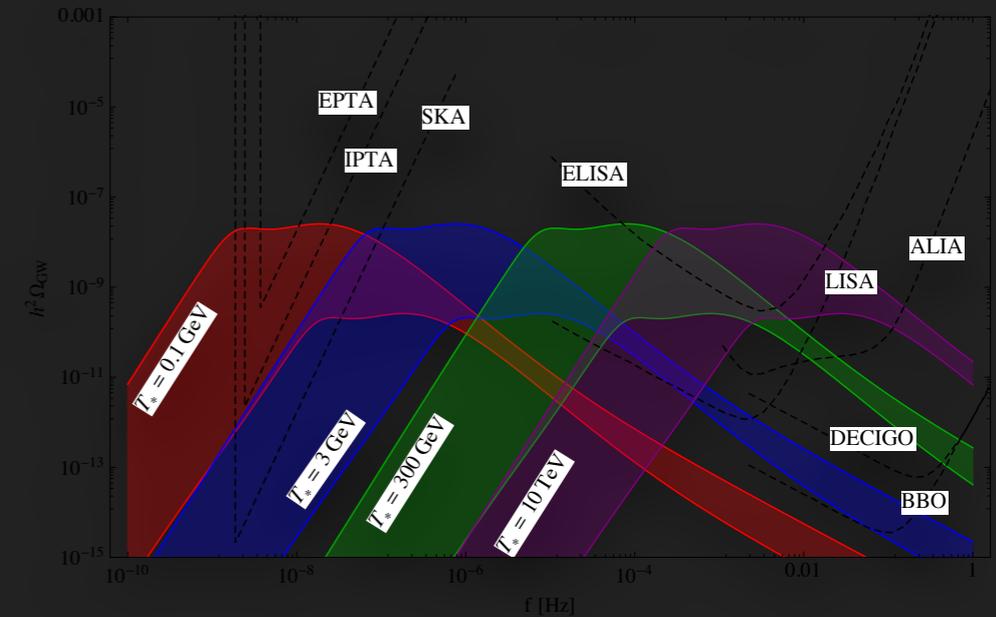
Dark sector gravitational wave signatures

- ◆ Spectrum of GW from a deconfinement 1st order phase transition in the dark sector



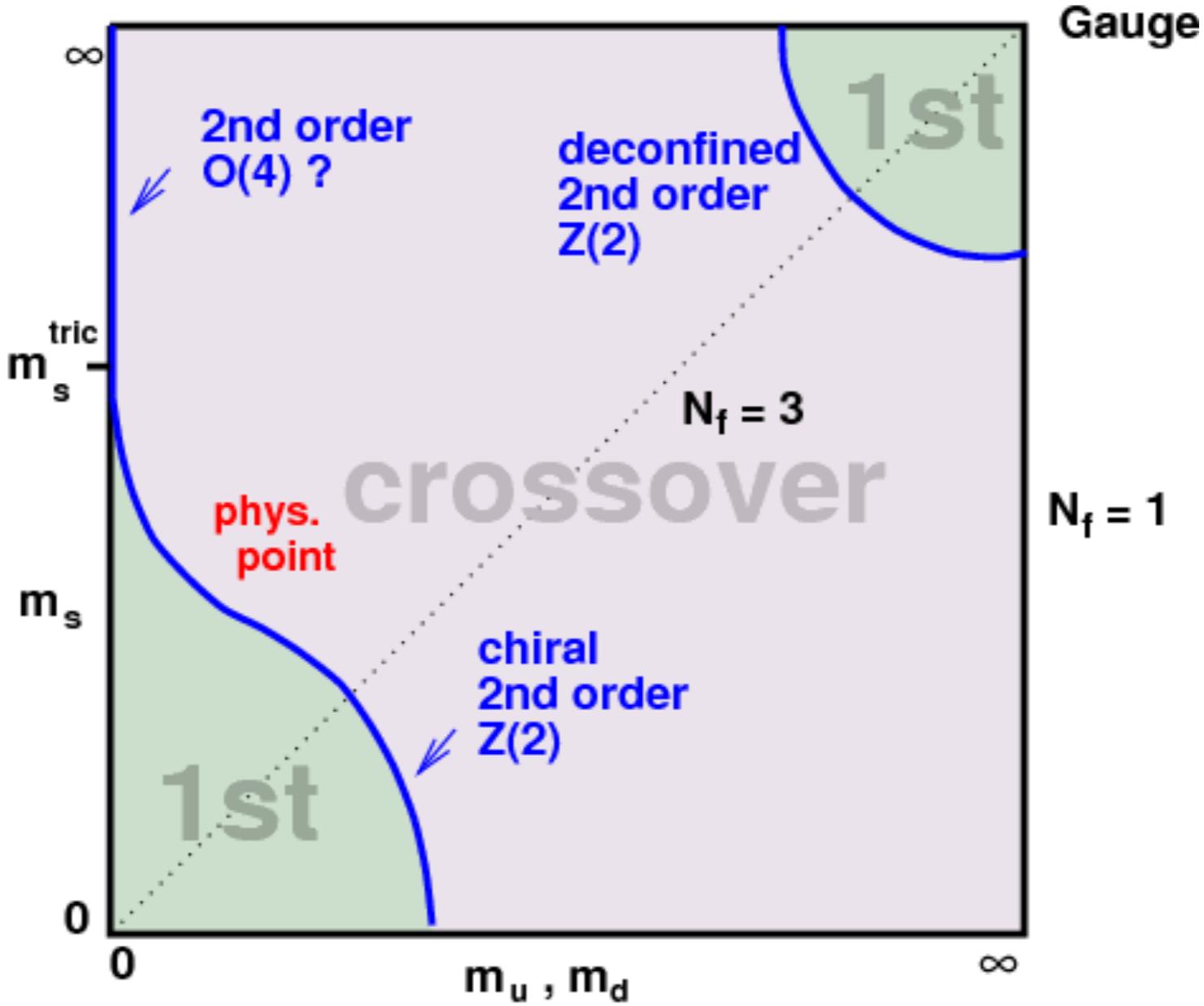
GRAVITATIONAL WAVES SPECTRUM

- ◆ Determined by 3 parameters:
 - ◆ α → relative energy density in the source (related to latent heat at the phase transition)
 - ◆ β → bubble nucleation rate proportional to inverse time of the transition (related to tunneling probability between vacua)
 - ◆ u → bubble velocity
- ◆ Plus we need to know the temperature of the phase transition $T_{\star} \simeq T_c$



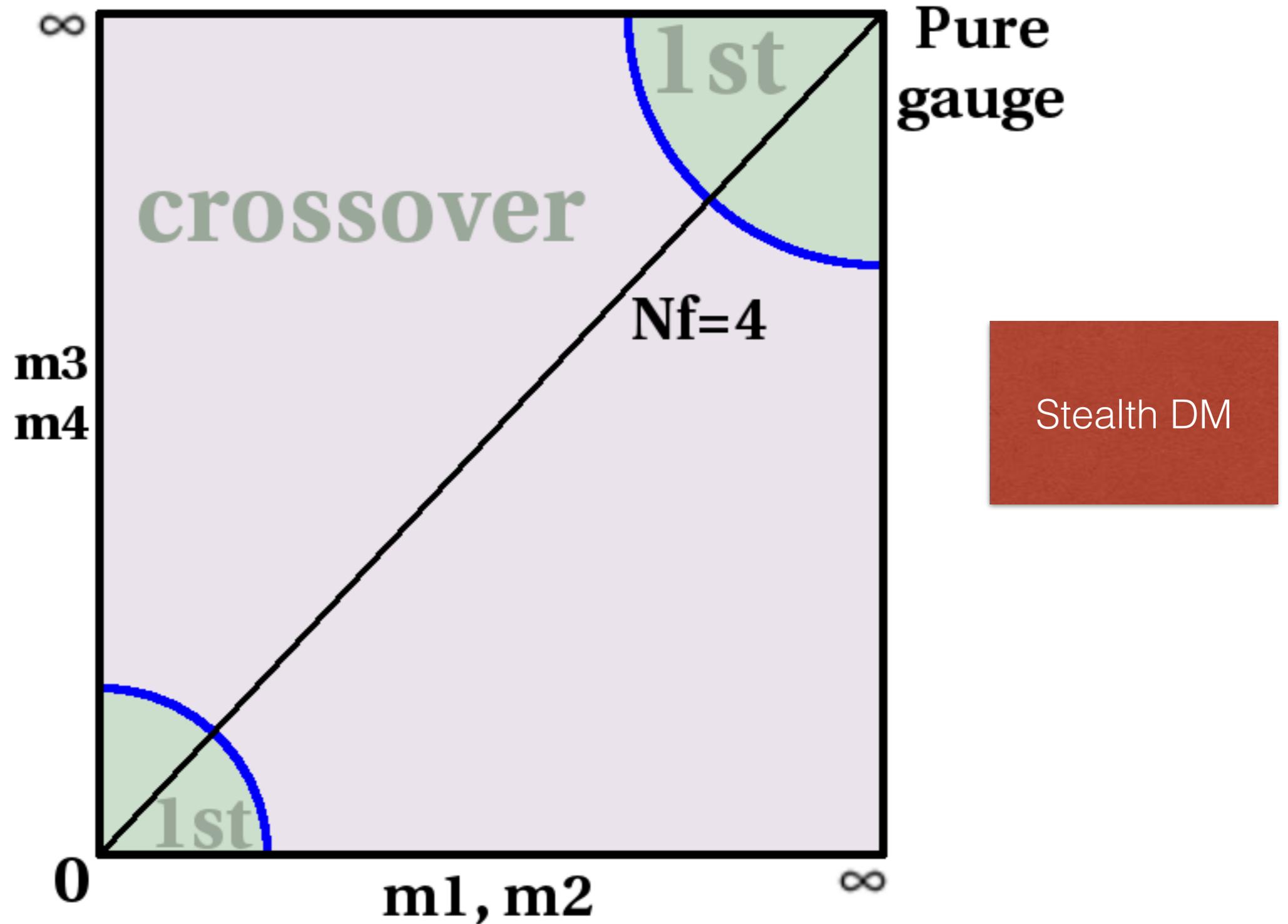
Phase Transitions in Strongly-coupled Theories

QCD

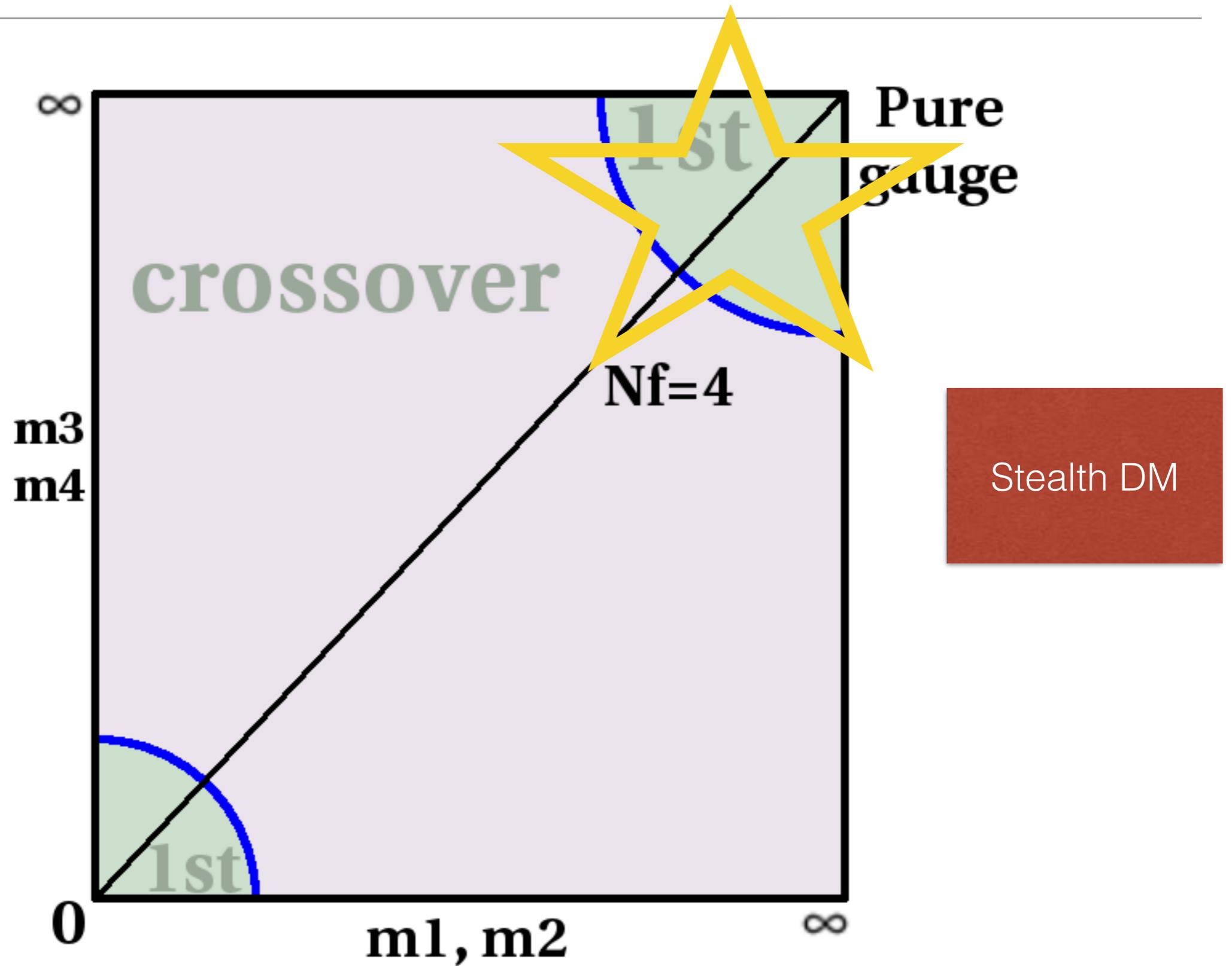


“Columbia” plot

Phase Transitions in Strongly-coupled Theories



Phase Transitions in Strongly-coupled Theories



BUILDING BLOCKS FOR THE PHASE DIAGRAM

$SU(4)$

with dynamical nHYP smeared
staggered fermions

$a \cdot m$

$\{0.05, 0.1, 0.2, 0.4, \infty\}$

N_t

$\{4, 6, 8, 12\}$

$\alpha = L/N_t$

$\{2, 3, 4, 6, 8\}$

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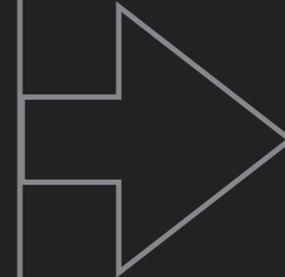
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Scan in

β_F

~ 1370 ensembles

LATTICE OBSERVABLES

$$|PL_W|$$

and $\chi_O = L^3 (\langle O^2 \rangle - \langle O \rangle^2)$

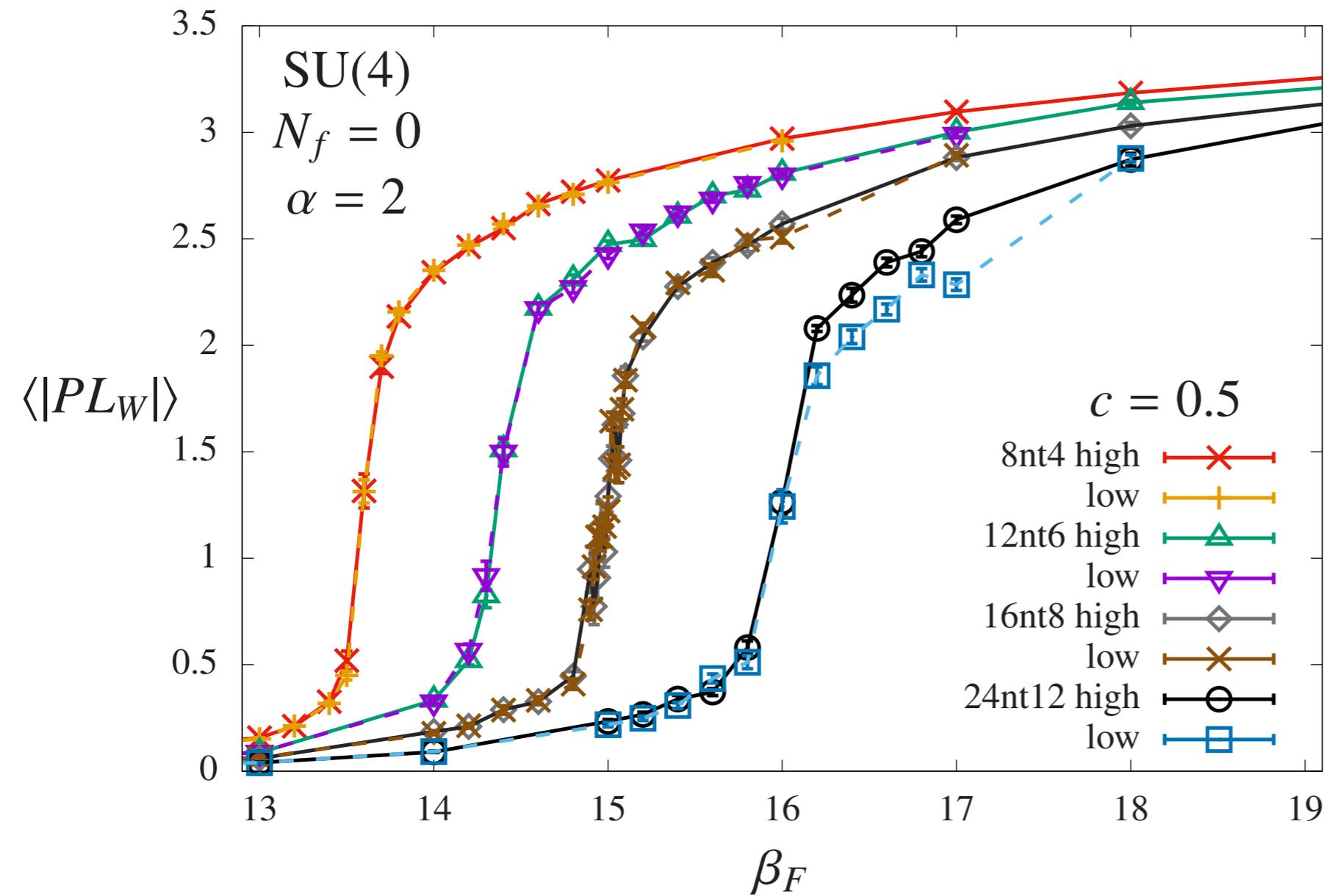
$$R_E(t)$$

$$\equiv \left\langle \frac{E_{ss}(t)}{E_{s\tau}(t)} \right\rangle \begin{matrix} \nearrow \approx 1 \\ \searrow > 1 \end{matrix}$$

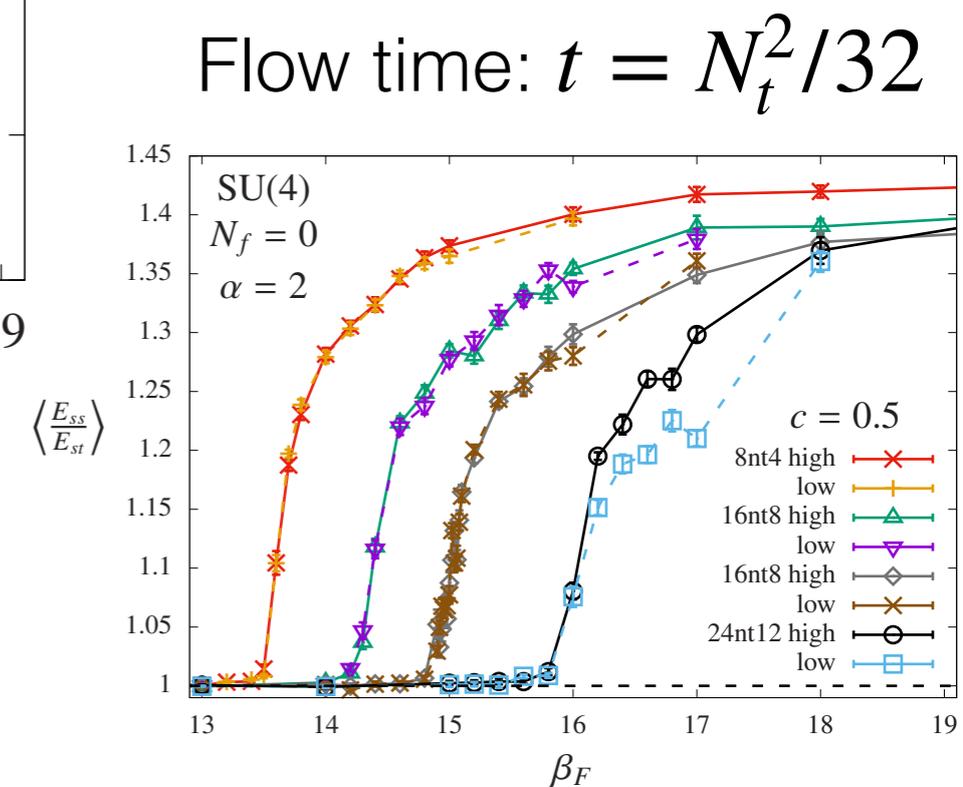
$$f(\theta)$$

$$\equiv \frac{\pi/4}{\pi/4 - \theta} \left[\frac{N_{\mathbf{in}}}{N_{\mathbf{tot}}} - \frac{\theta}{\pi/4} \right] \begin{matrix} \nearrow \rightarrow 0 \\ \searrow \rightarrow 1 \end{matrix}$$

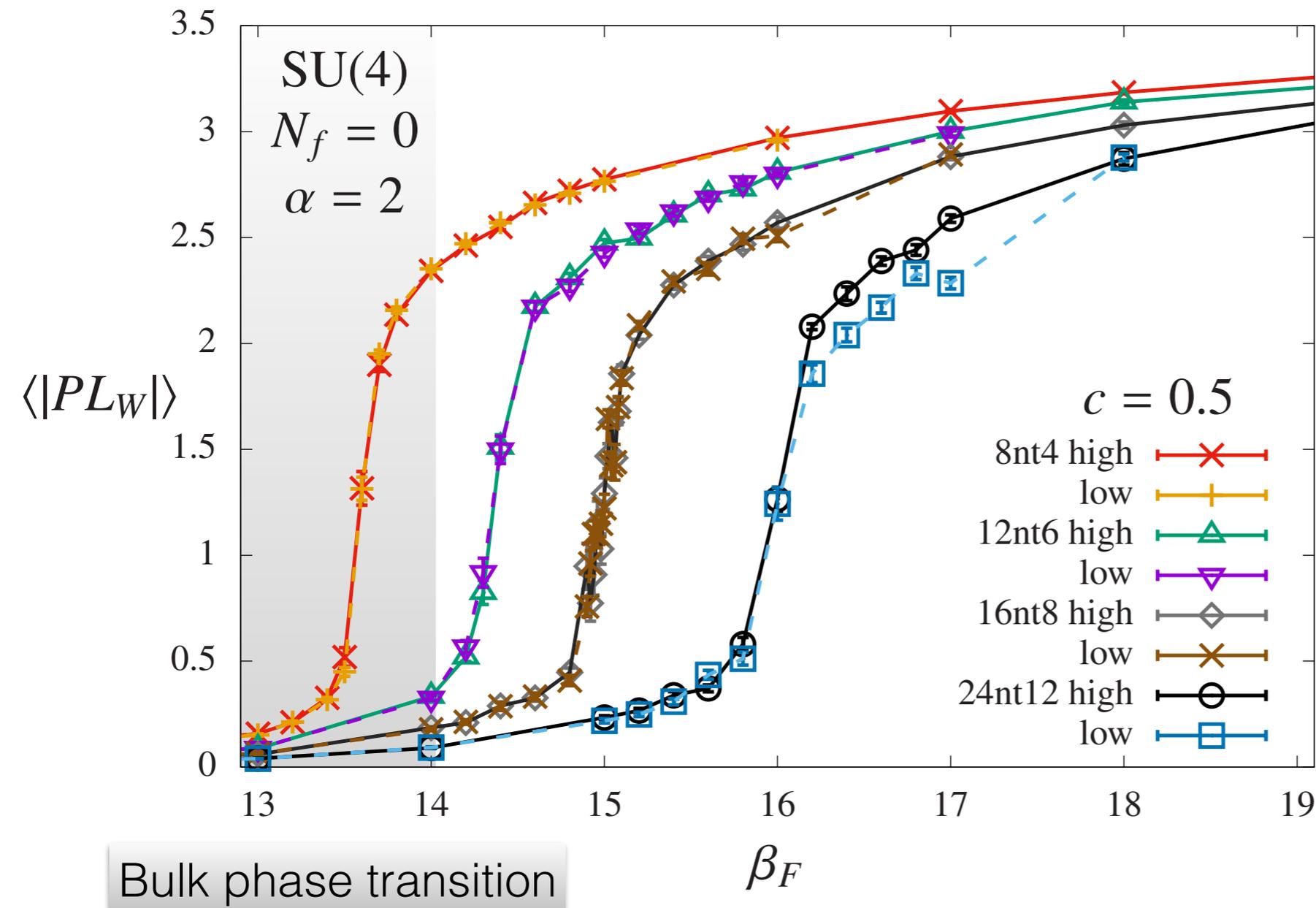
Results: Pure-Gauge system $a \cdot m = \infty$



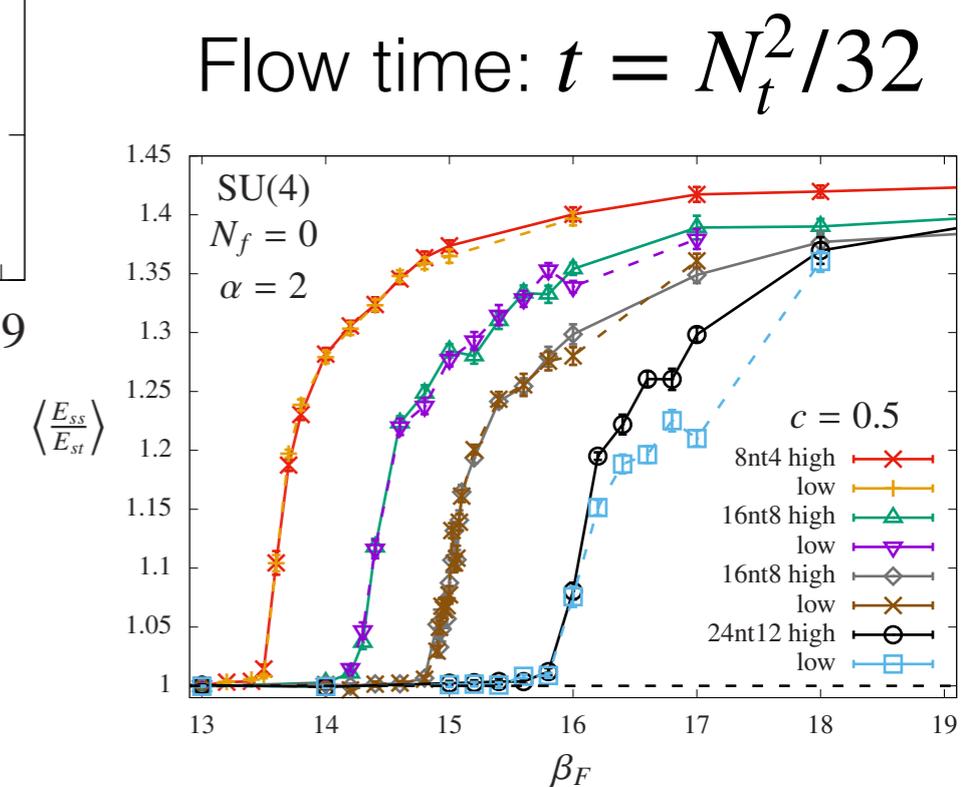
- Long autocorrelations
- Topological freezing
- high-low starts



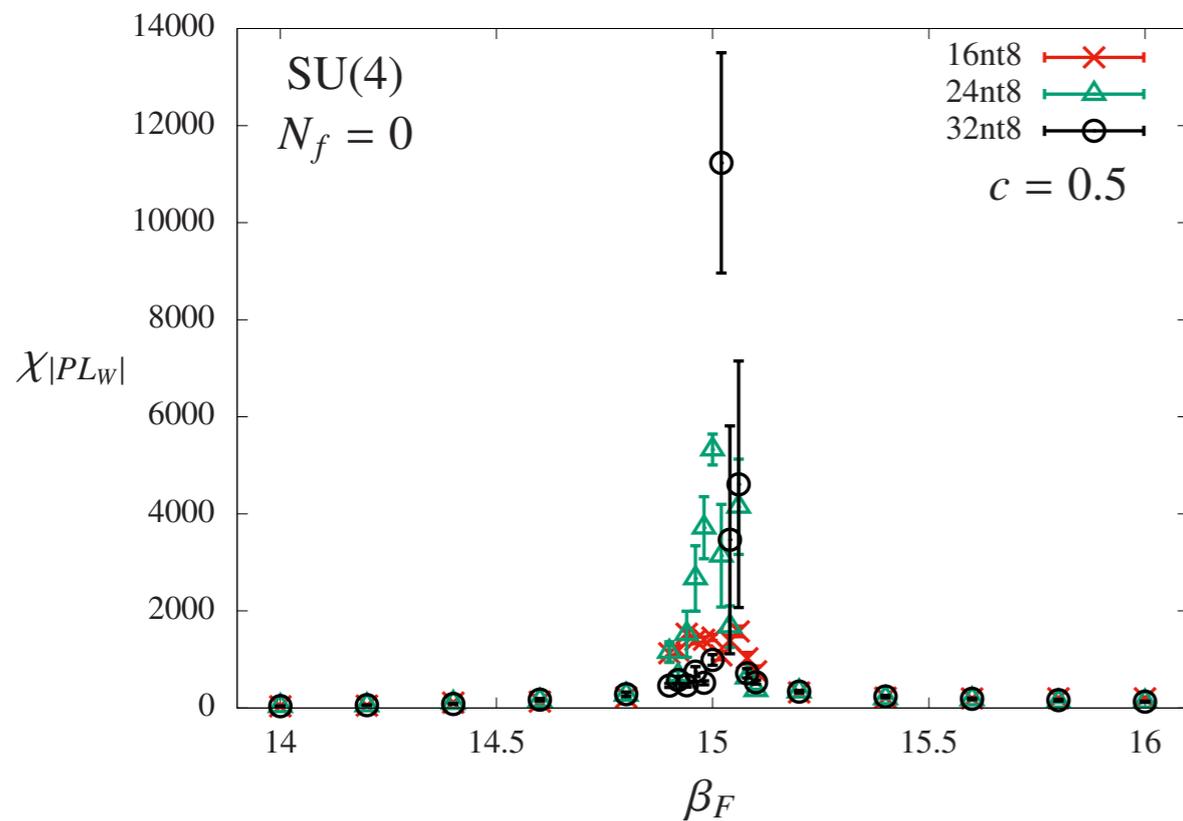
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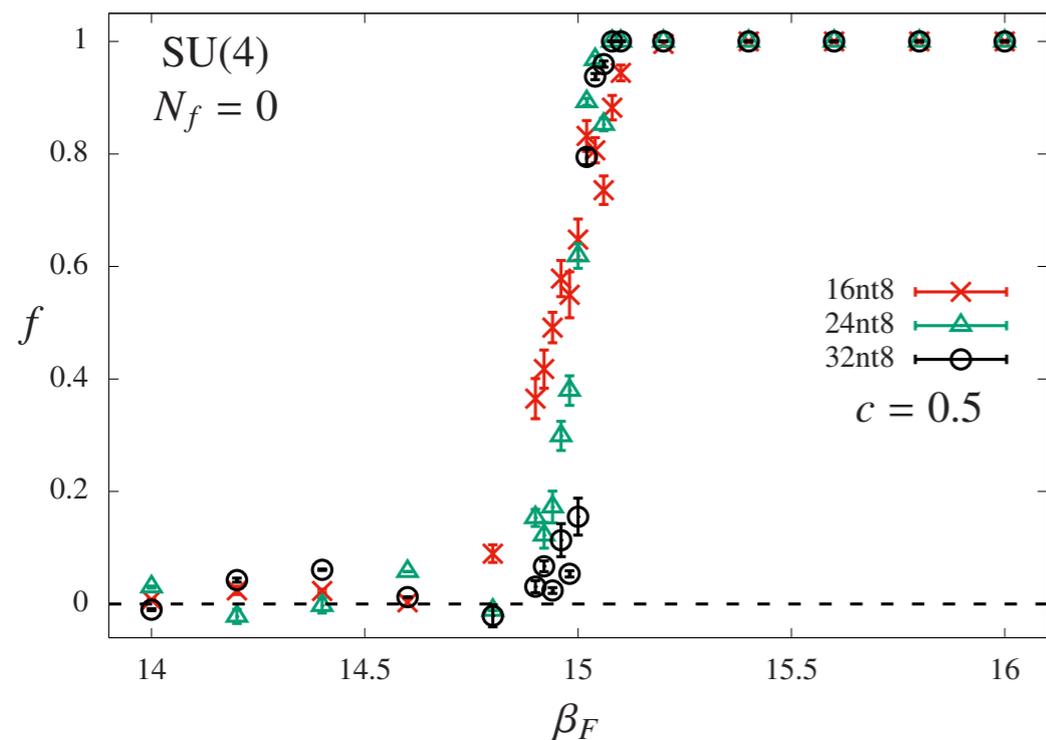
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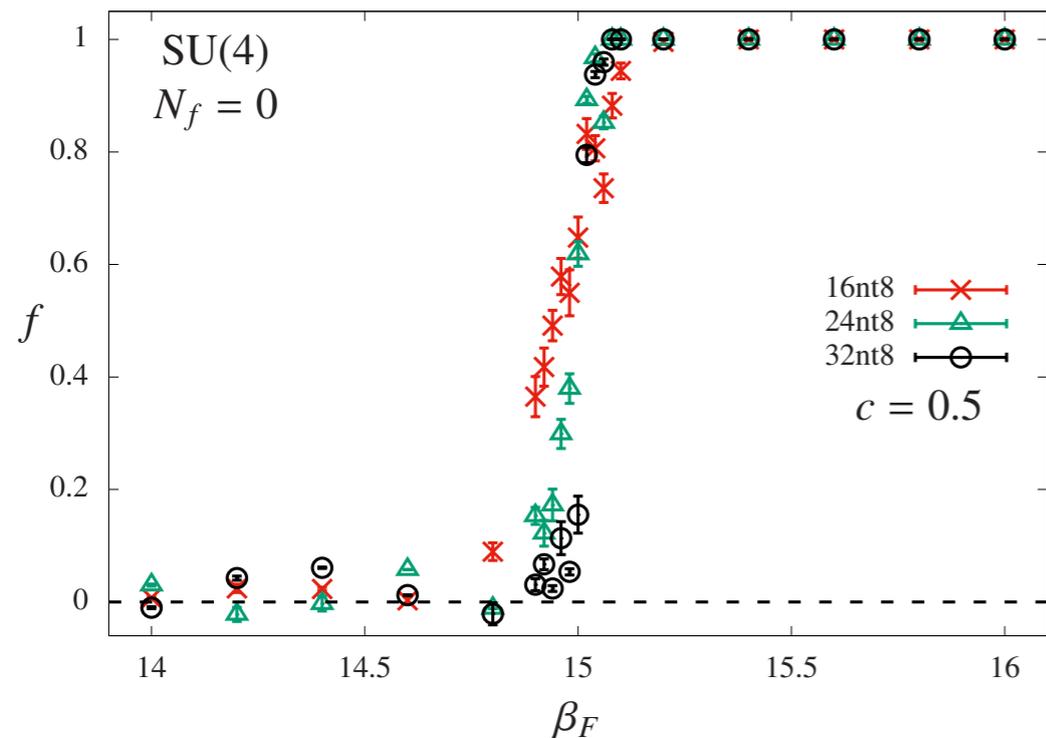
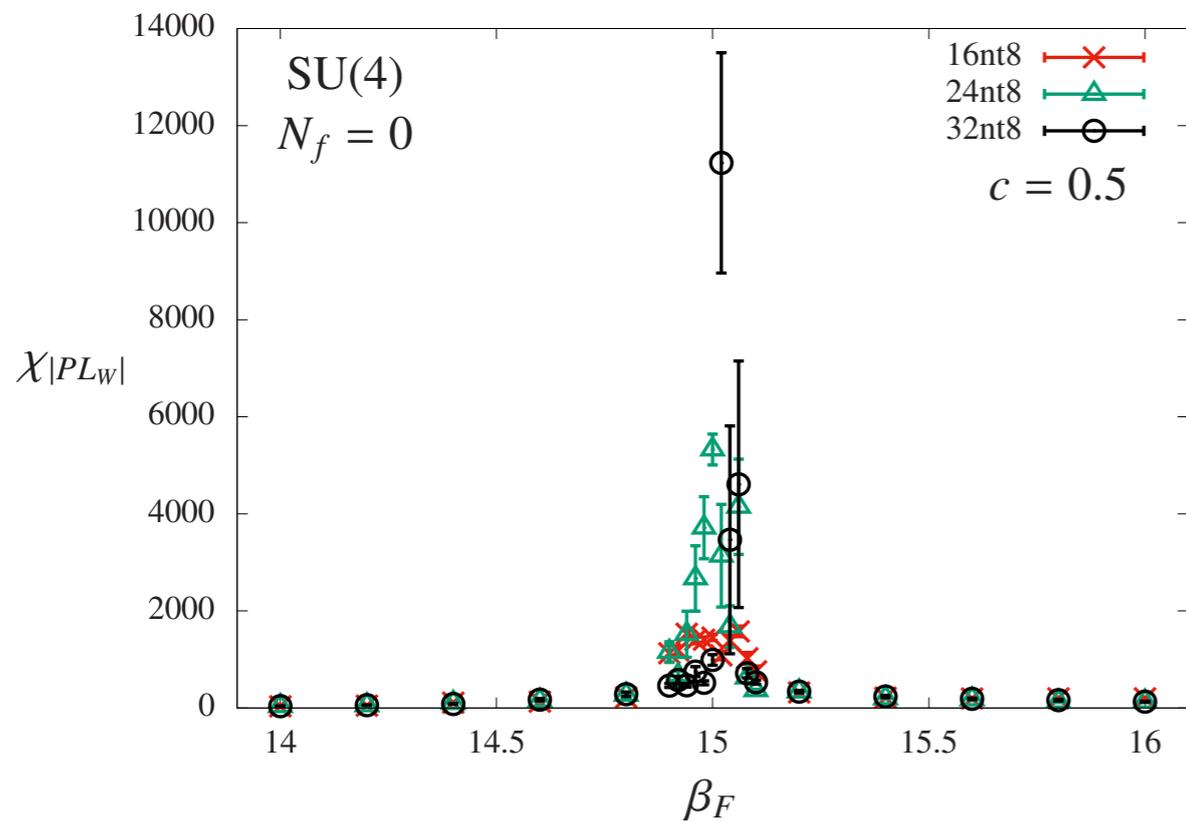
Results: Pure-Gauge system $N_t = 8$



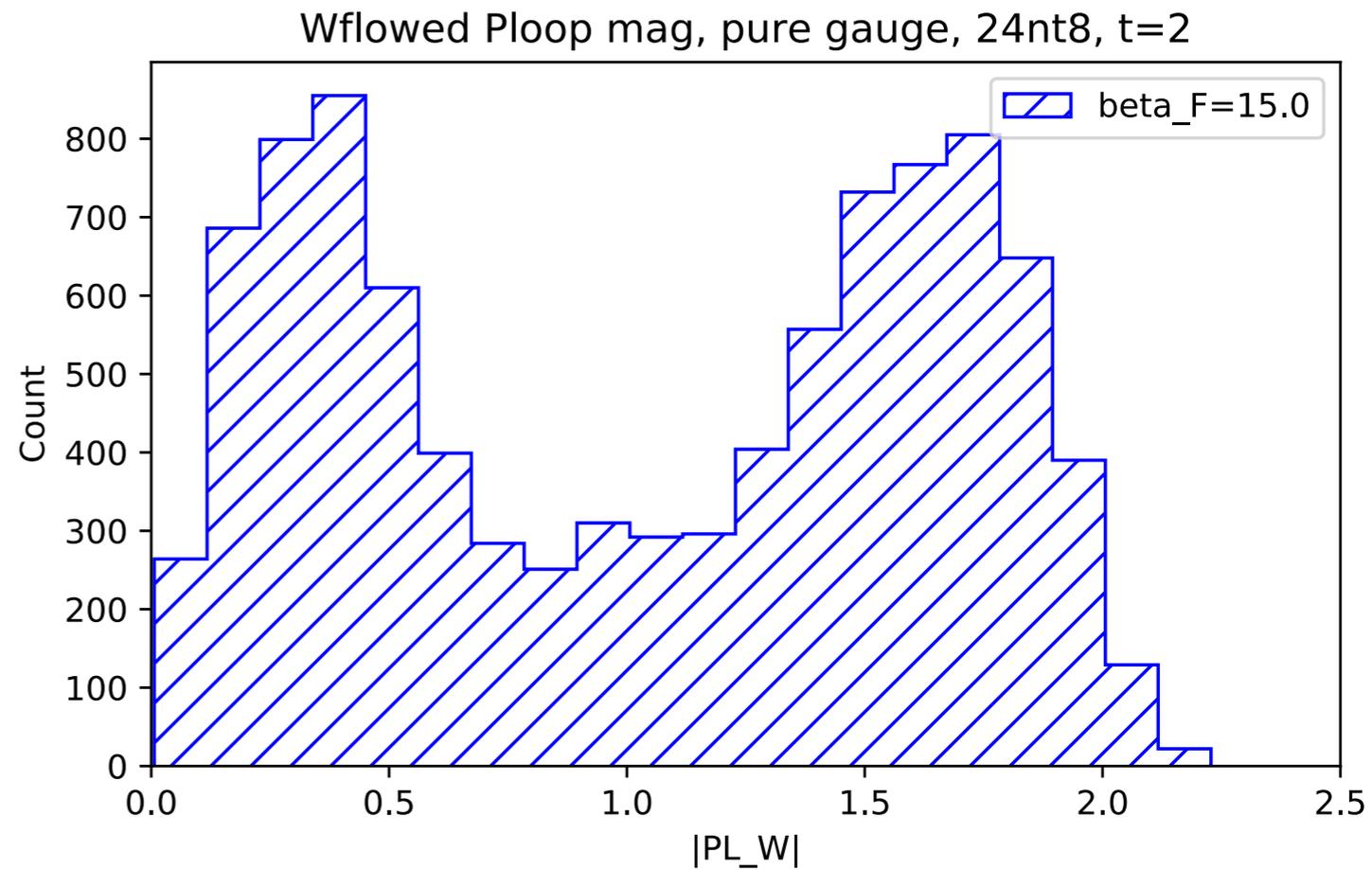
- Peak of χ grows with volume
- Deconfinement fraction gets steeper



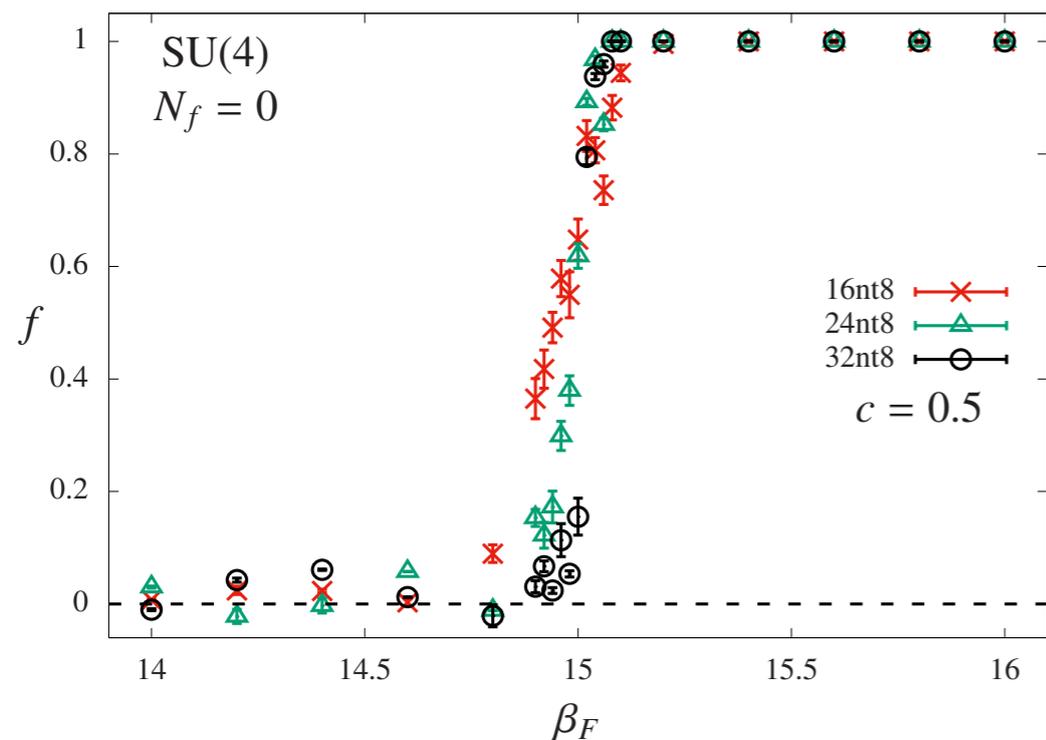
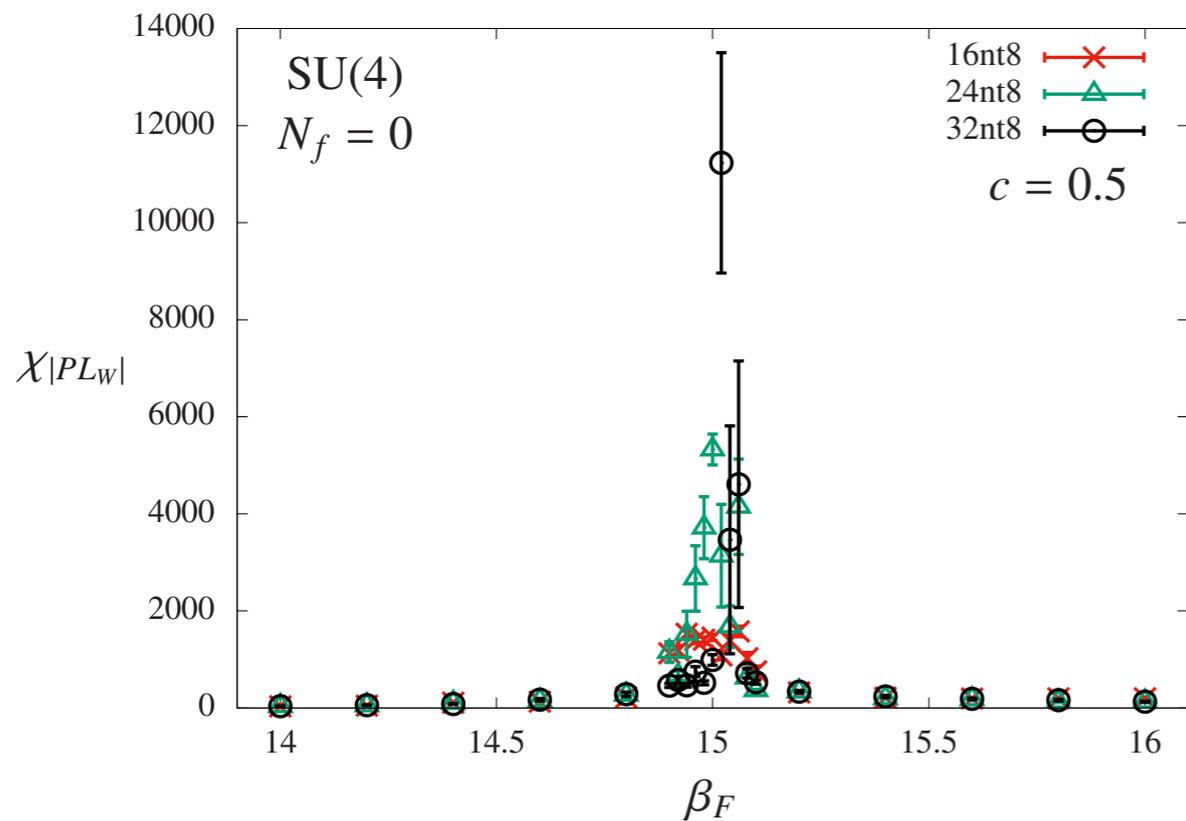
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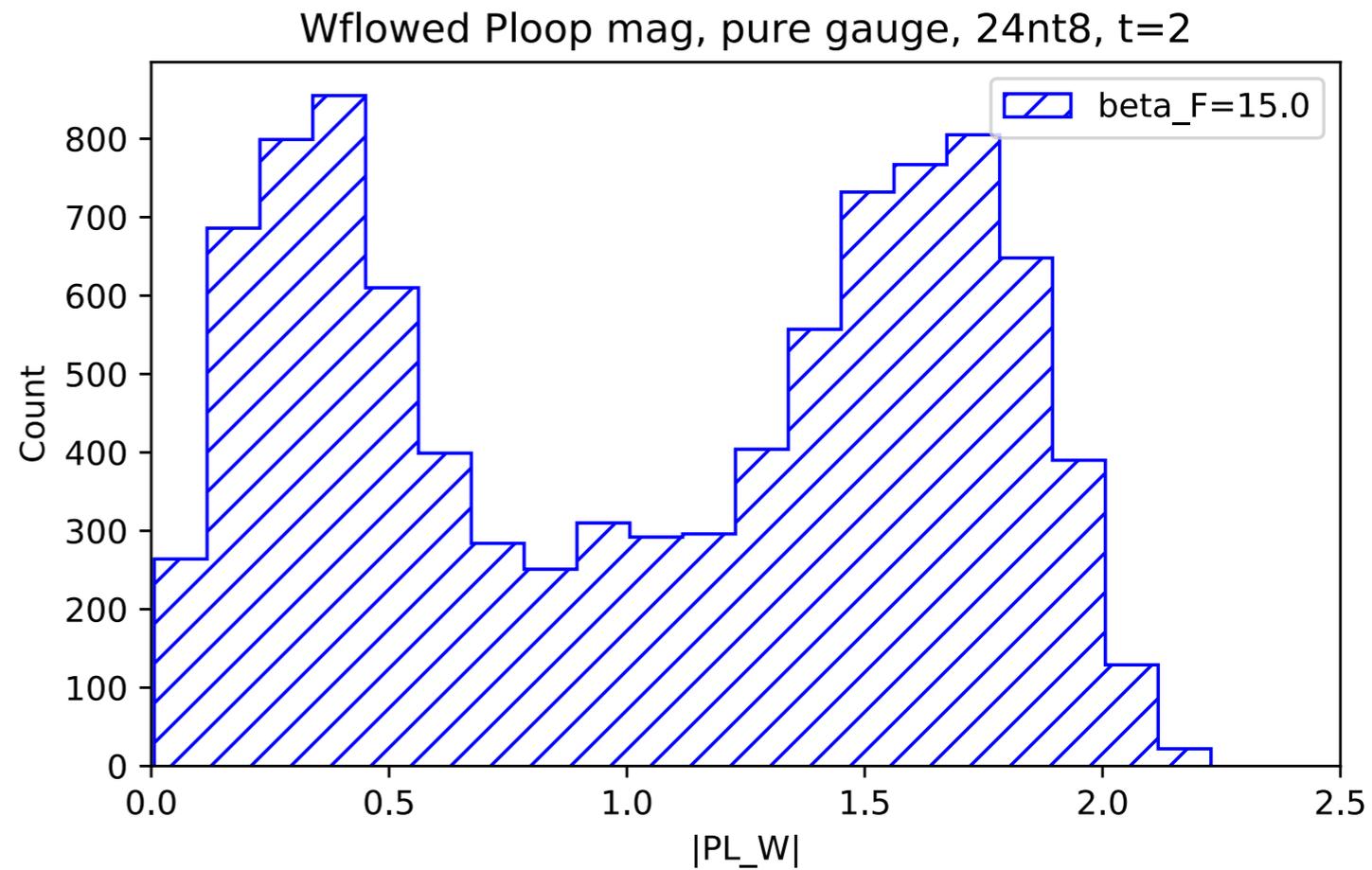
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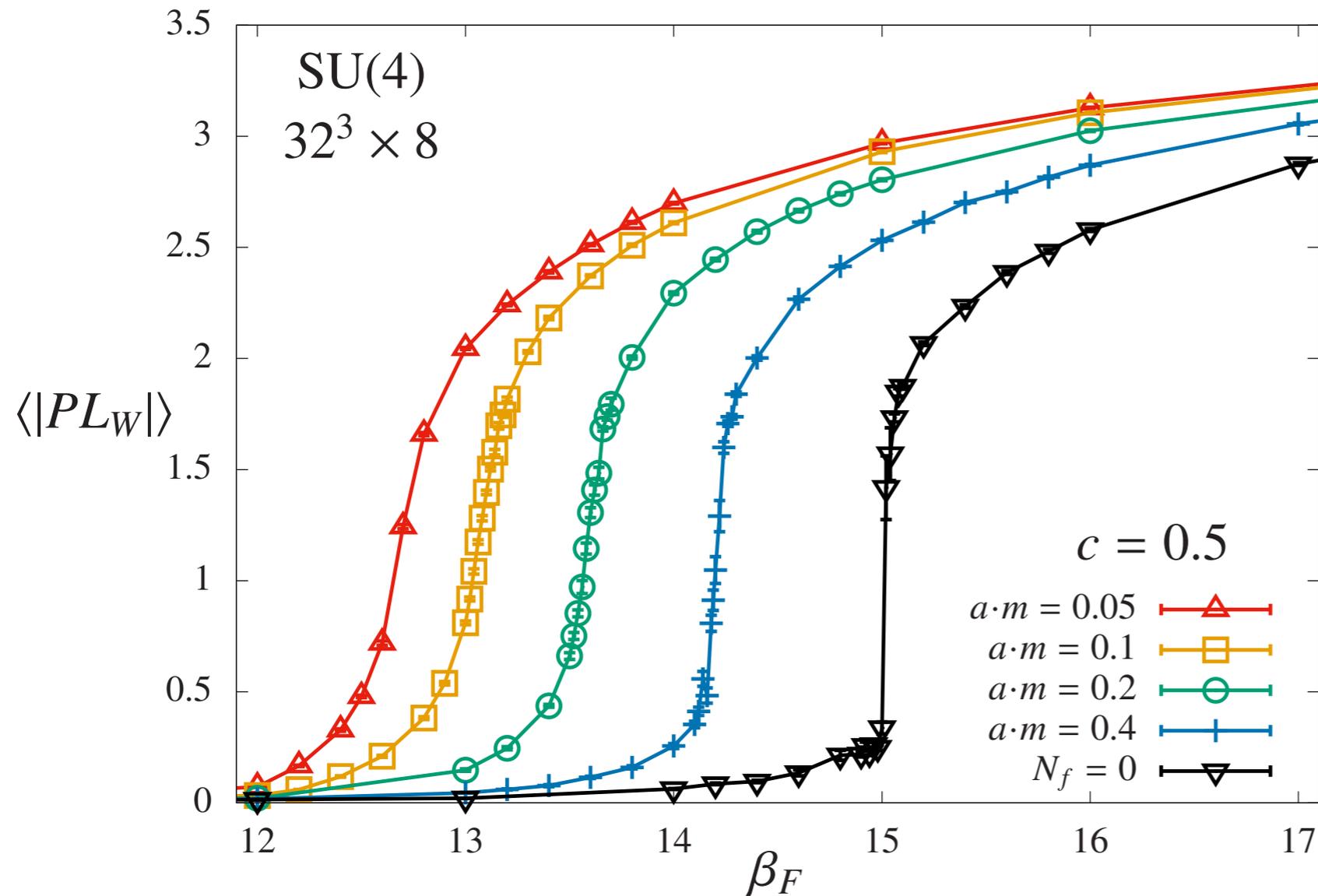


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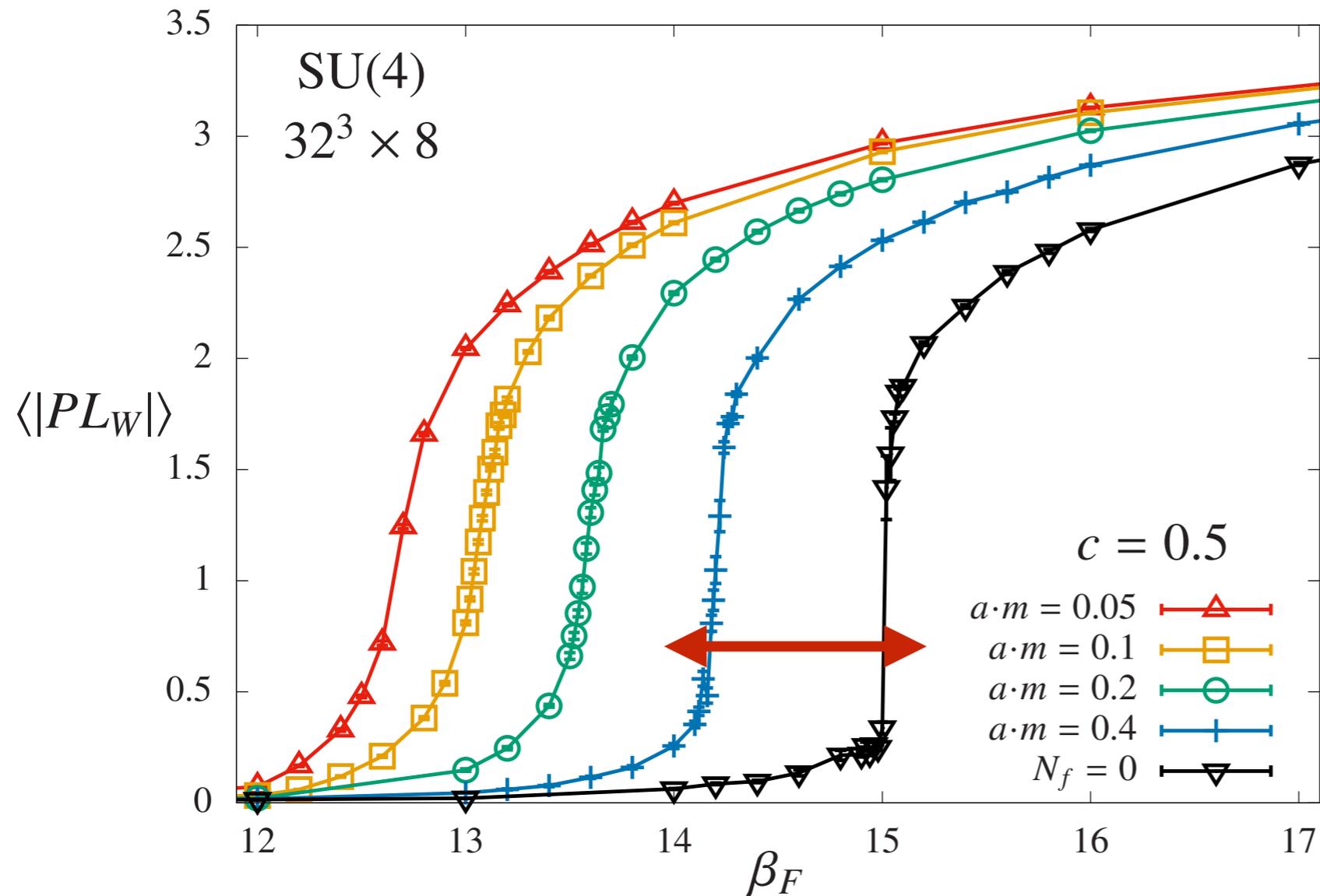
Signs of 1st order transition

Results: Dynamical Fermions $N_f = 4$



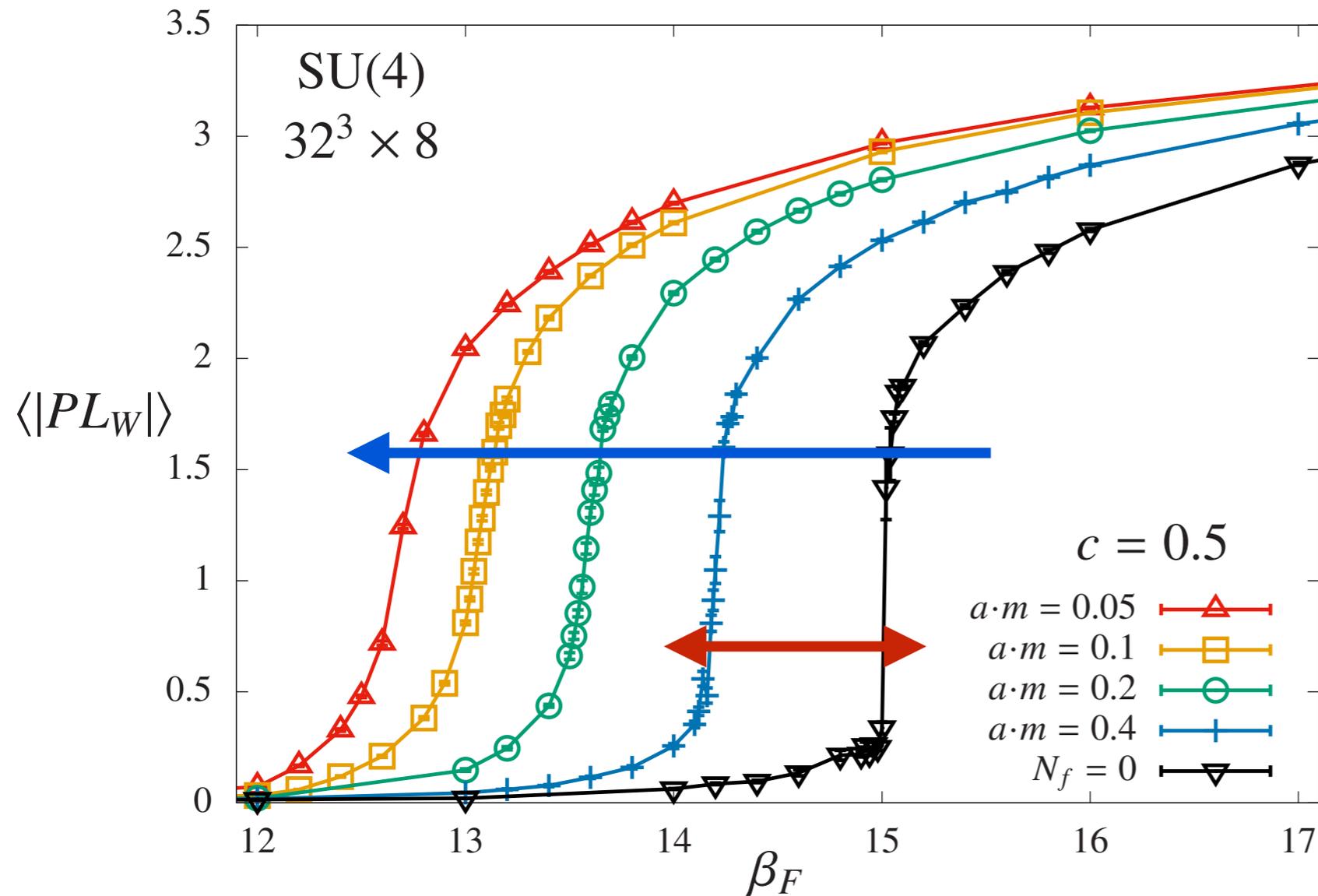
- Difference from quenched
- Stronger couplings needed at smaller mass
- Susceptibility scales for $a \cdot m > 0.2$
- Two-peak histogram
- Still no continuum limit

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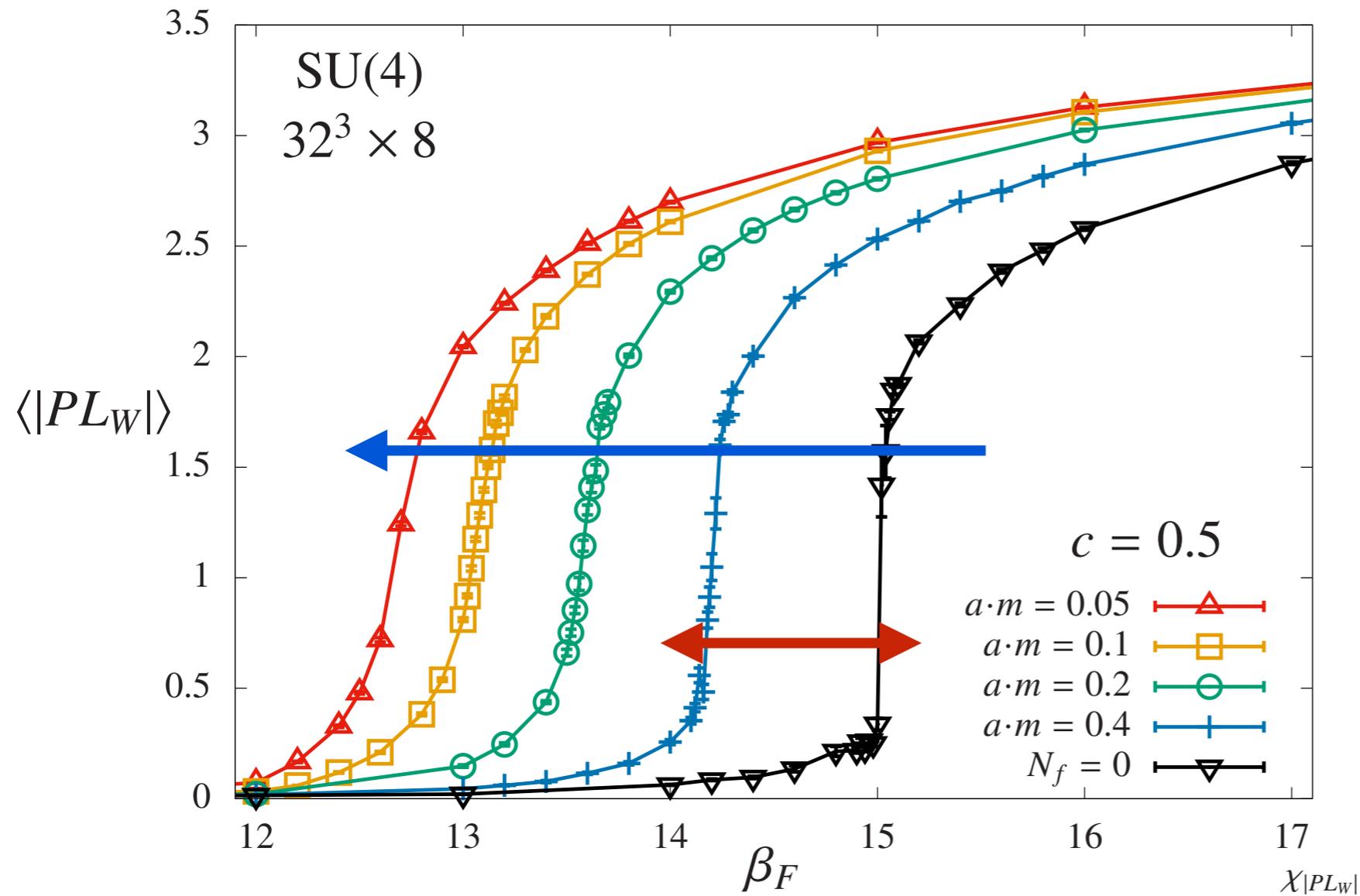
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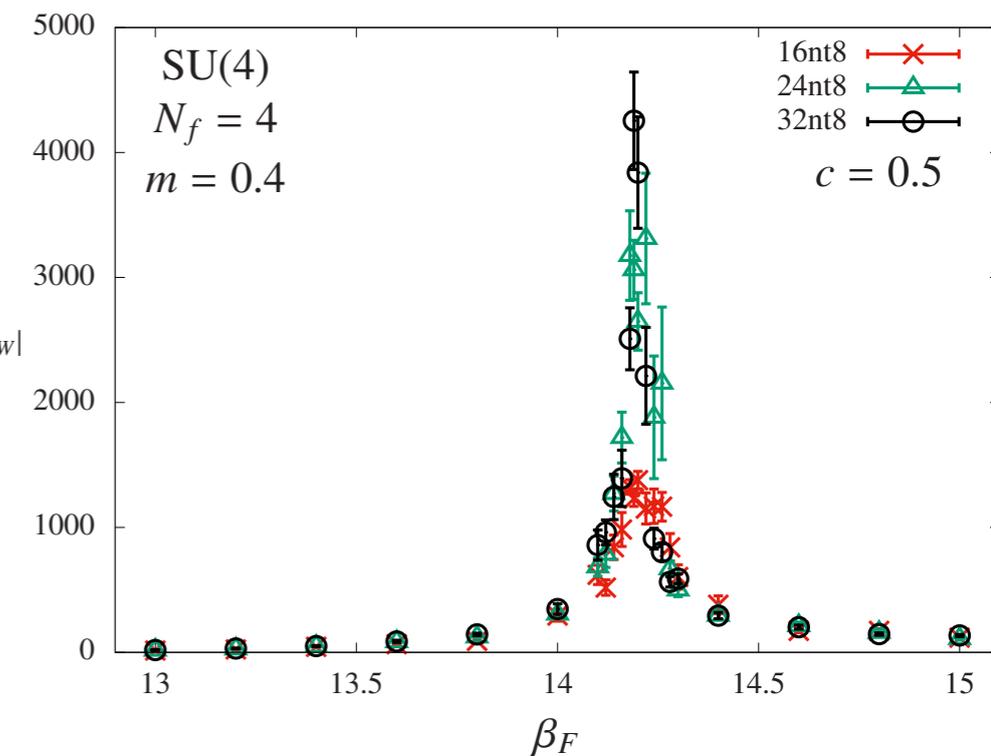


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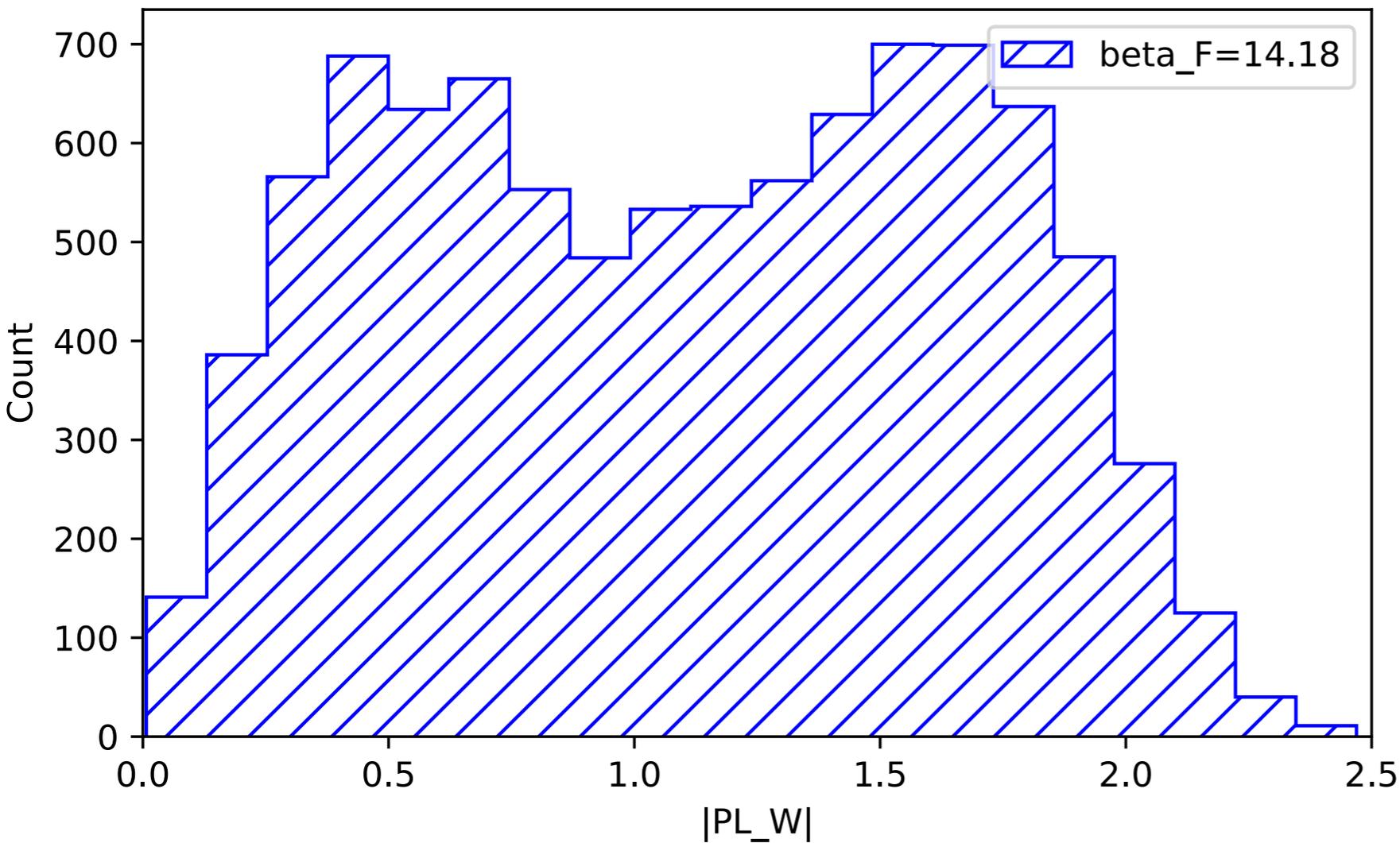


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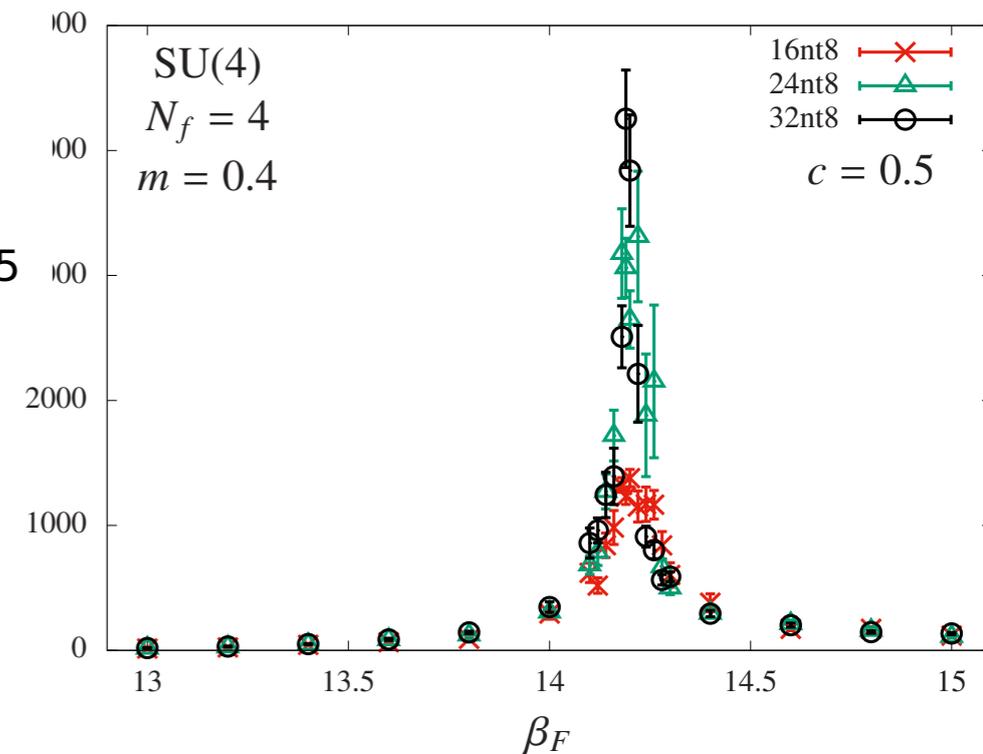
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Wflowed Ploop mag, $N_f=4$, 16nt8, $m=0.4$, $t=2$

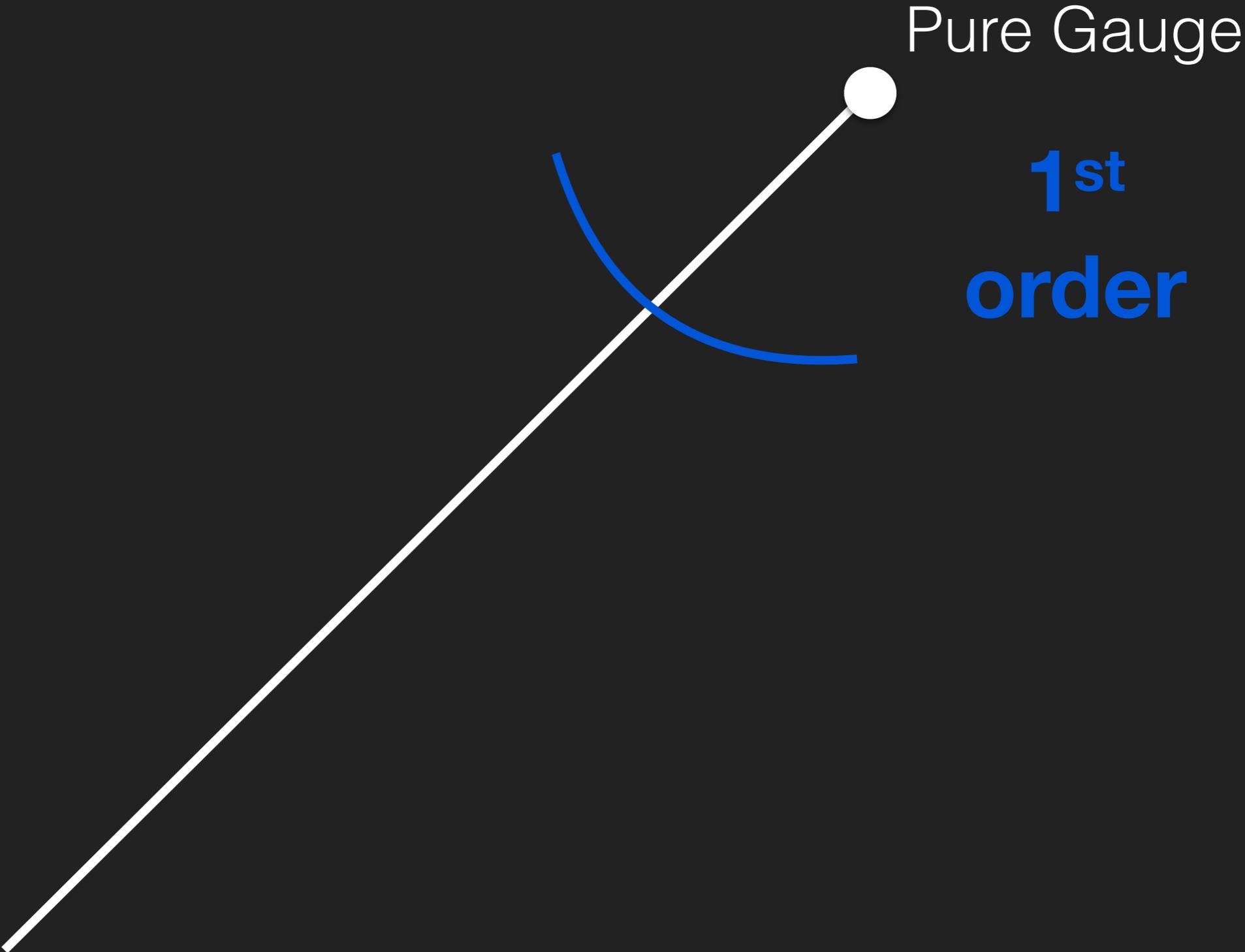


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SPECTRUM

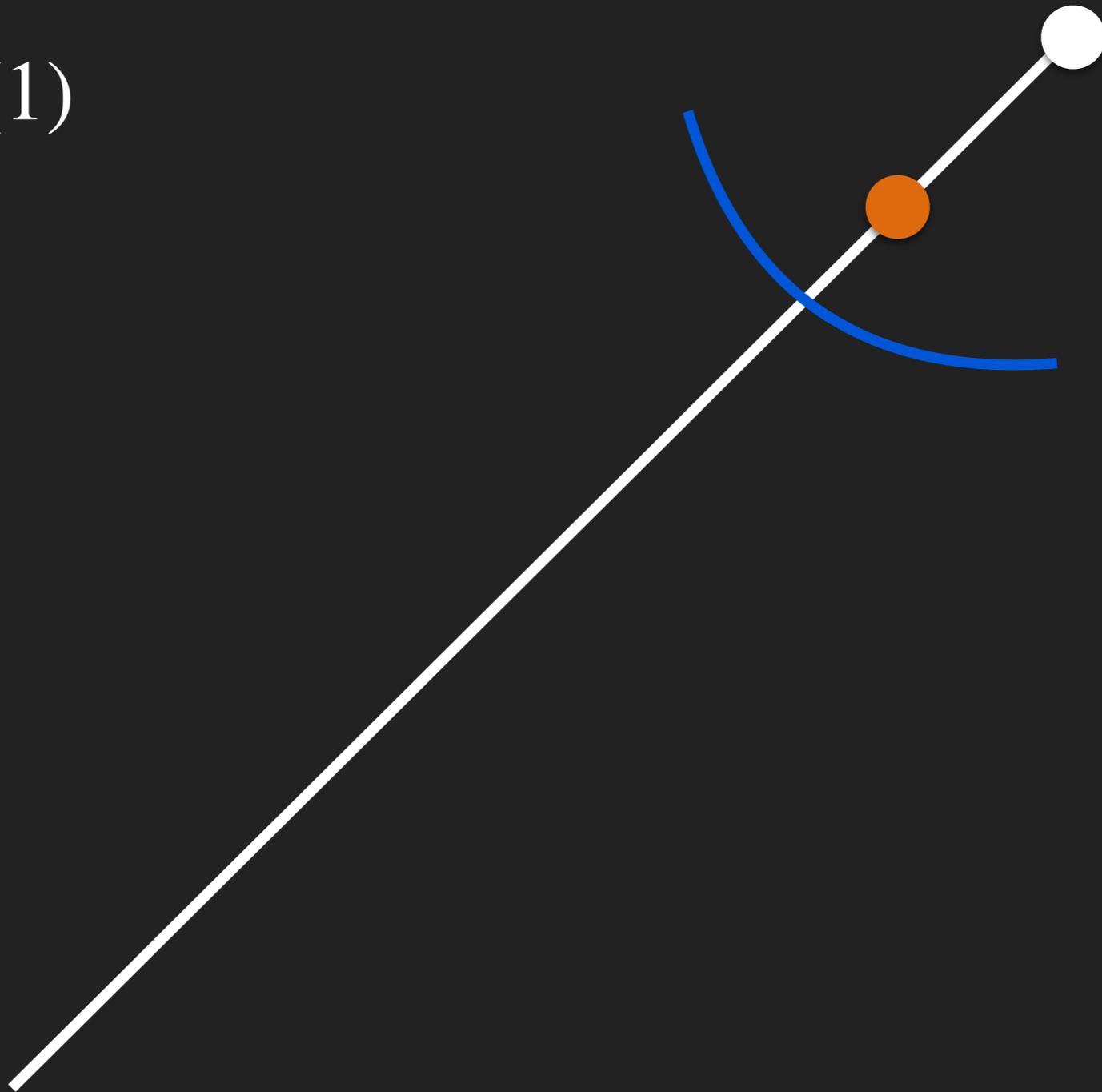


SPECTRUM

● $\frac{M_P}{M_V} = 0.96(1)$

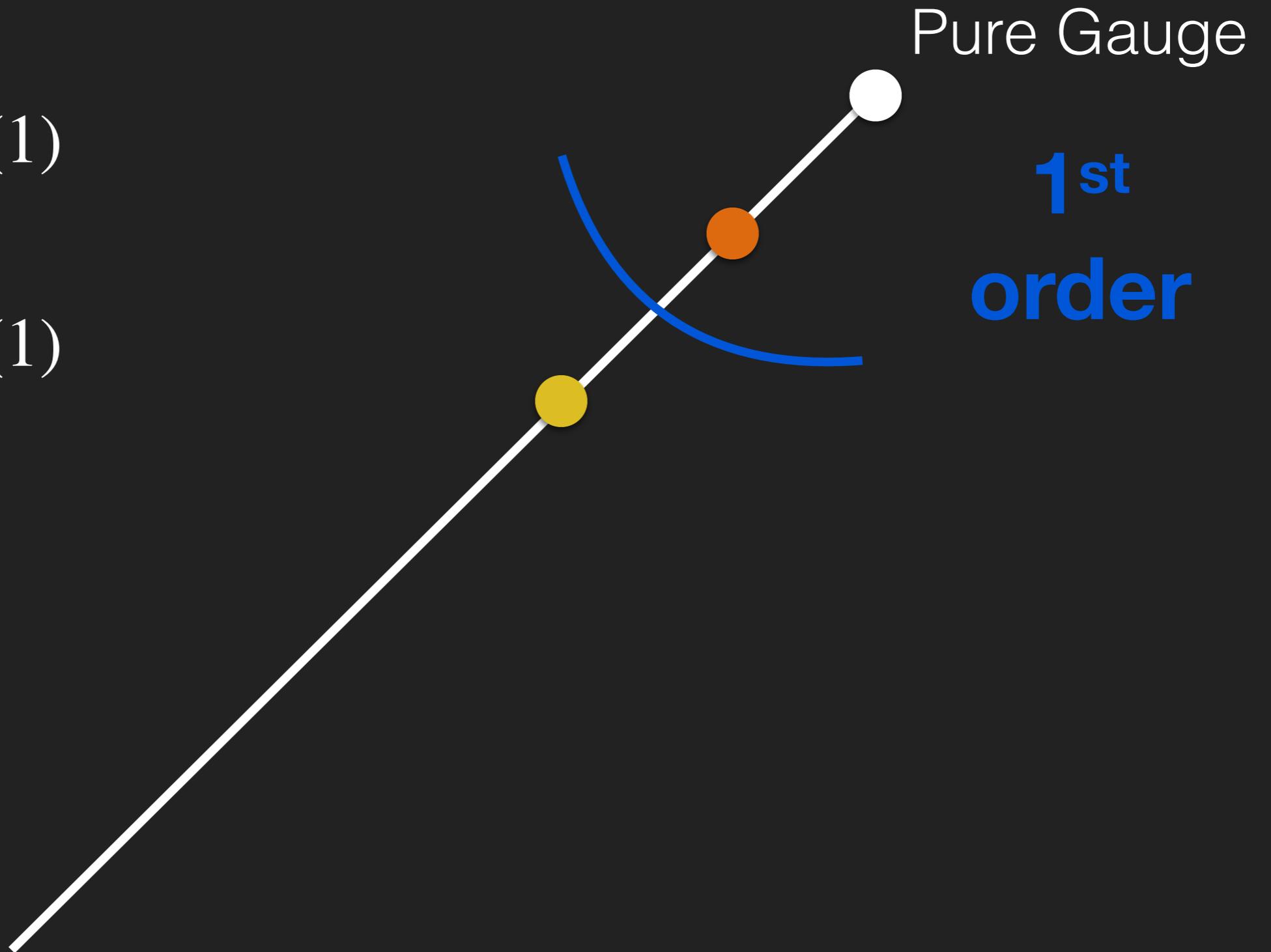
Pure Gauge

**1st
order**



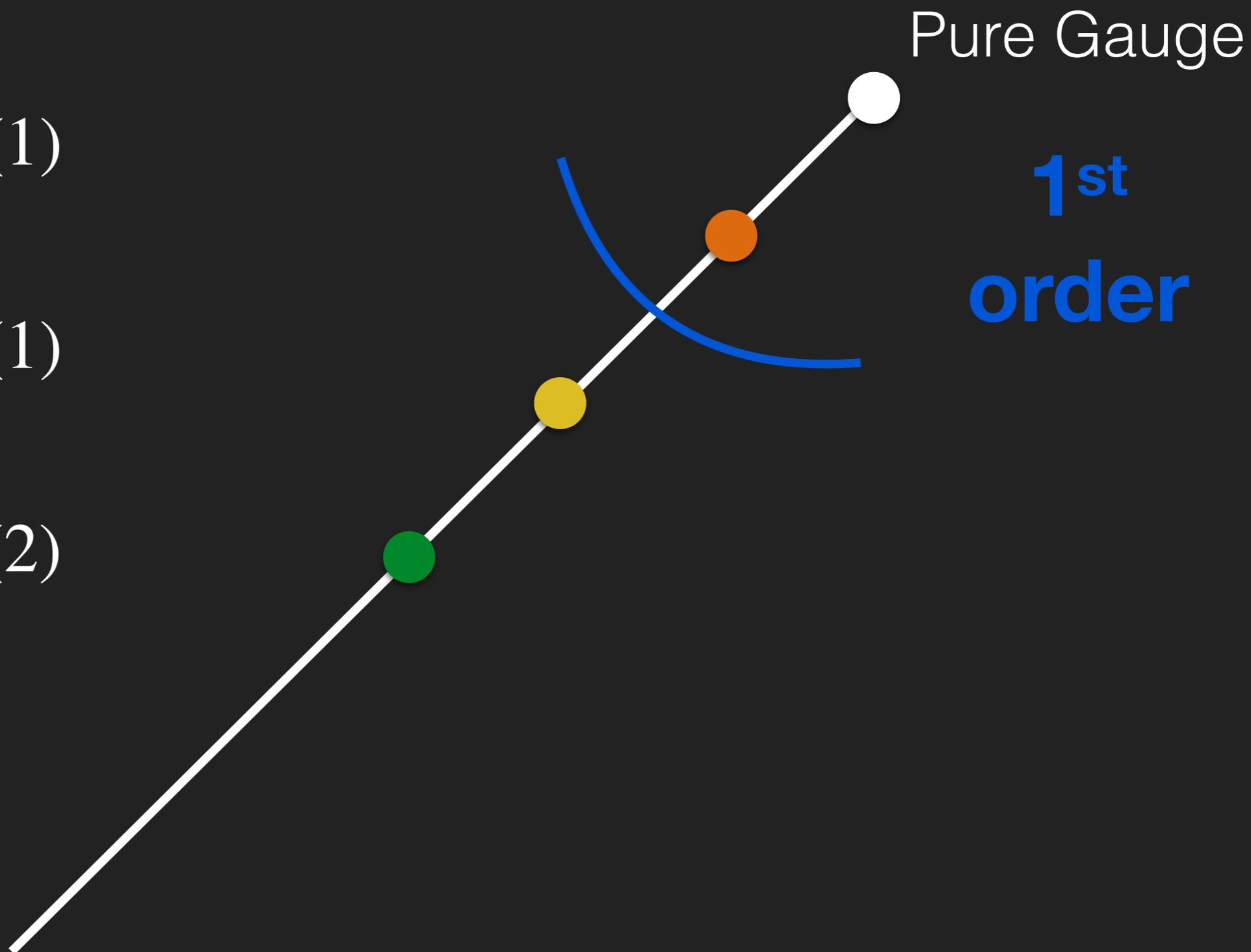
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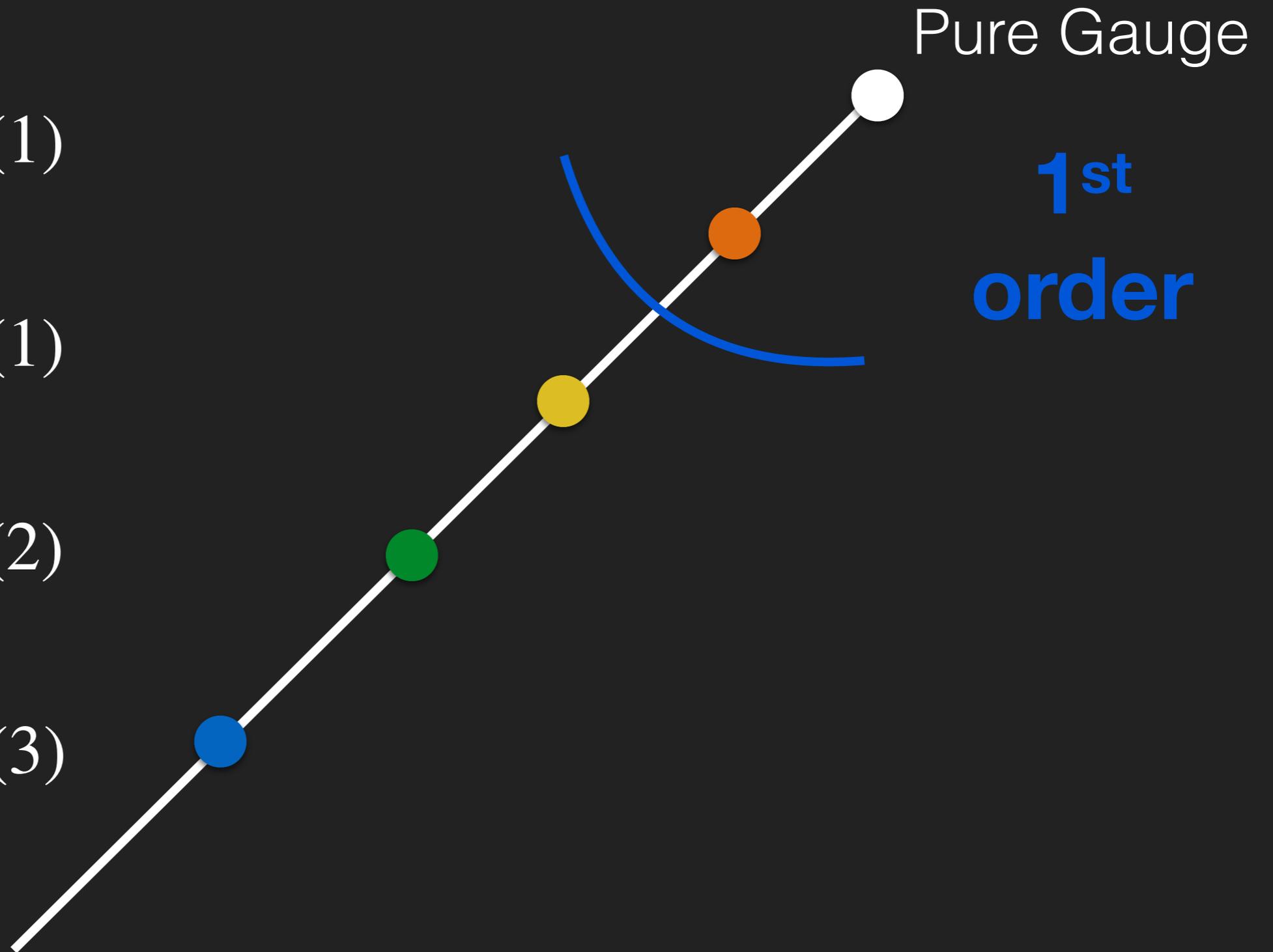
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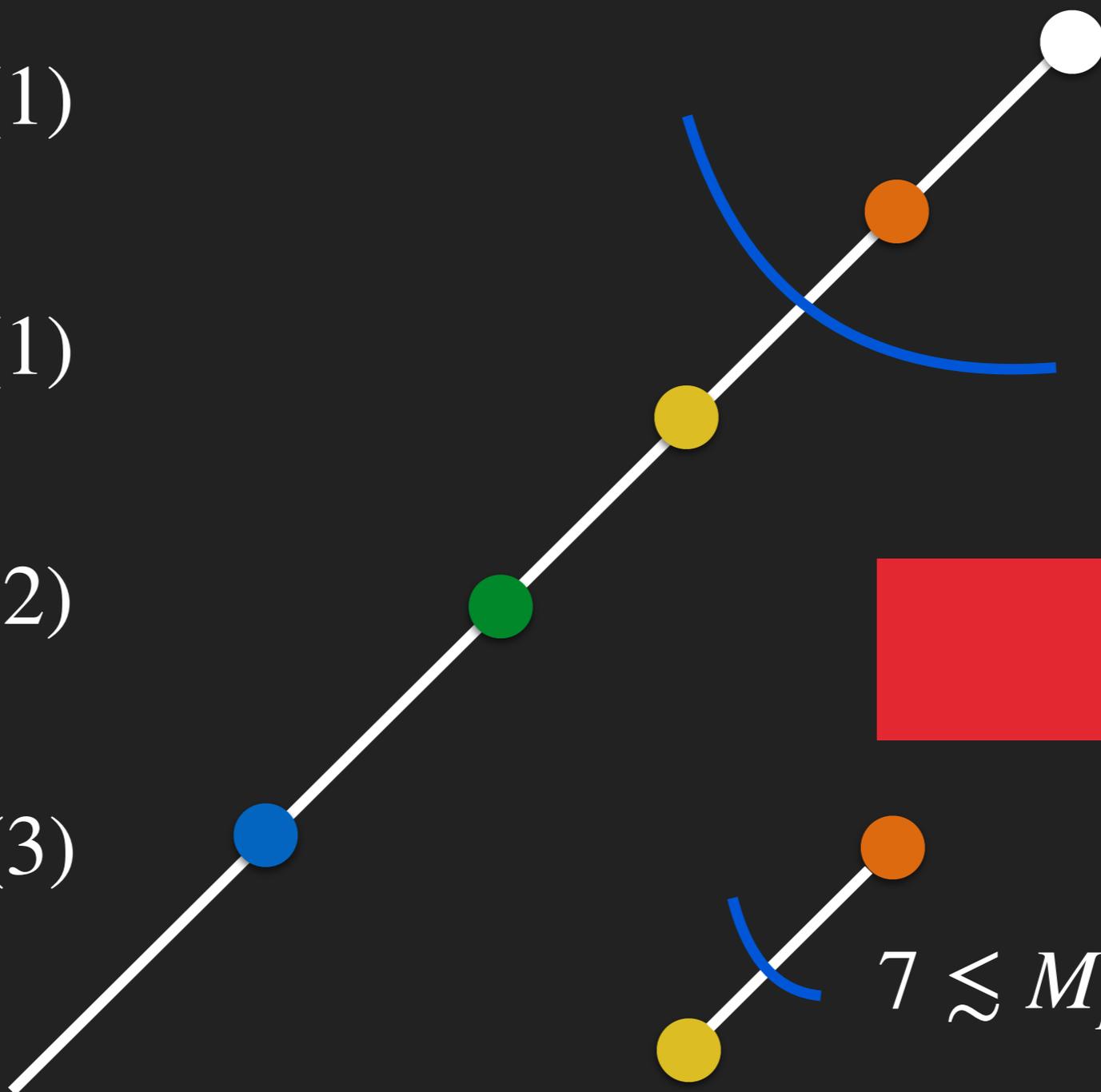
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**1st
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END-POINT

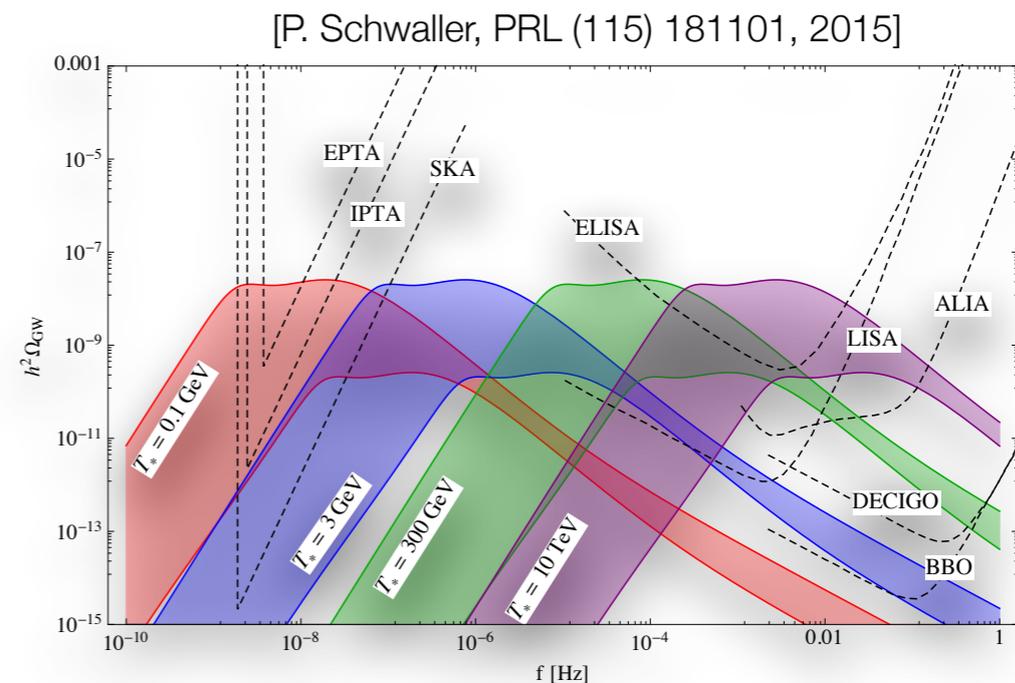
$7 \lesssim M_P/T_c \lesssim 10$



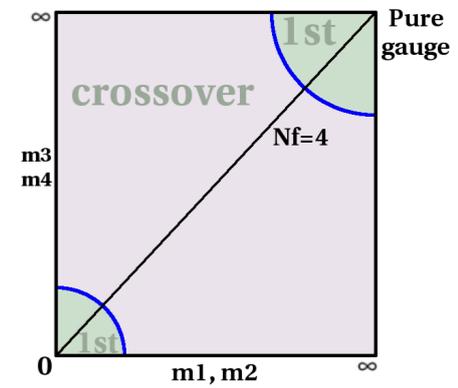
Conclusions

- ◆ Composite Dark Matter provides interesting signals for dark matter searches at colliders and in direct detection experiments
- ◆ With a 1st order confinement-deconfinement transition, the dark sector can be discovered and constrained using gravitational waves
- ◆ Stealth Dark Matter is a SU(4) dark sector model with 4 heavy fermions
- ◆ Our lattice exploration of the phase diagram shows a thermal phase transition of 1st order at sufficiently high masses
- ◆ Using current bounds from experimental searches at colliders and our spectrum results we can provide a lower bound for the critical temperature

$$T_c > 0.2 \text{ TeV}$$

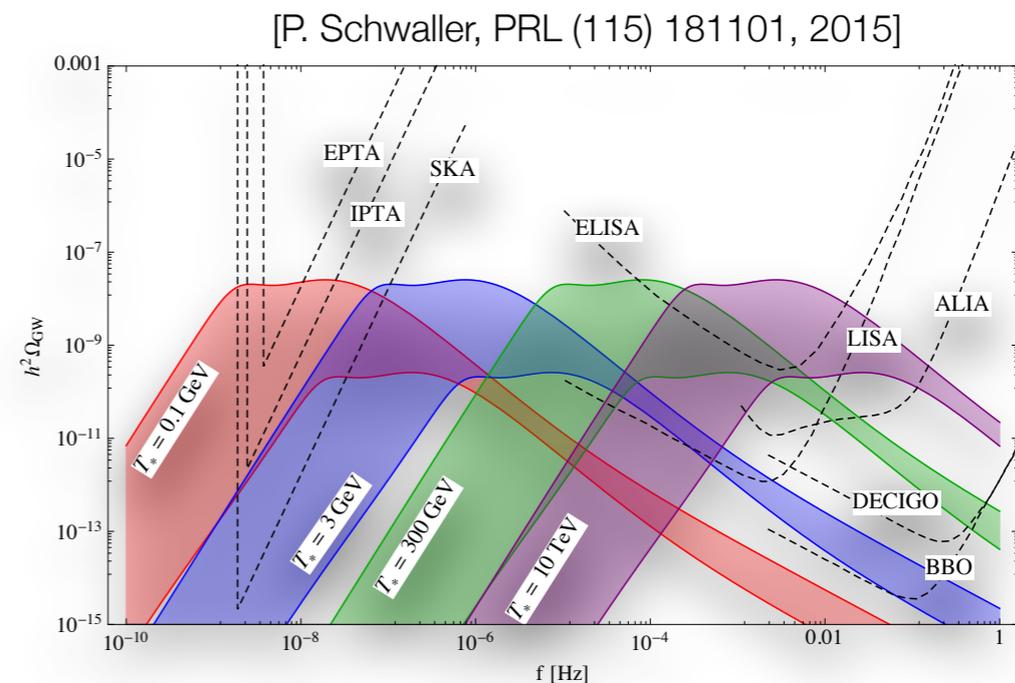


Conclusions

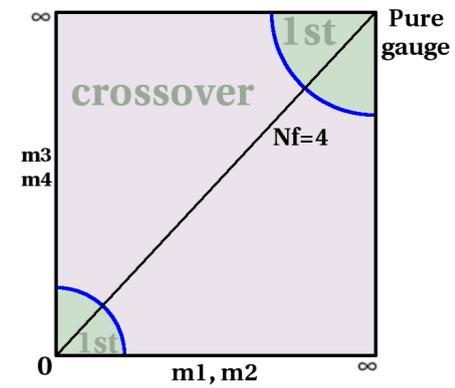


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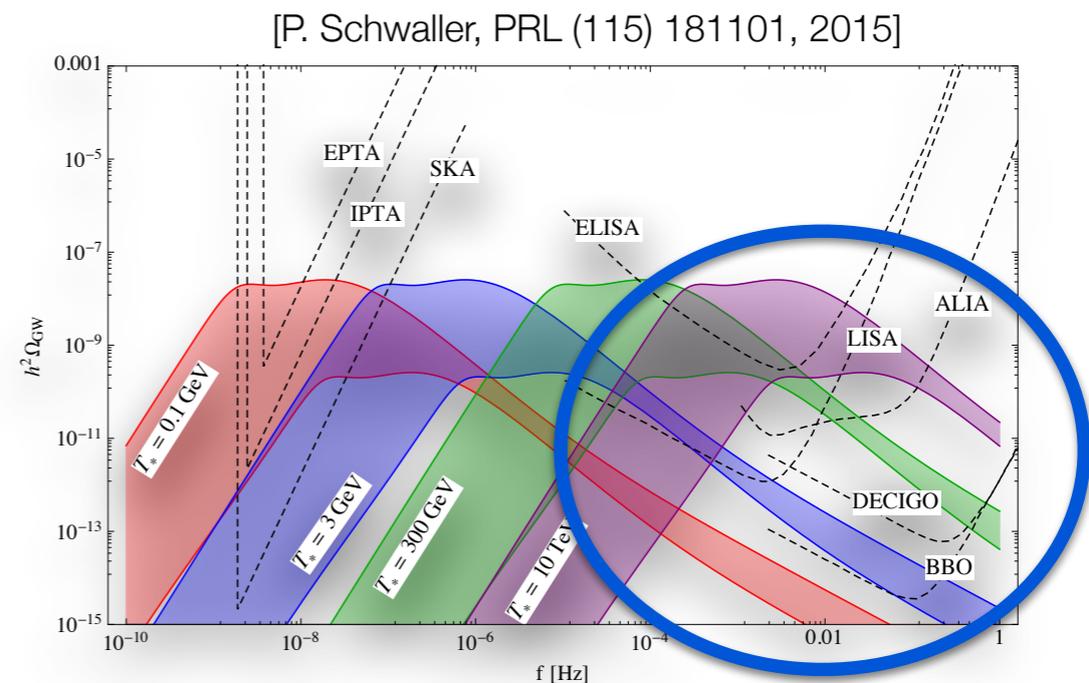


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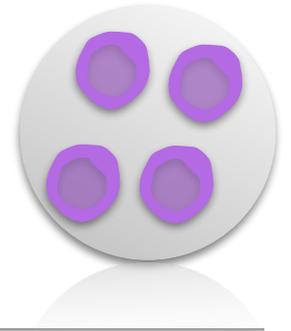




dark
matter

?

Backup Slides



“Stealth Dark Matter” model

- The field content of the model consists in **8 Weyl fermions**
- Dark fermions interact with the SM Higgs and obtain **current/chiral masses**
- Introduce **vector-like masses** for dark fermions that do not break EW symmetry
- Diagonalizing in the mass eigenbasis gives **4 Dirac fermions**
- Assume **custodial SU(2) symmetry** arising when **$u \leftrightarrow d$**

Field	$SU(N)_D$	$(SU(2)_L, Y)$	Q
$F_1 = \begin{pmatrix} F_1^u \\ F_1^d \end{pmatrix}$	\mathbf{N}	$(\mathbf{2}, 0)$	$\begin{pmatrix} +1/2 \\ -1/2 \end{pmatrix}$
$F_2 = \begin{pmatrix} F_2^u \\ F_2^d \end{pmatrix}$	$\overline{\mathbf{N}}$	$(\mathbf{2}, 0)$	$\begin{pmatrix} +1/2 \\ -1/2 \end{pmatrix}$
F_3^u	\mathbf{N}	$(\mathbf{1}, +1/2)$	$+1/2$
F_3^d	\mathbf{N}	$(\mathbf{1}, -1/2)$	$-1/2$
F_4^u	$\overline{\mathbf{N}}$	$(\mathbf{1}, +1/2)$	$+1/2$
F_4^d	$\overline{\mathbf{N}}$	$(\mathbf{1}, -1/2)$	$-1/2$

$$\mathcal{L} \supset + y_{14}^u \epsilon_{ij} F_1^i H^j F_4^d + y_{14}^d F_1 \cdot H^\dagger F_4^u - y_{23}^d \epsilon_{ij} F_2^i H^j F_3^d - y_{23}^u F_2 \cdot H^\dagger F_3^u + h.c.$$

$$\mathcal{L} \supset M_{12} \epsilon_{ij} F_1^i F_2^j - M_{34}^u F_3^u F_4^d + M_{34}^d F_3^d F_4^u + h.c.$$

$$y_{14}^u = y_{14}^d \quad y_{23}^u = y_{23}^d \quad M_{34}^u = M_{34}^d$$



Lattice results for Composite Dark Matter

Template Models	Spectrum	Higgs	Mag. Dip.	Charge r.	Polariz.
SU(2) $N_f=1$	★	★	[Francis, Hudspith, Lewis, Tulin 1809.09117]		
SU(2) $N_f=2$	★	★		★	★
SU(3) $N_f=2,6$	★		★	★	[Drach, et al. 1511.04370]
SU(3) $N_f=8$	★	★			
SU(3) $N_f=2$ (S)	★	[Fodor, et al. 1601.03302]			
SU(4) $N_f=4$	★	★			★
SO(4) $N_f=2$ (V)	★				
SU(N) $N_f=0$	★				



Lattice results for Composite Dark Matter

Template Models	Spectrum	Higgs	Mag. Dip.	Charge r.	Polariz.
SU(2) $N_f=1$	★	★	[Francis, Hudspith, Lewis, Tulin 1809.09117]		
SU(2) $N_f=2$	★	★	forbidden in pNGB DM	★	★
SU(3) $N_f=2,6$	★		★	★	[Drach, et al. 1511.04370]
SU(3) $N_f=8$	★	★			
SU(3) $N_f=2$ (S)	★	[Fodor, et al. 1601.03302]			
SU(4) $N_f=4$	★	★	forbidden in Stealth DM		★
SO(4) $N_f=2$ (V)	★				
SU(N) $N_f=0$	★	forbidden in SUNonia			