

Aspects of chiral transition in a Hadron Resonance Gas Model

Deeptak Biswas

The Institute of Mathematical Sciences

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In collaboration with Peter Petreczky, Sayantan Sharma

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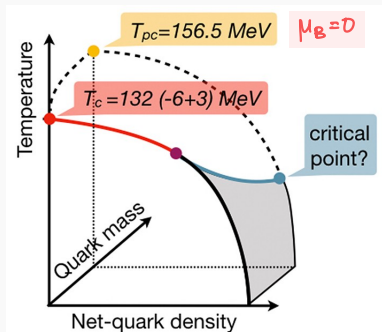
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Introduction

Motivation

- In QCD with 2 massless quarks, $SU(2)_V \times SU(2)_A \times U(1)_V \rightarrow$ Exact symmetry
- With physical mass, $SU(2)_V \times SU(2)_A \times U(1)_V \rightarrow$ Approximate (good)
- This symmetry is spontaneously broken to $SU(2)_V \times U(1)_V$
- For 2 flavors at non-zero mass, the chiral symmetry is restored via analytic crossover at $T_c = 156.5(1.5)$ MeV. [HotQCD 2018]
- How far can we estimate T_c and pseudo-critical line in Hadron Resonance Gas(HRG) model?



Chiral Condensate in the Hadron Resonance Gas model

Earlier results from χ_{PT} and HRG:

- ➡ The earliest estimation of pseudo-critical temperature, done within the NNLO chiral perturbation theory(χ_{PT}) gave $T_c = 250$ MeV.

[P. Gerber, H. Leutwyler 1989]

- ➡ Lowered to about 190 MeV with inclusion of heavier hadrons.
- ➡ Recent studies within the HRG have found a higher $T_c \sim 170$ MeV.

[J. Jankowski et al. 2013, A. N Tawfik, N. Magdy 2015]

Renormalized chiral condensate

- We can define the renormalized chiral condensate from the pressure as,

$$-m_s [\langle \bar{\psi}\psi \rangle_{I,T} - \langle \bar{\psi}\psi \rangle_{I,0}] = -m_s \frac{\partial P}{\partial m_l}$$

- The normalization is not unique [BMW 2010] ,

$$\langle \bar{\psi}\psi \rangle_R = -\frac{m_l}{m_\pi^4} [\langle \bar{\psi}\psi \rangle_{I,T} - \langle \bar{\psi}\psi \rangle_{I,0}] .$$

- A natural choice for dimensionless condensate [HotQCD 2012],

$$\Delta_R^I = d + m_s r_1^4 [\langle \bar{\psi}\psi \rangle_{I,T} - \langle \bar{\psi}\psi \rangle_{I,0}]$$

- Using low energy constant of $SU(2)$ χ_{PT} , $\Sigma^{1/3} = 272(5)$ MeV, $m_s = 92.2(1.0)$ MeV, and $r_1 = 0.3106$ fm, one gets $d = 0.022791$. [FLAG 2022],

⇒ The renormalized chiral condensate,

$$m_s \frac{\partial P}{\partial m_l} = -\frac{m_s}{m_l} \sum_{\alpha} \frac{g_{\alpha}}{2\pi^2} \int_0^{\infty} dp \, p^2 \, n_{\alpha}(E_{\alpha}) \frac{1}{2E_{\alpha}} m_l \frac{\partial M_{\alpha}^2}{\partial m_l}.$$

⇒ The non-trivial ingredient is $\partial M_{\alpha}^2 / \partial m_l$.

⇒ From $SU(2)$ χ_{PT} ,

$$M_\pi^2 = M^2 \left[1 - \frac{1}{2} \zeta \bar{l}_3 + \mathcal{O}(\zeta^2) \right] , \quad \zeta = \frac{M^2}{16\pi^2 F_\pi^2}$$

⇒ Kaon properties are predicted well from 2+1 χ_{PT}

[RBC 2014, Durr 2015]

$$M_K^2 = B_K(m_s) m_s \left[1 + \frac{\lambda_1(m_s) + \lambda_2(m_s)}{F^2} M^2 \right]$$

$$M^2 = 2Bm_l, \quad B = \Sigma/F^2$$

⇒ From LQCD the pion mass is consistent with LO result

$$M_\pi^2 \approx 2Bm_l. \quad [\text{RQCD Bali et al. 2016}] .$$

Sigma terms for Heavier hadrons

$$\rightarrow \sigma_\alpha = m_l \frac{\partial M_\alpha}{\partial m_l} \Big|_{m_l=m_l^{phys}} = m_l \langle \alpha | \bar{u}u + \bar{d}d | \alpha \rangle = M_\pi^2 \frac{\partial M_\alpha}{\partial M_\pi^2} \Big|_{M_\pi=M_\pi^{phys}}.$$

N	Λ	Σ	Ξ
44(3)(3)	31(1)(2)	25(1)(1)	15(1)(1)
Δ	Σ^*	Ξ^*	Ω^-
29(9)(3)	18(6)(2)	10(3)(2)	5(1)(1)

The sigma terms of ground state baryons have been only recently calculated with precision. [\[Copeland et al. 2021\]](#) .

New development from our work:

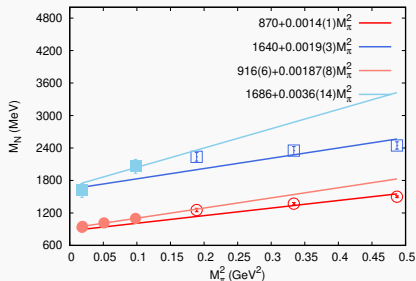
- ➡ We have done extensive compilation of the LQCD results to find $M_\pi^2 \frac{\partial M_\alpha}{\partial M_\pi^2}$ at a constant m_s , set at the physical value.
- ➡ For the first time, σ terms for η , $\rho(770)$, $K^*(892)$, and η' have been calculated from LQCD data.

[RQCD Bali et al. 2016, D. Guo et al. 2016, RQCD Bali et al. 2021] .

- ➡ We have assigned sigma terms for all meson resonances,
 - Iso-vector mesons $\rightarrow \sigma_{\rho(770)}$
 - Open strange mesons $\rightarrow \sigma_{K^*(892)}$.
 - Iso-scalar mesons \rightarrow corresponding ground states mesons σ terms.
- ➡ $f_0(500)$ is not considered as cancellation by the repulsive interactions [Broniowski et al. 2015].

Sigma terms for Baryon resonances: Nucleons

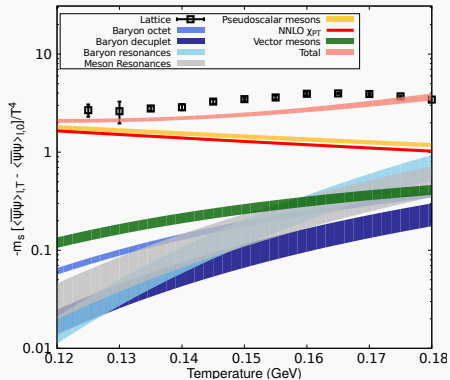
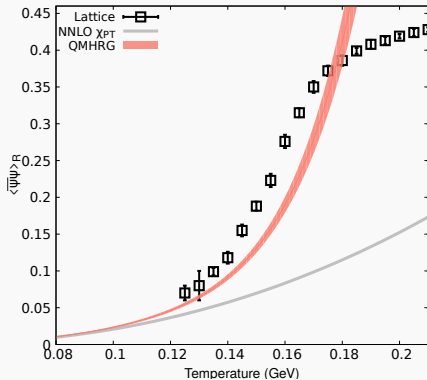
- ➡ It is difficult to measure baryon resonances in LQCD as they are close to the scattering state and resonances.
- ➡ For the excited N state, the fit to $2+1$ flavor LQCD data gives $\sigma = 68(27)$ MeV.
- ➡ Within large errors is consistent with the sigma term of its ground state.



- We have considered the σ terms for all resonances (even for strange baryons) to be same as the ground state.
- To reliably account for large uncertainty in σ of high mass resonances, we have taken the relative errors in the σ -terms of excited states to be 50%.
- However such large uncertainty contributes to only 10% of the total error in the renormalized chiral condensate as the dominant contribution comes from ground state pseudo-scalar mesons.

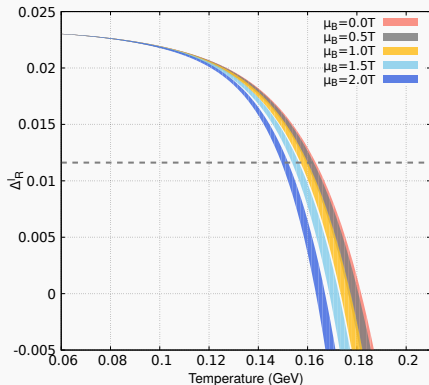
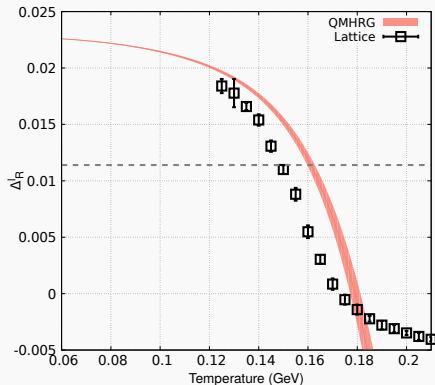
Results

Chiral condensate: LQCD vs. HRG model



HRG model calculations are consistent with LQCD continuum estimates till $T \sim 140\text{MeV}$.

$$\Delta_R' = d + m_s r_1^4 [\langle \bar{\psi} \psi \rangle_{I,T} - \langle \bar{\psi} \psi \rangle_{I,0}]$$



- On the lattice Δ_R' goes to half of its low-temperature value at T_c .
- We use this fact to estimate T_c from our HRG model calculations.

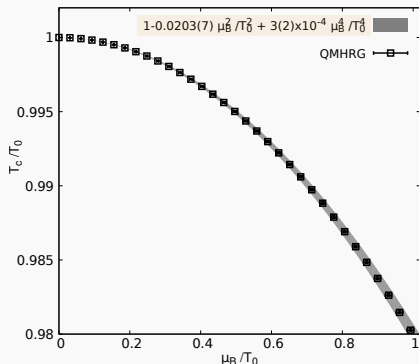
Summary of results

- Our improved HRG calculation gives $T_c = 161.2 \pm 1.7$ MeV at $\mu_B = 0$.
- Lattice QCD results on T_c in the continuum limit, $T_c = 156.5 \pm 1.5$ MeV [HotQCD 2018, BMW 2020]

Curvature of the pseudo-critical line

We extract κ_2 and κ_4 by fitting $T_c(\mu_B)$ for $0 < \mu_B/T_c(\mu_B = 0) < 1$.

- Our estimation $\kappa_2 = 0.0203(7)$.
- $\kappa_4 = -3(2) \times 10^{-4}$ is quite noisy.
- Highlight that our results are in very good agreement with LQCD estimates of $\kappa_2 = 0.012(4)$ [HotQCD 2018] and $0.0153(18)$ [BMW 2020], $\kappa_4 = 0$.

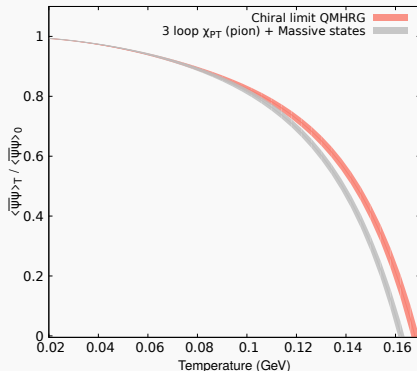


Transition in the chiral limit

- ➡ 3-loop χ_{PT} for pions + hadrons, gave a $T_c^0 \sim 170$ MeV .

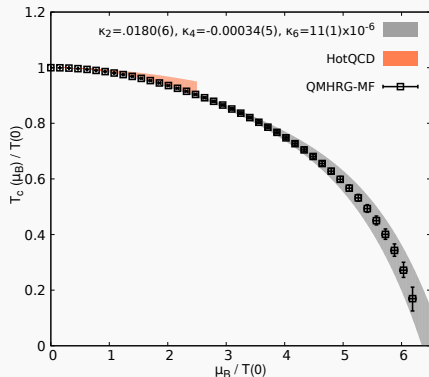
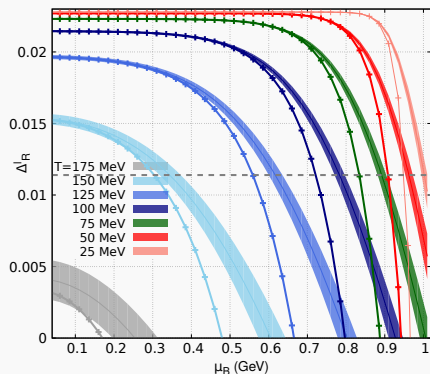
[P. Gerber, H. Leutwyler 1989]

- ➡ 3-loop χ_{PT} + our improved HRG estimates lower it to 162 MeV.
- ➡ LQCD predicts $T_c = 132_{-6}^{+3}$ MeV [HotQCD 2019].
- ➡ Need to go beyond 3-loop χ_{PT} and include temperature dependent width of the critical f_0 mode to improve the agreement



Extending our results to higher baryon density

For high density, we need to include repulsive interaction among baryons via mean field. [Huovinen, Petreczky 2018]



Summary and outlook

- We have studied chiral observables for physical hadrons within the HRG model.
- For the first time, precise values of σ terms for ρ, η, K^* , isoscalar mesons and ground state baryons have been included.
- This has successfully improved the T_c from HRG model, bringing it closer to the LQCD estimates.
- Curvature coefficients κ_2, κ_4 are very close to lattice results than previous estimates.
- With these excellent agreements at $\mu_B \approx 0$, we have extended our formalism to larger values of μ_B . [P.Petreczky, to appear]