Disentangling the gravity dual of Yang-Mills theory from the lattice



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• Wish to understand confinement

conformal symmetry must be broken

- Non-perturbative: defied analytic attempts. Numerical progress.
- Use non-AdS/non-CFT
- Tools at disposal:
 - Wilson etc loops
 - Entanglement entropy
 - C-functions
- Holography is never going to solve confinement, but

we may learn valuable lessons

- Holography score card by comparison to YM simulations
- We need to gain understanding of entanglement entropy in gauge theories ↔ this is our first goal

• Entanglement roughly probes confinement [Nishioka-Takayanagi hep-th/0611035,Klebanov-Kutasov-Murugan 0709.2140,NJ-Subils 2010.09392]

- Effective degrees of freedom:
 - deconfining phase: colorful (e.g. gluons) $\sim \mathcal{O}(N_c^2)$
 - confining phase: color singlets (e.g. glueballs) $\sim \mathcal{O}(1)$
- Derived quantities of EE capture the number of dofs
- Key idea: C-function constructed from EE acts as an order parameter for deconfinement at $N_c \sim \infty$

Entanglement entropy itself motivates

- Quantum information theory (related entanglement measures)
- Universal order parameter for quantum phase transitions [Kitaev-Preskill hep-th/0510092]
- "Measured" in cold atom systems

[Islam et al. 1509.01160]

- $N_c \ll \infty$: Computable from the lattice [Buividovich-Polikarpov 0802.4247, ...]
- $N_c = \infty$: Holography
 - simple prescription via minimal surfaces [Ryu-Takayanagi hep-th/0603001,Hubeny-Rangamani-Takayanagi 0705.0016]
 - $\bullet\,$ relationship to black hole thermo \rightarrow get thermo for QFT

Definition of EE

• Take (scalar) QFT in \mathbb{R}^{d+1} in vacuum state $|0\rangle$ and split $A = \mathbb{R}^d \times I_\ell$, $B = \mathbb{R}^d \times (\mathbb{R} - I_\ell)$ ℓ

a

↓ ¢

• Quantum entanglement (von Neumann) entropy

 $S_{A}=-{
m Tr}
ho_{
m A}\log
ho_{
m A}$, reduced density matrix $ho_{
m A}={
m Tr}_{
m B}|0
angle\langle 0|$

• EE for region A is entropy seen by an observer unable to access dofs in B

Holographic EE and $q\bar{q}$ potential

• Consider SU($N_c \sim \infty$) in 4d and 3d

 \rightsquigarrow rough duals from stacks of D3- and D2-branes

• Compute potentials for quarks separated by L and entanglement entropies for slabs of width ℓ using holography. [Rey-Yee hep-th/9803001,Maldacena hep-th/9803002,Ryu-Takayanagi

hep-th/0603001]

Holographic results for 4d SYM

 $V(L) \propto 1/L$ $S_{RT} \propto 1/\ell^2$

• Holographic results for 3d SYM

$$V(L) \propto 1/L^{2/3}$$

 $S_{RT} \propto 1/\ell^{4/3}$

Existing results in the literature

EE for slabs from the lattice in

• Relation of EE to free energy / path integral

[Calabrese-Cardy'04&'05]

- Prescription for free energy measurements [Fodor'07,Endrödi-Fodor-Katz-Szabo'07]
- Pure glue SU(2) in 4d

[Buividovich-Polikarpov'08]

 Pure glue SU(3) in 4d [Nakagawa-Nakamura-Motoki-Zakharov'09,Itou-Nagata-Nakagawa-Nakamura-

Zakharov'11&'15]

• Pure glue SU(2,3,4) in 4d [Rabenstein-Bodendorfer-Buividovich-Schäfer'18]

Our work, pure glue $SU(N_c)$

- in 4d for $N_c = 3$ and 5 at T = 0
- in 3d for $N_c=2$ (upto $N_c=7$ running) at $T\gtrsim T_c$

$SU(N_c = 3)$, 4d

- Compares well with previous results in the literature [here shown Nakagawa-Nakamura-Motoki-Zakharov 1104.1011]
- Our statistical errors are ~invisible to eye (black markers).



$SU(N_c = 3)$, 4d

Comparing the C-function



vs. holography $S_A = \frac{N_c^2}{\epsilon^2} - \frac{N_c^2}{\ell^2} + \dots$ [Ryu-Takayanagi hep-th/0605073]



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$SU(N_c = 2)$, 3d

• New results for $SU(N_c)$ in 3d at non-zero temperature



- We confirm predictions from D2-brane background
 - at UV (small ℓ) $dS/d\ell \sim \ell^{-7/3}$
 - at IR (large ℓ) $dS/d\ell \sim {\rm const.} = S_{BH} \propto T^{7/3}$

- $\bullet\,$ This is applied AdS/CFT in reverse
- Idea: take boundary measurements and attempt to "guess" what is bulk
- Need either lattice simulations or experimental measurements
 - they are *noisy*
 - \Rightarrow one cannot find precise dual geometry
- How to deal with inherently noisy data?

[NJ-Pönni 2007.00010]

Bulk reconstruction from EE

- Make as little assumptions on bulk as possible: use RT
 - can only reconstruct some components of the bulk metric
- Can be shown, for a given metric that:

$$\frac{dS_{RT}}{d\ell} = \text{function}(r_*) \ , \ r_* = \text{tip of RT surface}$$

• Use Hamiltonian Monte Carlo to reconstruct the metric (and errors!)

$$\left\{ \left(\frac{1}{T^2V}\frac{dS}{d\ell}\right)_i, \ell_i, \sigma_i \right\}, \ i \in 1, \dots, N_{data}$$

$$ds^{2} = \left(\frac{r_{p}}{z}\right)^{\frac{5}{2}} \left(-\frac{1 - (z/z_{h})^{5}}{a(z)^{2}} dt^{2} + d\bar{x}^{2}\right) + \left(\frac{z}{r_{p}}\right)^{\frac{5}{2}} \frac{a(z)^{2}}{1 - (z/z_{h})^{5}} dz^{2} + r_{p}^{\frac{3}{2}} \sqrt{z} d\Omega_{6}^{2}$$
$$a(z) = 1 + \sum_{i=1}^{N_{\text{basis}}} a_{i} f_{i}(z)$$

Bulk reconstruction from EE

• Bulk metric (posteriors *a_i*, Newton's constant) reconstructed from 3d lattice data



- predictions w/ AdS/CFT w/ confidence intervals!
- static qq̄ potential using D2-brane Ansatz for the dilaton to switch frames



Quark-anti-quark potentials

We also tested D2-background further by computing the $q\bar{q}$ potential.



- If $ReV \sim L^{-10/3}$ then $ImV \sim 1/L$
- Can test following the method [Burnier-Kaczmarek-Rothkopf 1410.2546]



- Lattice data for EE or Wilson loops can be generated $N_c \ll \infty$
- This data can be used to reconstruct the dual geometry
- Allows new predictions in same QFT w/ confidence intervals
- New era has begun

Precision holography!

Thank you!