

Twist-2 and twist-3 GPDs from Lattice QCD

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Outline:

Introduction Quasi-GPDs: – how it works – twist-2 GPDs – twist-3 GPDs Prospects/conclusion

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Many thanks to my Collaborators for work presented here:

C. Alexandrou, S. Bhattacharya, M. Constantinou, J. Dodson

K. Hadjiyiannakou, K. Jansen, A. Metz, A. Scapellato, F. Steffens

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Nucleon structure

One of the central aims of hadron physics: to understand better nucleon structure.





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Nucleon structure

One of the central aims of hadron physics: to understand better nucleon structure.

- This is one of the crucial expectations from the Electron-Ion Collider (EIC)!
- In particular, we want to probe the 3D structure.
- Thus, we need to access new kinds of functions: GPDs, TMDs.
- Also higher-twist is of growing importance for the full picture.
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Lattice can provide *qualitative* and eventually *quantitative* knowledge of different functions and their moments:

- 1D: form factors
- 1D: parton distribution functions (PDFs)
- 3D: generalized parton distributions (GPDs)
- 3D: transverse momentum dependent PDFs (TMDs)
- 5D: Wigner function







Lattice QCD – what one should keep in mind



Introduction				
Nucleon	structure			

- x-dependence Quasi-PDFs
- GPDs
- GPDs
- Quasi-GPDs
- Results
- Summary

- Lattice QCD offers a way for a careful *ab initio* study of non-perturbative aspects of QCD.
- Its huge strength: possibility to control all systematic effects: *cut-off effects, finite volume effects, quark mass effects, isospin breaking, excited states, ...*
- For many aspects, already precision results with percent/per mille total uncertainty.
- However, many aspects (the difficult ones!) with only exploratory studies.





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- However, many aspects (the difficult ones!) with only exploratory studies.
- Difficult problems need time to:
 - \star find the proper way to address
 - ★ prove computational feasibility
 - \star optimize the computational method
 - acquire all data (long computations...) *
 - analyze all systematics \star
- Nucleon structure is mostly difficult... and very expensive computationally.
- Thus, do not expect miracles.





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- Nucleon structure is mostly difficult... and very expensive computationally.
- Thus, do not expect miracles.
- Overall, expect complementary role of lattice.
- Robust quantitative statements: *low moments, form factors.*
- *x*-dependence: breakthrough in recent years, but a long way to go to solid quantitative statements.







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$$Q(x,\mu_R) = \int_{-1}^{1} \frac{dy}{y} C\left(\frac{x}{y},\mu_F,\mu_R\right) q(y,\mu_F),$$
 some lattice observable





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 some lattice observable

- Matrix elements: $\langle N | \overline{\psi}(z) \Gamma F(z) \Gamma' \psi(0) | N \rangle$ with different choices of Γ, Γ' Dirac structures and objects F(z).
 - * hadronic tensor K.-F. Liu, S.-J. Dong, 1993
 - * auxiliary scalar quark U. Aglietti et al., 1998
 - * auxiliary heavy quark W. Detmold, C.-J. D. Lin, 2005
 - * auxiliary light quark V. Braun, D. Müller, 2007
 - * quasi-distributions X. Ji, 2013
 - * "good lattice cross sections" Y.-Q. Ma, J.-W. Qiu, 2014,2017
 - * **pseudo-distributions** A. Radyushkin, 2017
 - ★ "OPE without OPE" QCDSF, 2017

Lattice PDFs/GPDs: dynamical progress



Lattice PDFs/GPDs: dynamical progress



K. Cichy, Progress in x-dependent partonic distributions from lattice QCD, plenary talk LATTICE 2021, 2110.07440

- K. Cichy, Overview of lattice calculations of the x-dependence of PDFs, GPDs and TMDs, plenary talk of Virtual Tribute to Quark Confinement 2021, 2111.04552
- K. Cichy, M. Constantinou, A guide to light-cone PDFs from Lattice QCD: an overview of approaches, techniques and results, invited review for a special issue of Adv. High Energy Phys. 2019 (2019) 3036904, 1811.07248
- M. Constantinou, The x-dependence of hadronic parton distributions: A review on the progress of lattice QCD (would-be) plenary talk of LATTICE 2020, EPJA 57 (2021) 77, 2010.02445
- X. Ji et al., Large-Momentum Effective Theory, Rev. Mod. Phys. 93 (2021) 035005
- M. Constantinou et al., Parton distributions and LQCD calculations: toward 3D structure, PPNP 121 (2021) 103908





X. Ji, Parton Physics on a Euclidean Lattice, Phys. Rev. Lett. 110 (2013) 262002





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Main idea:







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Correlation along the ξ^- -direction: $q(x) = \frac{1}{2\pi} \int d\xi^- e^{-ixp^+\xi^-} \langle N | \overline{\psi}(\xi^-) \Gamma \mathcal{A}(\xi^-, 0) \psi(0) | N \rangle$ $|N\rangle - \text{nucleon at rest in the light-cone frame}$





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Main idea: ξ^{-} ξ^{-} ξ^{+} $\xi^{3} \equiv z$

Correlation along the ξ^- -direction: $q(x) = \frac{1}{2\pi} \int d\xi^- e^{-ixp^+\xi^-} \langle N | \overline{\psi}(\xi^-) \Gamma \mathcal{A}(\xi^-, 0) \psi(0) | N \rangle$ $|N \rangle$ – nucleon at rest in the light-cone frame Correlation along the $\xi^3 \equiv z$ -direction: $\tilde{q}(x) = \frac{1}{2\pi} \int dz \, e^{ixP_3z} \langle N | \overline{\psi}(z) \Gamma \mathcal{A}(z, 0) \psi(0) | N \rangle$ $|N \rangle$ – nucleon at rest in the standard frame Correlation along the ξ^3 -direction: $\tilde{q}(x) = \frac{1}{2\pi} \int dz \, e^{ixP_3z} \langle P | \overline{\psi}(z) \Gamma \mathcal{A}(z, 0) \psi(0) | P \rangle$ $|P \rangle$ – boosted nucleon





X. Ji, Parton Physics on a Euclidean Lattice, Phys. Rev. Lett. 110 (2013) 262002



Matching (Large Momentum Effective Theory (LaMET) X. Ji, Parton Physics from Large-Momentum Effective Field Theory, Sci.China Phys.Mech.Astron. 57 (2014) 1407 \rightarrow brings quasi-distribution to the light-cone distribution, up to power-suppressed effects:

$$\begin{split} \tilde{q}(x,\mu,P_3) &= \int_{-1}^{1} \frac{dy}{|y|} C\left(\frac{x}{y},\frac{\mu}{P_3}\right) q(y,\mu) + \mathcal{O}\left(\Lambda_{\rm QCD}^2/P_3^2, M_N^2/P_3^2\right) \\ \text{quasi-PDF} & \text{pert.kernel} \quad \text{PDF} & \text{higher-twist effects} \end{split}$$

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Twist-2 and twist-3 GPDs from LQCD – XVth CONF 2022 – 6 / 26





- Parton distribution functions (PDFs) formal definition: $f(x,\mu) = \frac{1}{2\pi} \int d\xi^- e^{-ixp^+\xi^-} \langle P | \overline{\psi}(\xi^-) \Gamma \mathcal{A}(\xi^-,0) \psi(0) | P \rangle$
- Generalized parton distributions (GPDs): $F(x,\xi,t,\mu) = \frac{1}{2\pi} \int d\xi^- e^{-ixp^+\xi^-} \langle P'' | \overline{\psi}(\xi^-) \Gamma \mathcal{A}(\xi^-,0) \psi(0) | P' \rangle$ The only difference: momentum transfer i.e. $P'' \neq P'$ (P'' = P' + Q, $t = -Q^2$).



Generalized parton distributions (GPDs)

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- GPDs reduce to PDFs in the forward limit, e.g. H(x, 0, 0) = q(x)
- Moments of GPDs are form factors, e.g. $\int dx H(x,\xi,t) = F_1(t)$
- Experimental access:
 - * PDFs Deep Inelastic Scattering (DIS) $ep \longrightarrow eX$
 - * GPDs Deeply Virtual Compton Scattering (DVCS) $ep \rightarrow e'p'\gamma$







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Quasi-GPDs: similar procedure to quasi-PDFs Important new aspect: 2 or 4 GPDs need to be disentangled, e.g. H and E:

 $\mathcal{M}(z,t,\xi;\,\mu_R;\,\Gamma,\overline{\Gamma}) = \mathcal{K}_H(\Gamma,\overline{\Gamma})H(z,t,\xi;\mu_R) + \mathcal{K}_E(\Gamma,\overline{\Gamma})E(z,t,\xi;\mu_R).$







x-dependence

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Quasi-GPDs lattice procedure





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most costly part of the procedure! needs several \vec{Q} vectors

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GPDs

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the final desired object!

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Renorm ME

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Setup Bare ME

Setup



Lattice setup:

- fermions: $N_f = 2$ twisted mass fermions + clover term
- gluons: Iwasaki gauge action, $\beta = 1.778$
- gauge field configurations generated by ETMC
- lattice spacing $a \approx 0.093$ fm,
- $32^3 \times 64 \Rightarrow L = 3$ fm,
- $m_{\pi} \approx 260$ MeV.

P_3	P_3 [GeV]	$N_{ m meas}$
$4\pi/L$	0.83	4152
$6\pi/L$	1.25	42080
$8\pi/L$	1.67	112192

Always: u - d flavor combination

ETMC, Phys. Rev. Lett. 125 (2020) 262001ETMC, Phys. Rev. D105 (2022) 034501S. Bhattacharya et al., 2112.05538





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- three nucleon boosts ($\xi = 0$): $P_3 = 0.83, 1.25, 1.67$ GeV,
- momentum transfer ($\xi = 0$): $-t = 0.69 \text{ GeV}^2$,



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- momentum transfer ($\xi = 0$): $-t = 0.69 \text{ GeV}^2$,
- nucleon boost ($\xi \neq 0$): $P_3 = 1.25$ GeV,
- momentum transfer $(\xi \neq 0)$: $-t = 1.02 \text{ GeV}^2$.



ETMC, Phys. Rev. Lett. 125 (2020) 262001ETMC, Phys. Rev. D105 (2022) 034501S. Bhattacharya et al., 2112.05538



Bare matrix elements



Lattice matrix elements need to be computed with 2 different projections (unpolarized/polarized). Below for the unpolarized Dirac insertion (for unpolarized GPDs)







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Removal of divergences and disentangling of H- and E-GPDs. Unpolarized Dirac insertion (for unpolarized GPDs)







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Light-cone distributions



Reconstruction of x-dependence and matching to light cone. Unpolarized Dirac insertion (for unpolarized GPDs)






Reconstruction of x-dependence and matching to light cone. Unpolarized Dirac insertion (for unpolarized GPDs)



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Comparison of PDFs and *H*-GPDs





$\frac{\text{unpolarized}}{\left[--H(r)-\text{GPD} \ f=0\right]} \quad \text{ETMC, Phys. Rev. Lett. 125 (2020) 262001}$



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Comparison of PDFs and *H*-**GPDs**





unpolarized ETMC, Phys. Rev. Lett. 125 (2020) 262001 $\begin{array}{c} & & \\ \hline & - & H(x) - \text{GPD}, \xi = 0 \\ \hline & - & H(x) - \text{GPD}, \xi = |1/3| \\ \hline & - & f_1(x) \\ \hline & P_3 = 1.25 \text{ GeV} \\ \hline & P_3 = 1.25 \text{ GeV} \\ \hline & \xi = 0, 1/3 \\ 0 \end{array}$

0.5

Important insights from models: S. Bhattacharya, C. Cocuzza, A. Metz

0

-0.5

Phys. Lett. B788 (2019) 453 Phys. Rev. D102 (2020) 054201

-1



Comparison of PDFs and *H*-GPDs





unpolarized ETMC, Phys. Rev. Lett. 125 (2020) 262001 3 - H(x)-GPD, $\xi = 0$ - H(x)-GPD, $\xi = |1/3|$ $-f_1(x)$ $P_3 = 1.25 \,\, {\rm GeV}$ 2 $P_3 = 1.25 \text{ GeV}$ $-t \models 0, 0.69, 1.02 \text{ GeV}^2$ $\xi = 0, 1/3$ -0.5 0.5 -1 0 x- H(x)-GPD $-f_1(x)$ $P_3 = 1.67 \text{ GeV}$ $P_3 = 1.67 \text{ GeV}$ $2 - t = 0, 0.69 \text{ GeV}^2$ $\xi = 0$ 0 -0.5 0.5 -1 0 x

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Transversity GPDs: ETMC, Phys. Rev. D105 (2022) 034501 $\stackrel{\downarrow}{4}$ 4 GPDs: H_T , E_T , \tilde{H}_T , \tilde{E}_T



Three nucleon boosts ($\xi = 0$): $P_3 = 0.83, 1.25, 1.67$ GeV Nucleon boost ($\xi \neq 0$): $P_3 = 1.25$ GeV

Momentum transfer ($\xi = 0$): $-t = 0.69 \text{ GeV}^2$ Momentum transfer ($\xi \neq 0$): $-t = 1.02 \text{ GeV}^2$





Transversity GPDs: ETMC, Phys. Rev. D105 (2022) 034501





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4 GPDs: H_T , E_T , \tilde{H}_T , \tilde{E}_T

Transversity GPDs





More fundamental quantity: $E_T + 2 \tilde{H}_T$









ETMC, Phys. Rev. D105 (2022) 034501

More fundamental quantity: $E_T + 2\tilde{H}_T$

- related to the transverse spin structure of the proton
- physically interpreted as lateral deformation in the distribution of transversely polarized quarks in an unpolarized proton
- lowest Mellin moment in the forward limit: transverse spin-flavor dipole moment in an unpolarized target (k_T)
- second moment related to the transverse-spin quark angular momentum in an unpolarized proton







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Moments of transversity GPDs



n = 0 Mellin moments:

$$\int_{-1}^{1} dx \, H_{T}(x,\xi,t) = \int_{-\infty}^{\infty} dx \, H_{Tq}(x,\xi,t,P_{3}) = A_{T10}(t) ,$$

$$\int_{-1}^{1} dx \, E_{T}(x,\xi,t) = \int_{-\infty}^{\infty} dx \, E_{Tq}(x,\xi,t,P_{3}) = B_{T10}(t) ,$$

$$\int_{-1}^{1} dx \, \widetilde{H}_{T}(x,\xi,t) = \int_{-\infty}^{\infty} dx \, \widetilde{H}_{Tq}(x,\xi,t,P_{3}) = \widetilde{A}_{T10}(t) ,$$

$$\int_{-1}^{1} dx \, \widetilde{E}_{T}(x,\xi,t) = \int_{-\infty}^{\infty} dx \, \widetilde{E}_{Tq}(x,\xi,t,P_{3}) = 0 ,$$
(1)

- lowest moments of GPDs skewness-independent,
- lowest moments of quasi-GPDs boost-independent.

n = 1 Mellin moments (related to GFF of one-derivative tensor operator):

$$\int_{-1}^{1} dx \, x \, H_{T}(x,\xi,t) = A_{T20}(t) ,$$

$$\int_{-1}^{1} dx \, x \, E_{T}(x,\xi,t) = B_{T20}(t) ,$$

$$\int_{-1}^{1} dx \, x \, \widetilde{H}_{T}(x,\xi,t) = \widetilde{A}_{T20}(t) , \qquad (3)$$

$$\int_{-1}^{1} dx \, x \, \widetilde{E}_{T}(x,\xi,t) = 2\xi \widetilde{B}_{T21}(t) , \qquad (2)$$

• skewness-dependence only in for \tilde{E}_T (only ξ -odd GPD).

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Moments of transversity GPDs



Moments of	$H_T(x,\xi=0,t=-0.69{ m GeV}^2)$			$H_T(x,\xi = 1/3, t = -1.02 \mathrm{GeV}^2)$
	$P_3 = 0.83 \text{ GeV}$	$P_3 = 1.25 \text{ GeV}$	$P_3 = 1.67 \text{ GeV}$	$P_3 = 1.25 \mathrm{GeV}$
H_{Tq}	0.65(4)	0.64(6)	0.81(10)	0.49(5)
H_T	0.69(4)	0.67(6)	0.84(10)	0.45(4)
xH_T	0.20(2)	0.21(2)	0.24(3)	0.15(2)
$A_{T10} (z = 0)$	0.65(4)	0.65(6)	0.82(10)	0.49(5)



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Mellin moments P_3 -independent, preserved by matching, suppressed with increasing -t.





Moments of	$H_T(x,\xi=0,t=-0.69{ m GeV}^2)$			$H_T(x,\xi = 1/3, t = -1.02 \mathrm{GeV}^2)$
	$P_3 = 0.83 \text{ GeV}$	$P_3 = 1.25 \text{ GeV}$	$P_3 = 1.67 \text{ GeV}$	$P_3 = 1.25 \mathrm{GeV}$
H_{Tq}	0.65(4)	0.64(6)	0.81(10)	0.49(5)
H_T	0.69(4)	0.67(6)	0.84(10)	0.45(4)
xH_T	0.20(2)	0.21(2)	0.24(3)	0.15(2)
$A_{T10} (z = 0)$	0.65(4)	0.65(6)	0.82(10)	0.49(5)

Mellin moments P_3 -independent, preserved by matching, suppressed with increasing -t.

Moments of	$E_T(x,\xi=0,t=-0.69{\rm GeV}^2)$		$H_T(x,\xi = 1/3, t = -1.02 \mathrm{GeV}^2)$	
	$P_3 = 0.83 \text{ GeV}$	$P_3 = 1.25 \text{ GeV}$	$P_3 = 1.67 {\rm GeV}$	$P_3 = 1.25 \mathrm{GeV}$
E_{Tq}		1.20(42)	2.05(65)	0.67(19)
E_T		1.15(43)	2.10(67)	0.73(19)
xE_T		0.06(4)	0.13(5)	0.11(11)
$B_{T10} \ (z=0)$	1.71(28)	1.22(43)	2.10(67)	0.68(19)
Moments of	$\widetilde{H}_T(x,\xi=0,t=-0.69\mathrm{GeV}^2)$			$\widetilde{H}_T(x,\xi = 1/3, t = -1.02 \mathrm{GeV}^2)$
	$P_3 = 0.83 \text{ GeV}$	$P_3 = 1.25 \text{ GeV}$	$P_3 = 1.67 {\rm GeV}$	$P_3 = 1.25 \mathrm{GeV}$
\widetilde{H}_{Tq}		-0.44(20)	-0.90(32)	-0.26(9)
\widetilde{H}_T		-0.42(21)	-0.92(33)	-0.27(9)
$x\widetilde{H}_T$		-0.17(8)	-0.30(10)	-0.05(5)
$\widetilde{A}_{T10} \ (z=0)$	-0.67(14)	-0.45(21)	-0.92(33)	-0.24(8)

Similar conclusions (but very large errors).



Comparison of PDFs and GPDs



ETMC, Phys. Rev. Lett. 125 (2020) 262001 ETMC, Phys. Rev. D105 (2022) 034501







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Krzysztof Cichy

Twist-2 and twist-3 GPDs from LQCD – XVth CONF 2022 – 21 / 26





PDFs/GPDs can be classified according to their twist, which describes the order in 1/Q at which they appear in the factorization of structure functions.

LT: twist-2 – probability densities for finding partons carrying fraction x of the hadron momentum.





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- no density interpretation,
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 - S. Bhattacharya et al., Phys. Rev. D102 (2020) 114025

BC-type sum rules S. Bhattacharya, A. Metz, 2105.07282

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a = 0.093 fm

Twist-3:

- $m_{\pi} = 260 \text{ MeV}$ TMF QUASI no density interpretation,
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- first exploration of twist-3 GPDs

S. Bhattacharya et al., 2112.05538



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First exploration of twist-3 GPDs



Very recently, we combined our explorations of GPDs and of twist-3 distributions S. Bhattacharya et al., 2112.05538

$$\begin{split} \text{Twist-3 axial GPDs:} \ \ \widetilde{G}_1, \ \widetilde{G}_2, \ \widetilde{G}_3, \ \widetilde{G}_4 \\ h_{\gamma^j \gamma_5} &= \langle \langle \frac{g_{\perp}^{j\rho} \Delta_\rho \gamma_5}{2m} \rangle \rangle [F_{\widetilde{E}} + F_{\widetilde{G}_1}] + \langle \langle g_{\perp}^{j\rho} \gamma_\rho \gamma_5 \rangle \rangle [F_{\widetilde{H}} + F_{\widetilde{G}_2}] + \langle \langle \frac{g_{\perp}^{j\rho} \Delta_\rho \gamma^+ \gamma_5}{P^+} \rangle \rangle F_{\widetilde{G}_3} + \langle \langle \frac{i\epsilon_{\perp}^{j\rho} \Delta_\rho \gamma^+}{P^+} \rangle \rangle F_{\widetilde{G}_4} \,. \end{split}$$

Bare ME: (same lattice setup)



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First exploration of twist-3 GPDs



Contributions from different insertions and projectors $(\vec{Q} = (Q_x, 0, 0))$:

```
\Pi(\gamma^2\gamma^5,\Gamma_0): \widetilde{H} + \widetilde{G}_2 \text{ and } \widetilde{G}_4,

\Pi(\gamma^2\gamma^5,\Gamma_2): \widetilde{H} + \widetilde{G}_2 \text{ and } \widetilde{G}_4,

\Pi(\gamma^1\gamma^5,\Gamma_1): \widetilde{H} + \widetilde{G}_2 \text{ and } \widetilde{E} + \widetilde{G}_1,

\Pi(\gamma^1\gamma^5,\Gamma_3): \widetilde{G}_3.
```



S. Bhattacharya et al., 2112.05538

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0.5

0

x

-10

S. Bhattacharya et al., 2112.05538

-1

-0.5



-4

-1

-0.5

 $P_3 = 1.25 \text{ GeV}$

 $-P_3 = 1.67 \,\,\mathrm{GeV}$

0.5

1

0

x



First exploration of twist-3 GPDs





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Twist-2 and twist-3 GPDs from LQCD – XVth CONF 2022 – 25 / 26



Conclusions and prospects



Introduction

Results

Summary

• Enormous progress in lattice calculations of GPDs!



Conclusions and prospects



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- GPDs much more challenging than PDFs:
 - \star signal decays with increasing -t,
 - \star separate calculations for different -t,
 - \star discreteness of -t,
 - \star several projectors needed to disentangle GPDs,
 - non-zero skewness enhanced power corrections at the ERBL-DGLAP boundary.




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Thank you for your attention!