

Charge dependent flow, early magnetic field and chiral magnetic effect in heavy ion collisions at ALICE



Shi Qiu



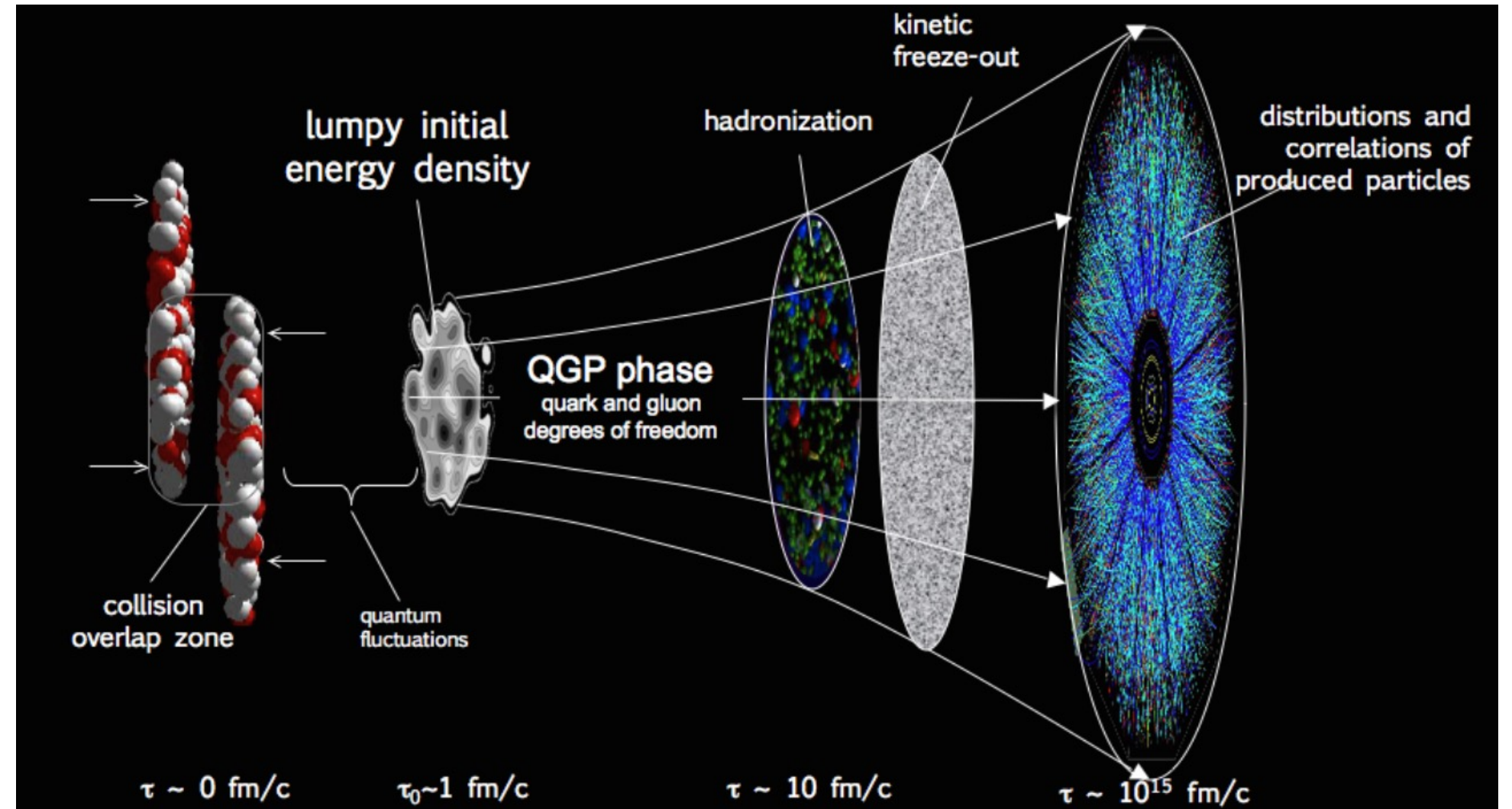
Nikhef

The XVth Quark Confinement and the Hadron Spectrum Conference, 1/8/2022

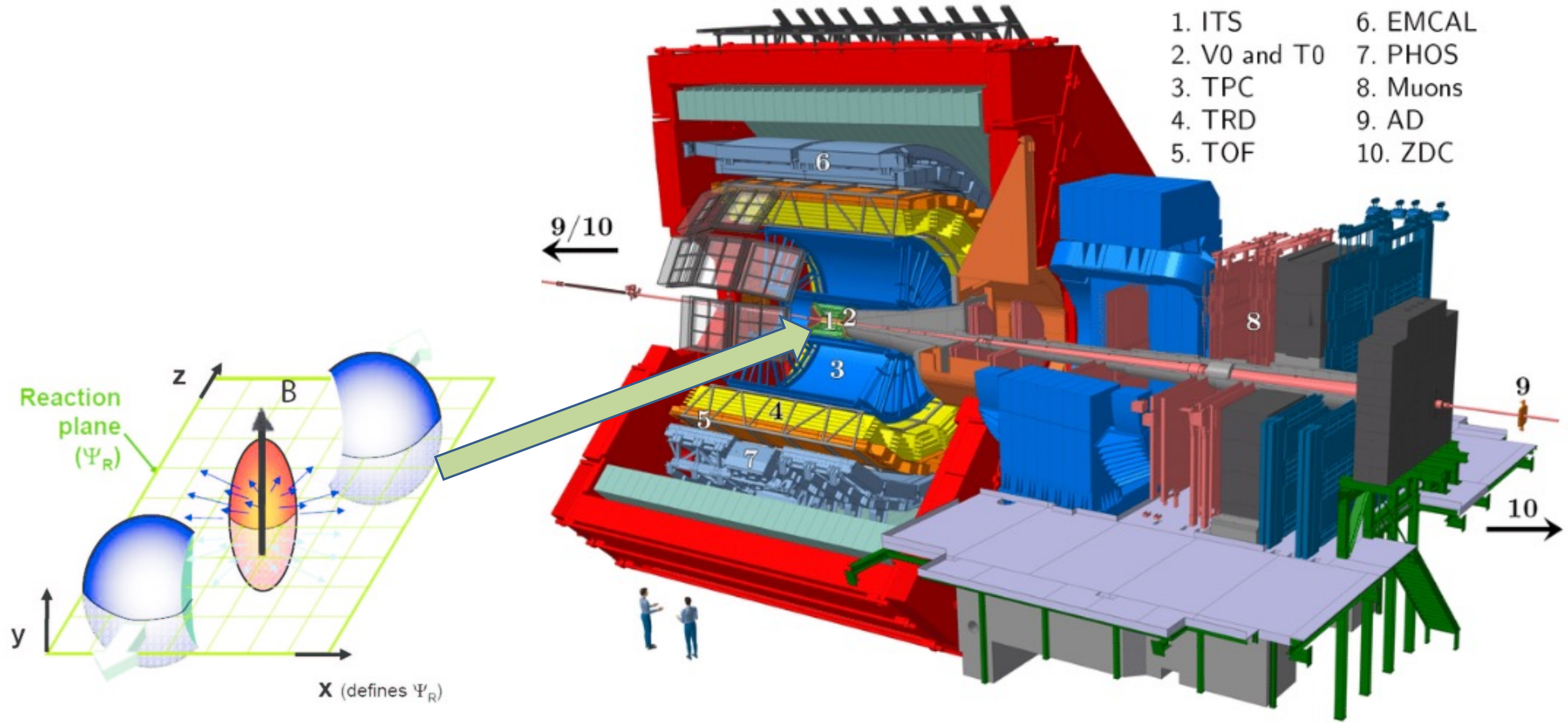


Content of the talk

- Probe electromagnetic field
 - Charge-dependent directed flow
 - Inclusive hadron measurements
 - Heavy flavor measurements (D^0 & \bar{D}^0)
 - Hyperon Polarization
- Searches for chiral anomalies
 - Chiral magnetic effect studies
 - Mix harmonic
 - Event shape engineering

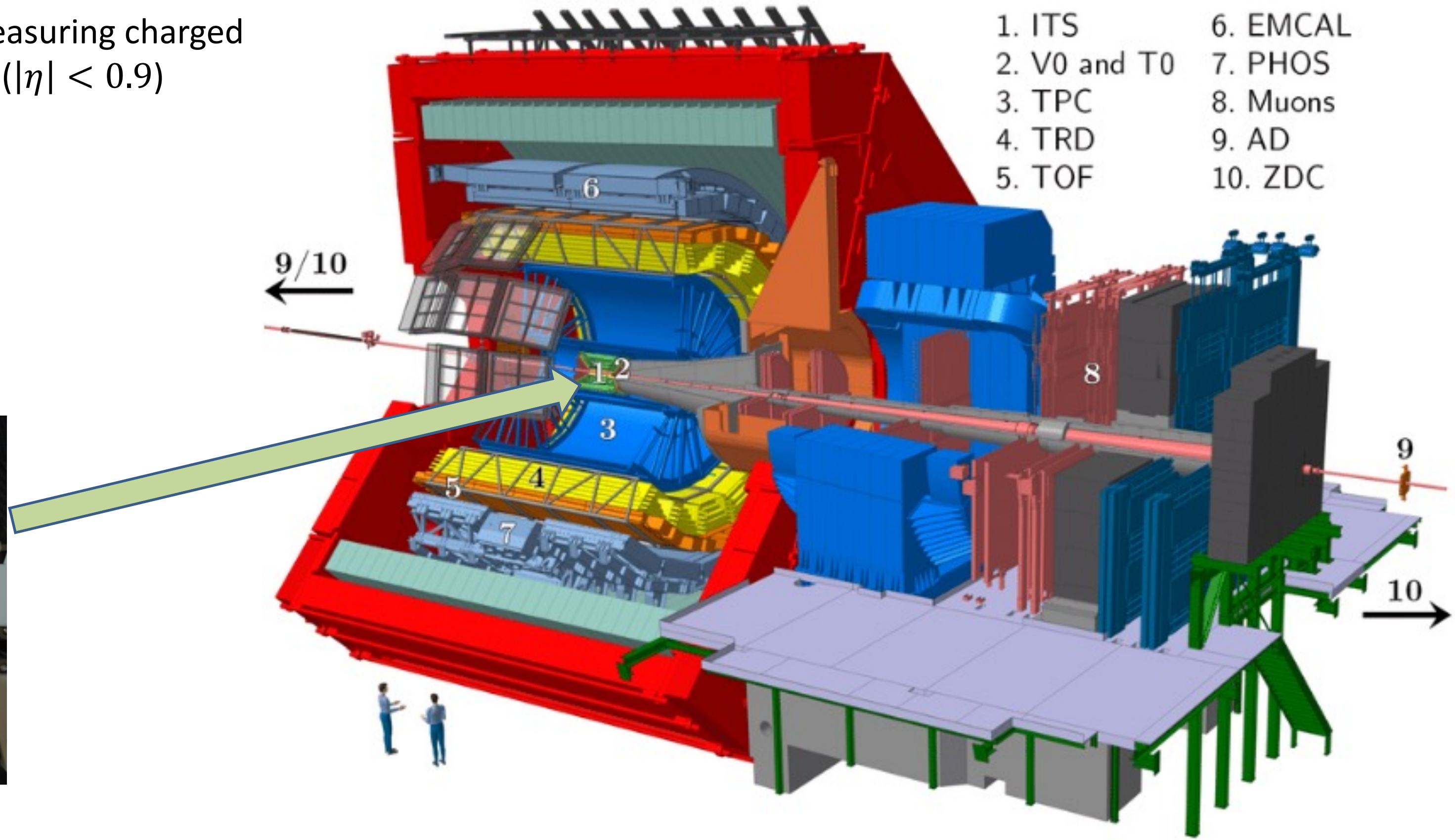
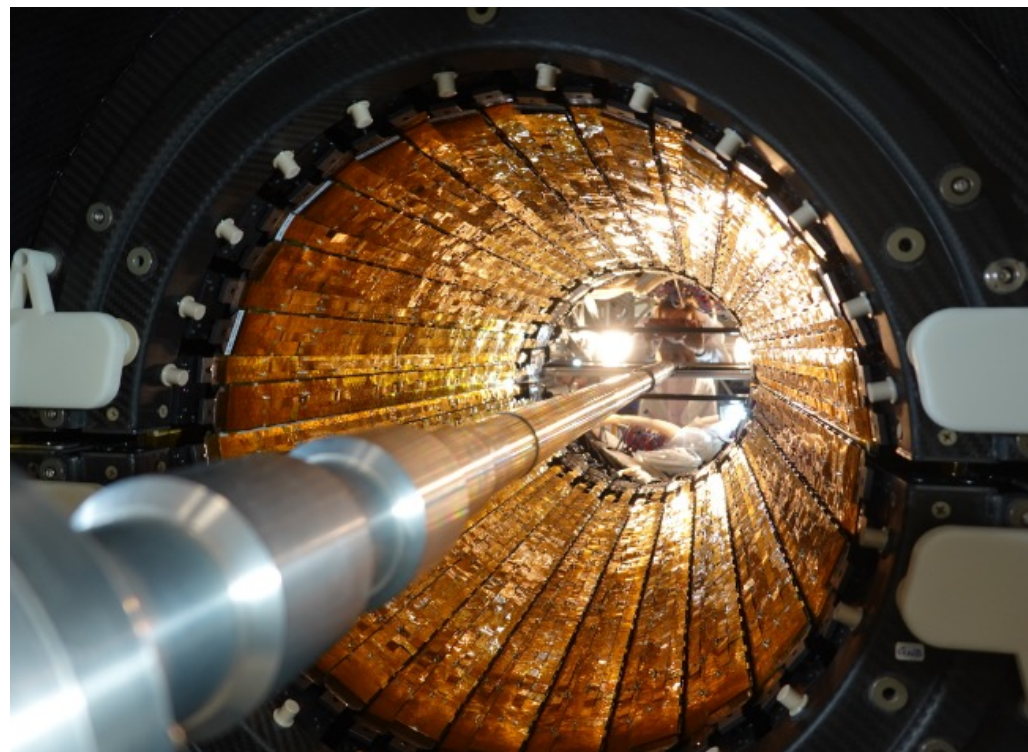


A Large Ion Collider Experiment



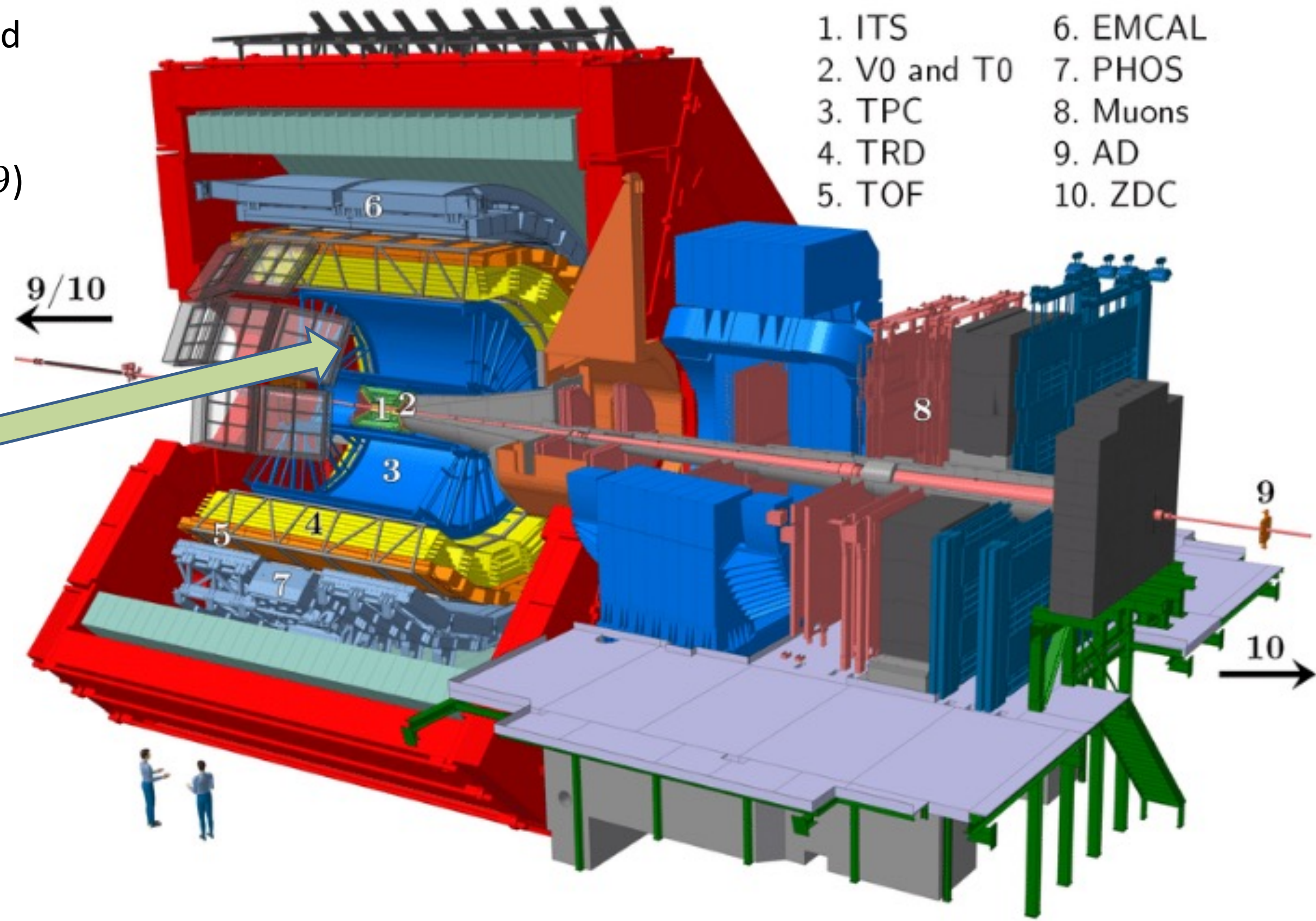
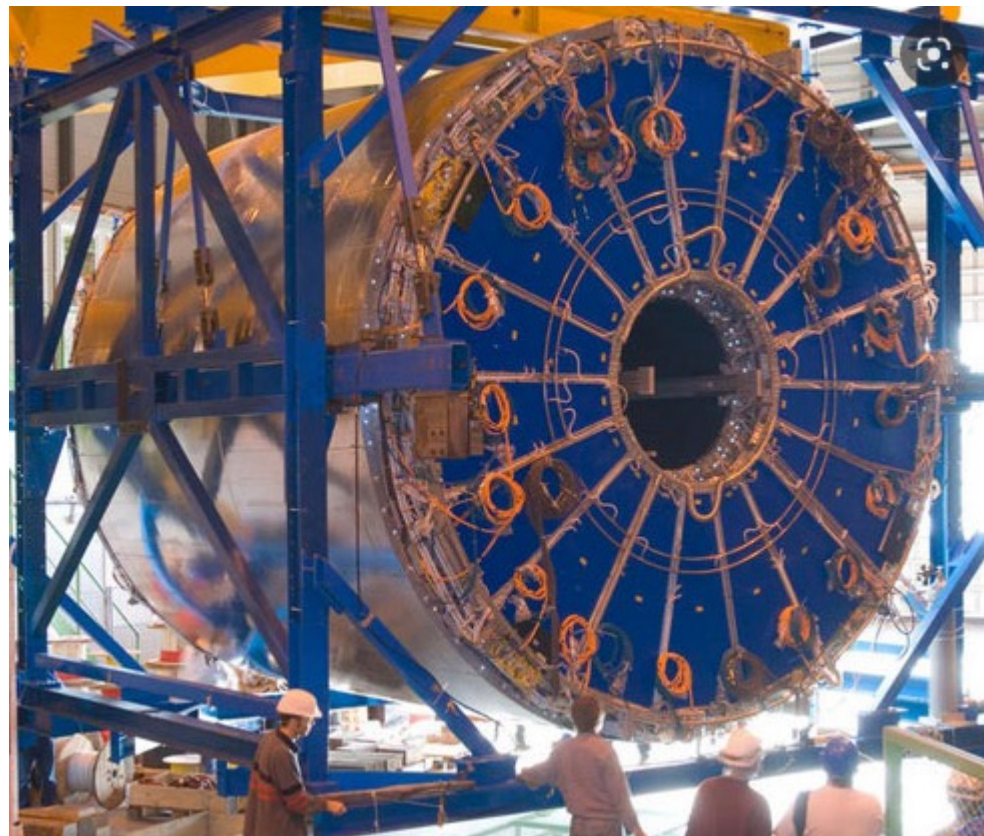
A Large Ion Collider Experiment

- ITS: semiconductor detector measuring charged particle traversing its segments ($|\eta| < 0.9$)



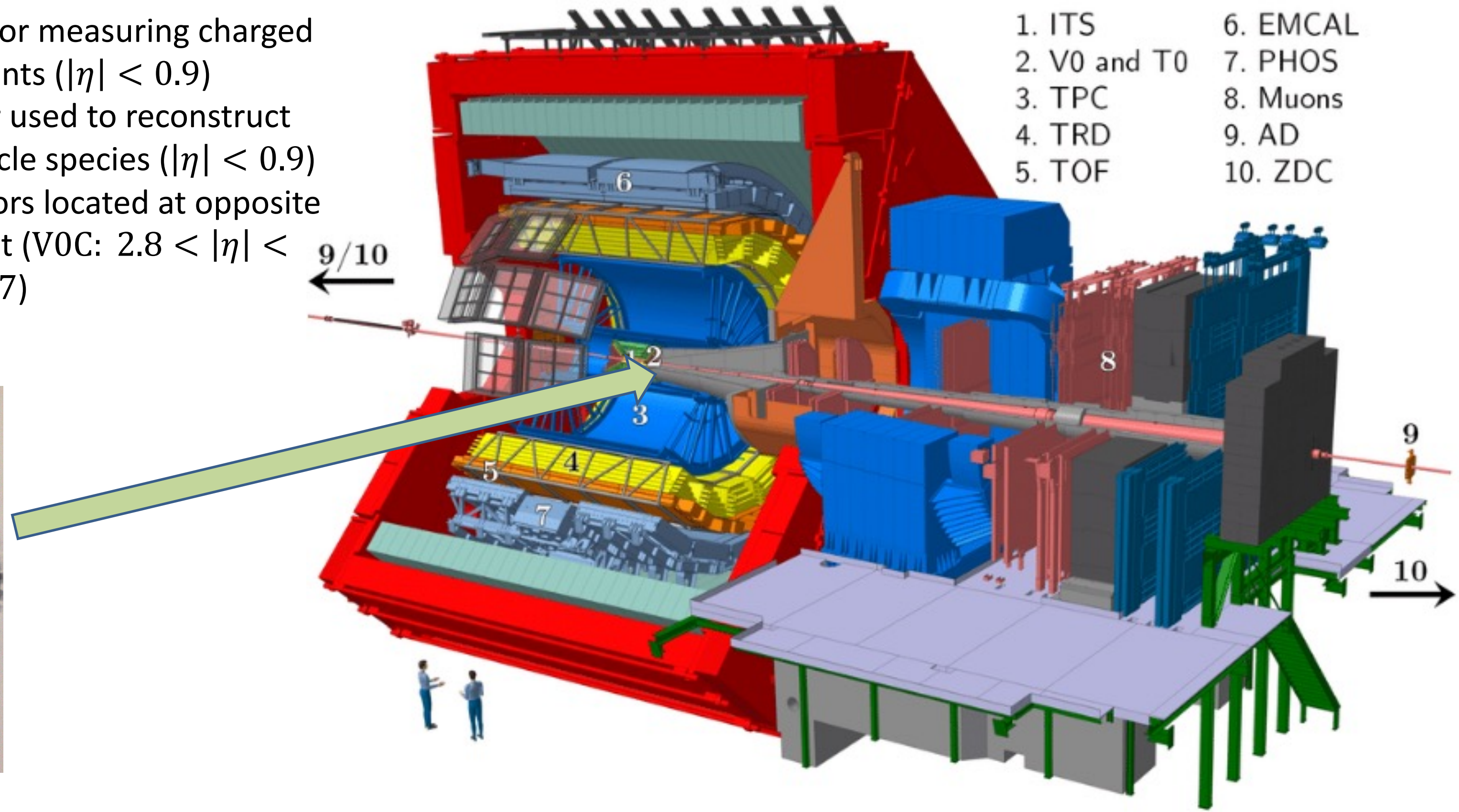
A Large Ion Collider Experiment

- ITS: a semiconductor detector measuring charged particle traversing its segments ($|\eta| < 0.9$)
- TPC: cylindrical gas detector used to reconstruct tracks and identify the particle species ($|\eta| < 0.9$)



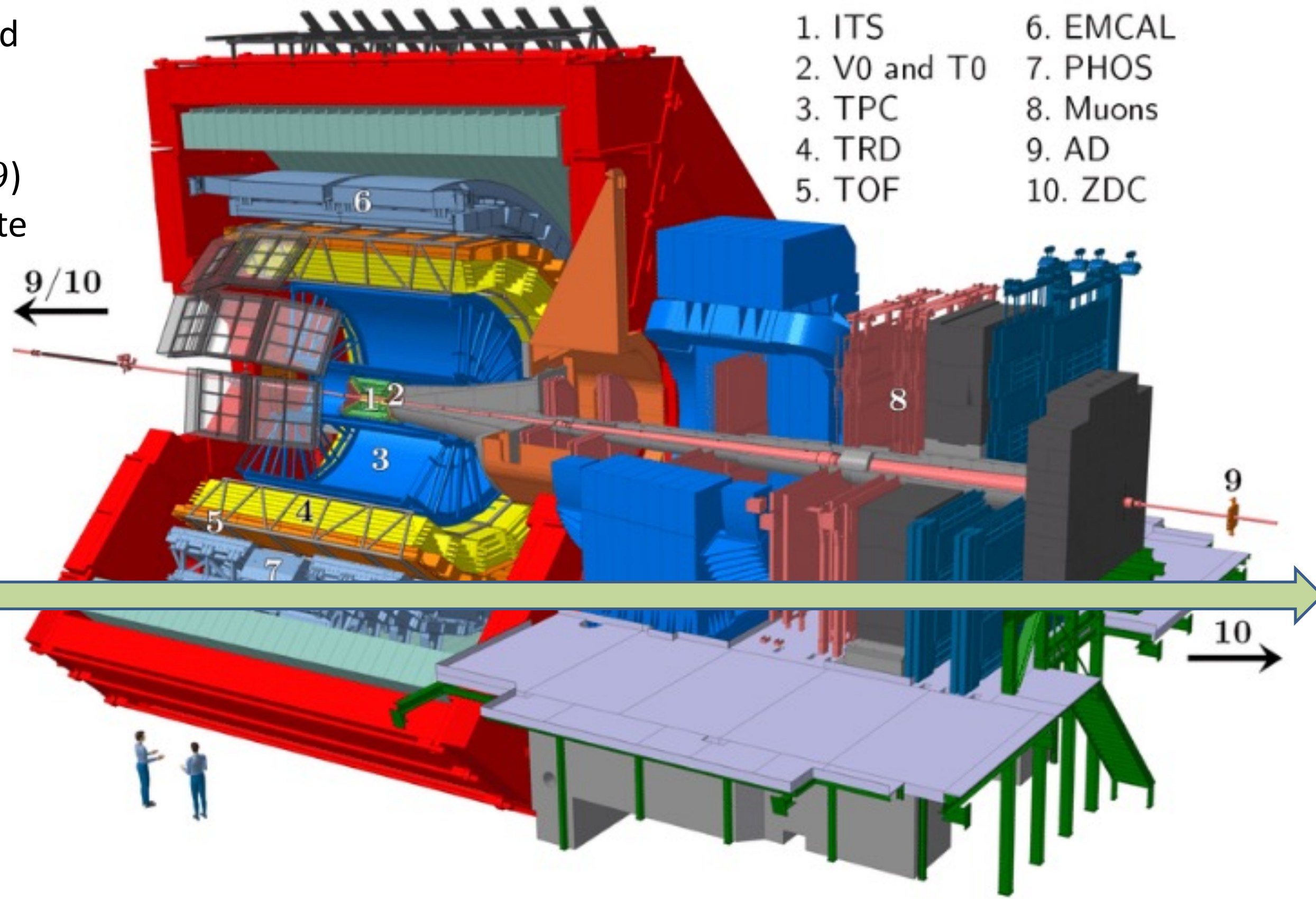
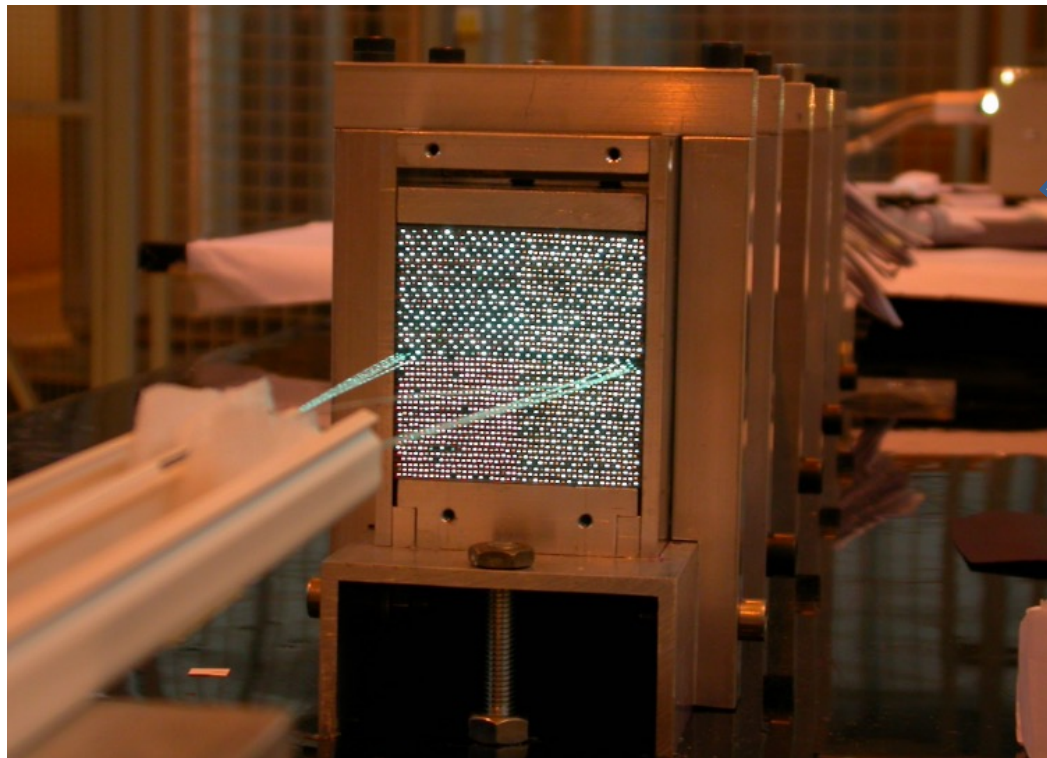
A Large Ion Collider Experiment

- ITS: a semiconductor detector measuring charged particle traversing its segments ($|\eta| < 0.9$)
- TPC: cylindrical gas detector used to reconstruct tracks and identify the particle species ($|\eta| < 0.9$)
- V0: two cylindrical scintillators located at opposite sides of the interaction point (V0C: $2.8 < |\eta| < 5.1$, V0A: $-3.7 < |\eta| < -1.7$)



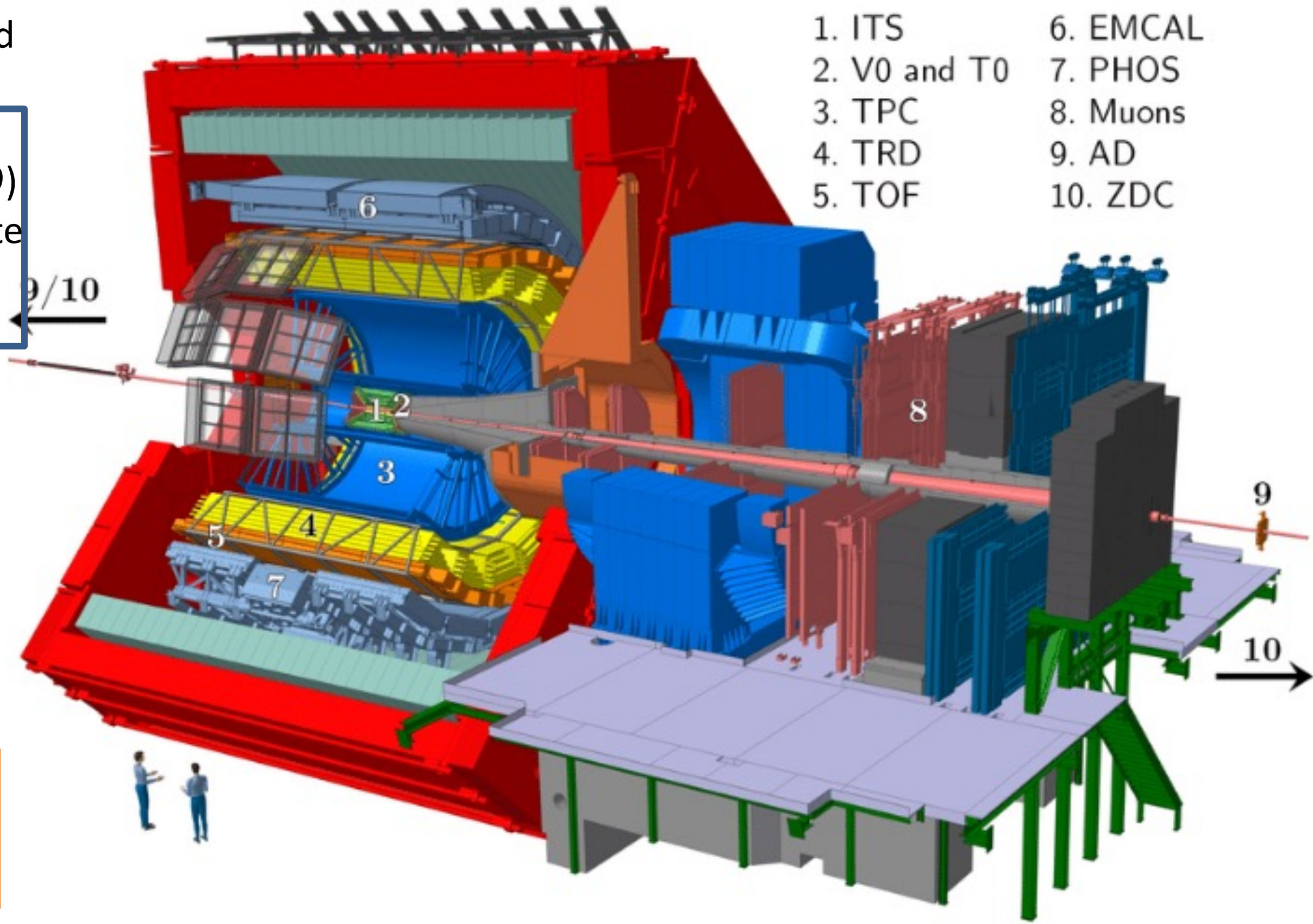
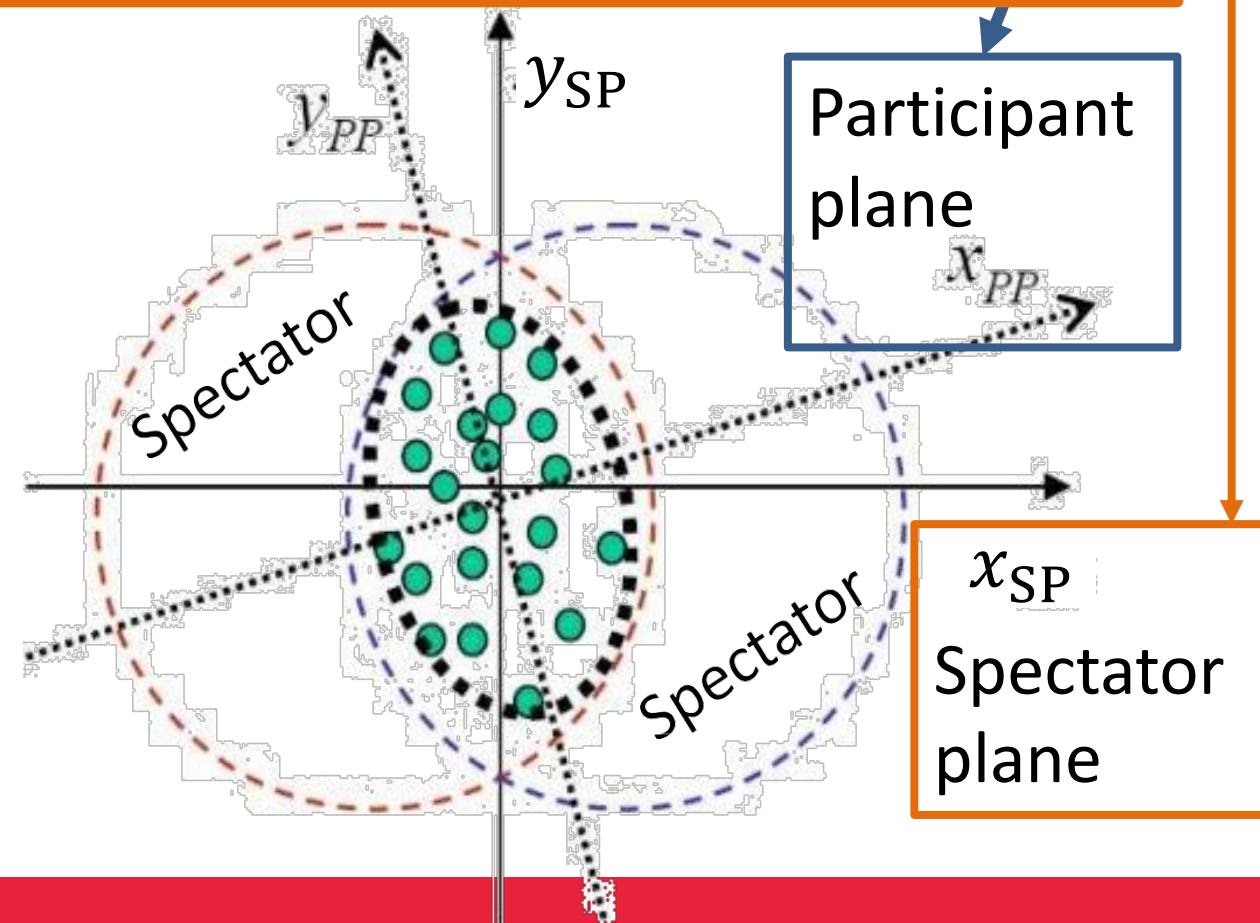
A Large Ion Collider Experiment

- ITS: a semiconductor detector measuring charged particle traversing its segments ($|\eta| < 0.9$)
- TPC: cylindrical gas detector used to reconstruct tracks and identify the particle species ($|\eta| < 0.9$)
- V0: two cylindrical scintillators located at opposite sides of the interaction point (V0C: $2.8 < \eta < 5.1$, V0A: $-3.7 < \eta < -1.7$)
- ZDC: hadronic calorimeters located at beam rapidities ($4.8 < |\eta| < 5.7$)



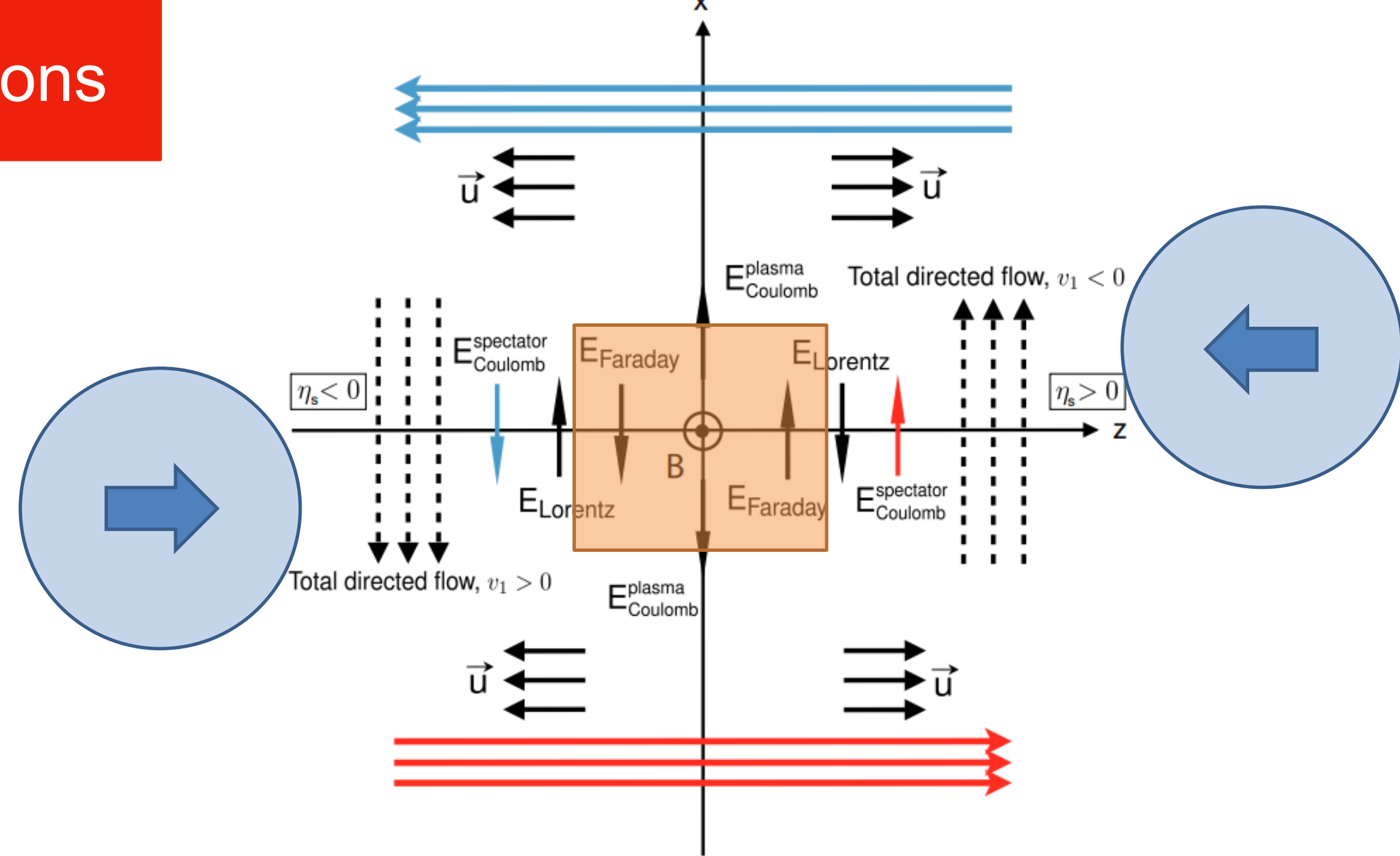
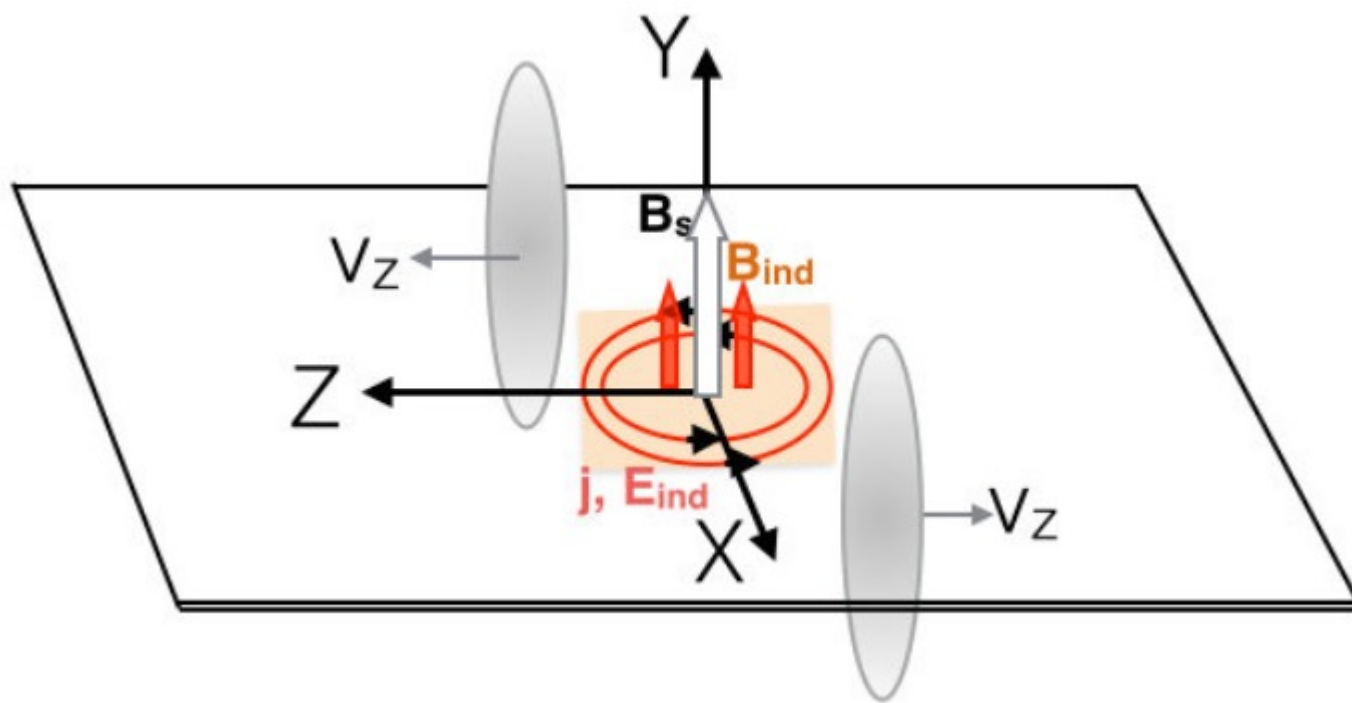
A Large Ion Collider Experiment

- ITS: a semiconductor detector measuring charged particle traversing its segments ($|\eta| < 0.9$)
- TPC: cylindrical gas detector used to reconstruct tracks and identify the particle species ($|\eta| < 0.9$)
- V0: two cylindrical scintillators located at opposite sides of the interaction point (V0C: $2.8 < |\eta| < 5.1$, V0A: $-3.7 < |\eta| < -1.7$)
- ZDC: hadronic calorimeters located at beam rapidities ($4.8 < |\eta| < 5.7$)



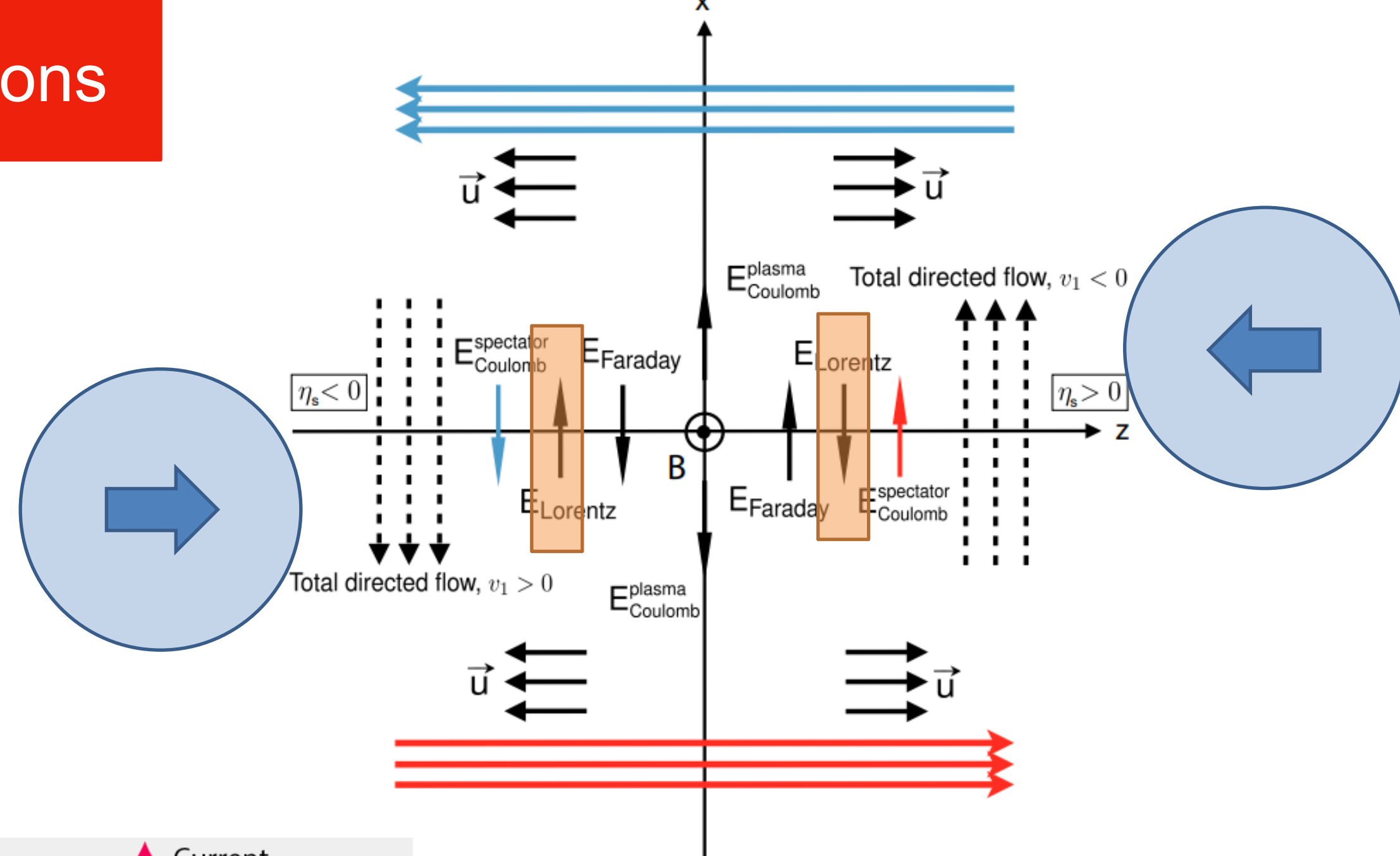
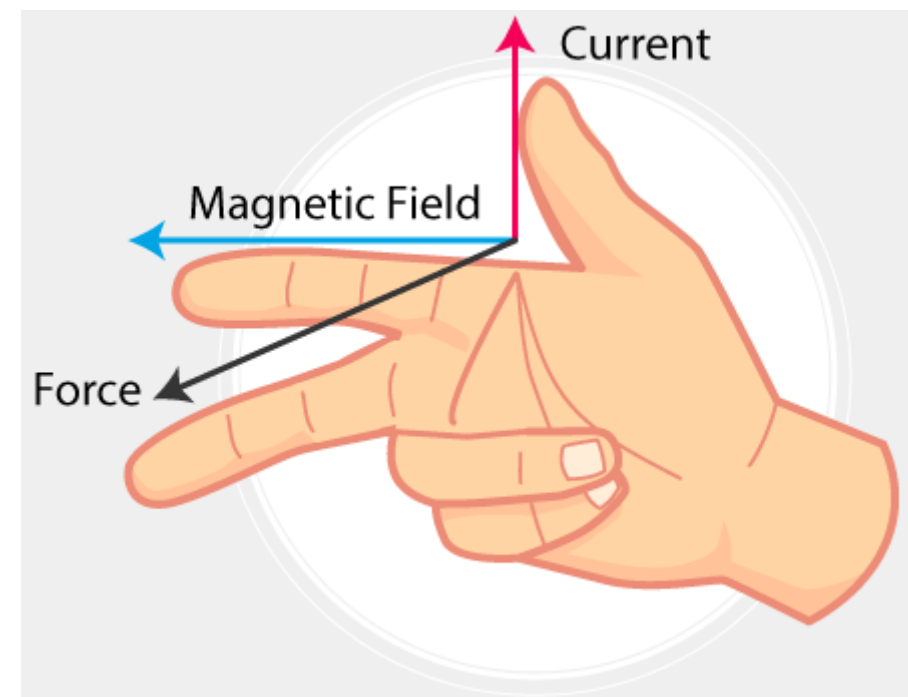
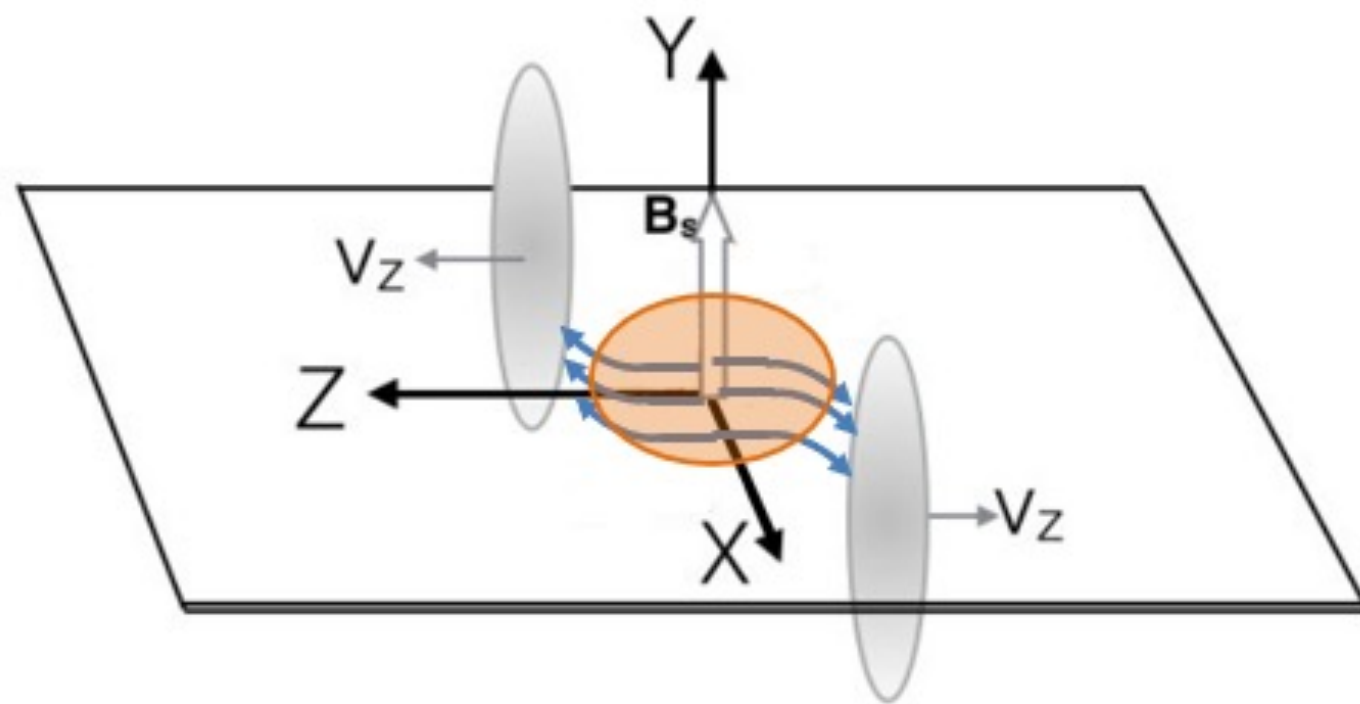
Electromagnetic field in heavy-ion collisions

- Complicated interplay of E and B fields gives rise to four types of currents
 - Faraday current:** induced current from the decrease of the B field in the medium



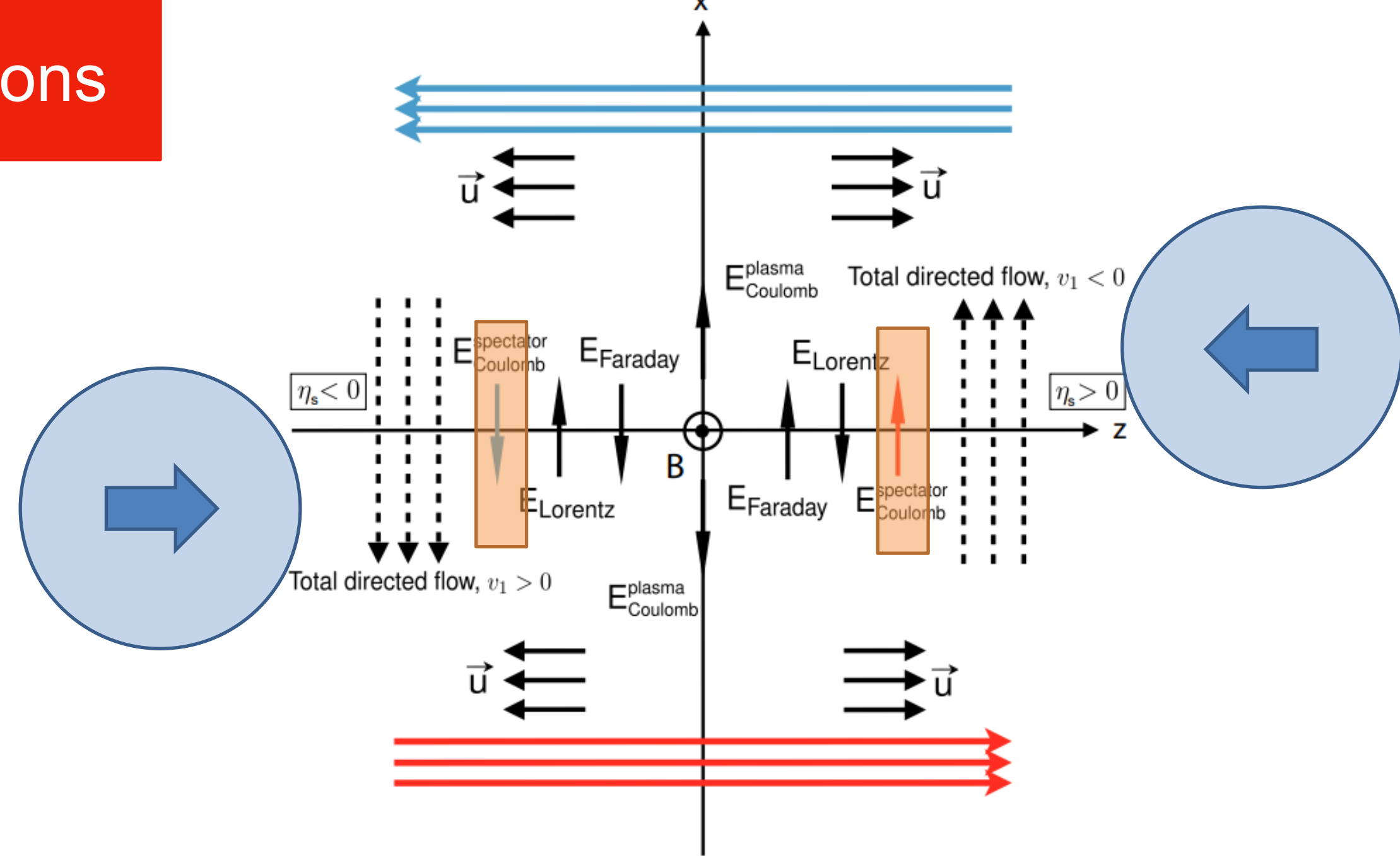
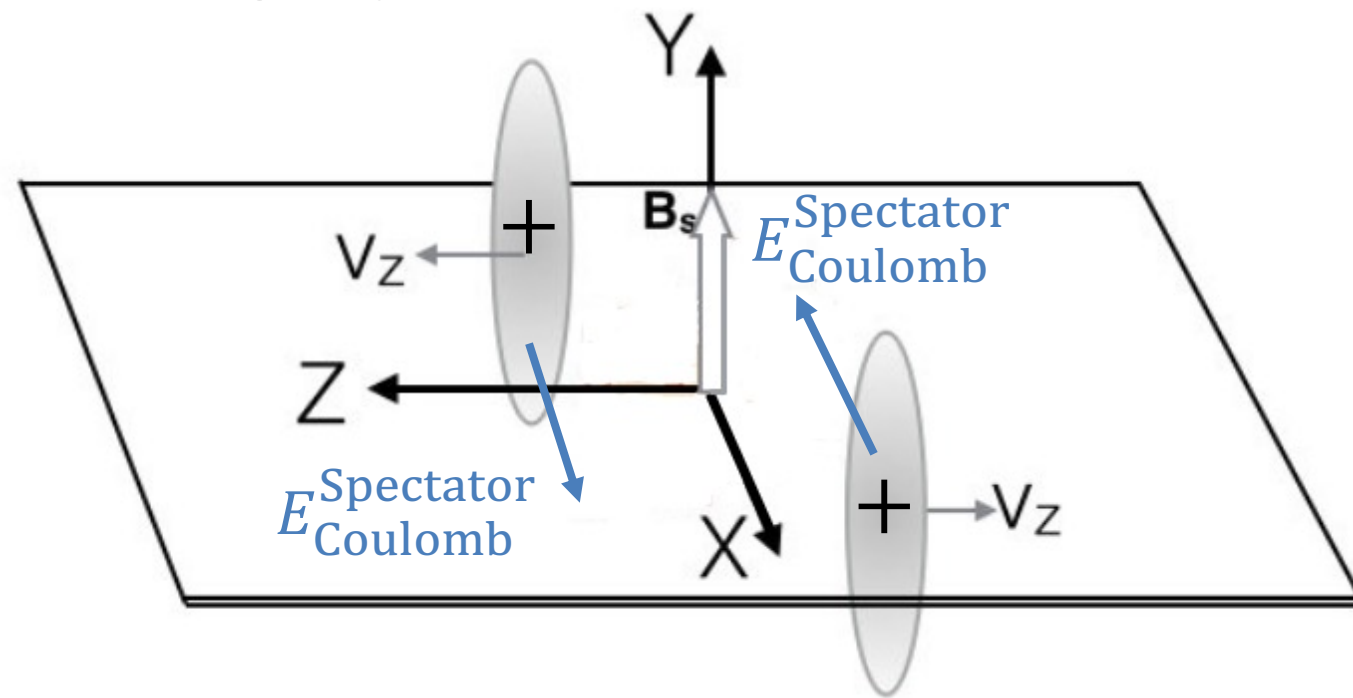
Electromagnetic field in heavy-ion collisions

- Complicated interplay of E and B fields gives rise to four types of currents
 - Faraday current:** induced current from the decrease of the B field in the medium
 - Lorentz current:** conducting medium expanding in longitudinal direction (\vec{u}) experiences Lorentz force result in an electric current perpendicular to \vec{u} and \vec{B} (Hall effect)



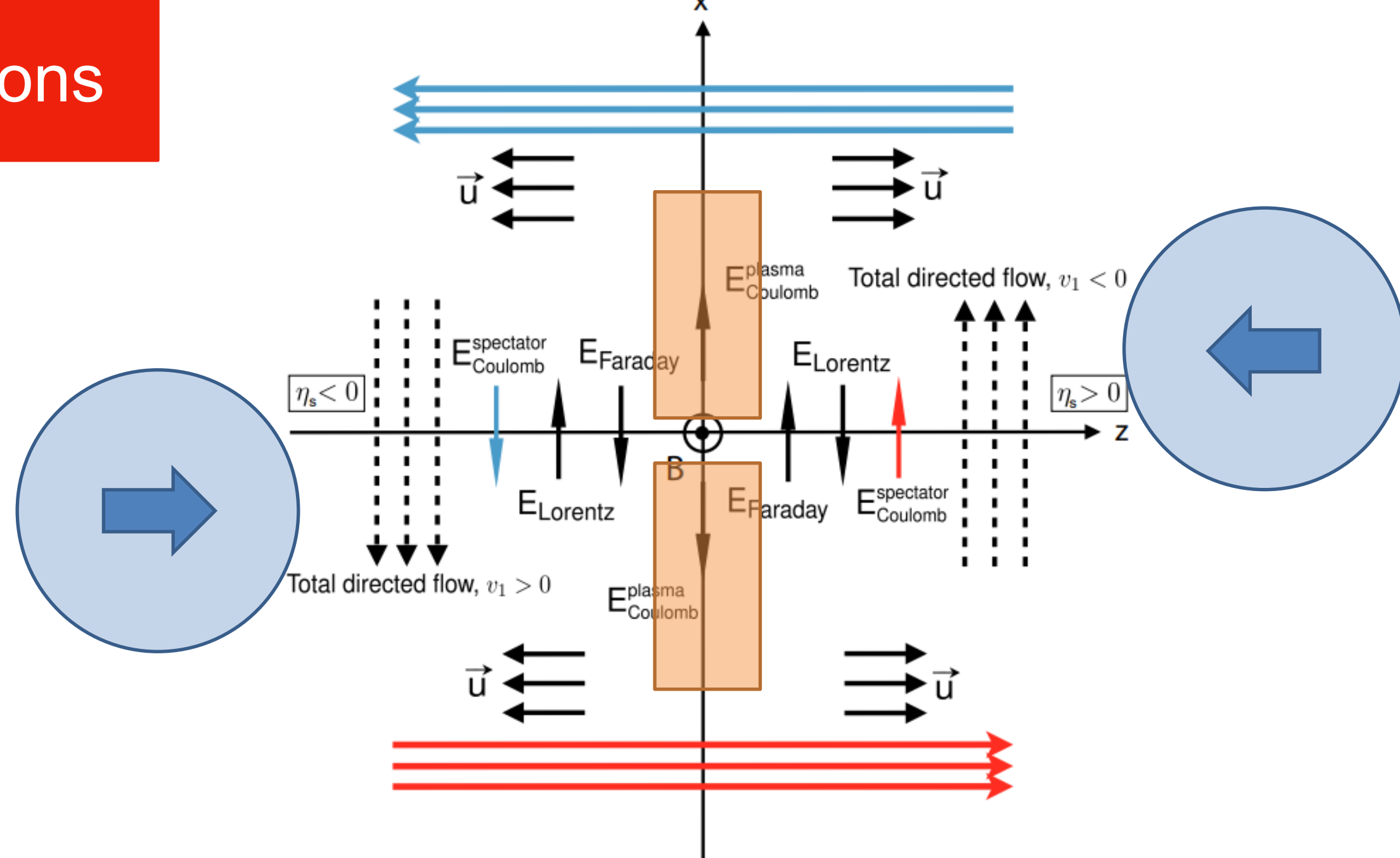
Electromagnetic field in heavy-ion collisions

- Complicated interplay of E and B fields gives rise to four types of currents
 - Faraday current:** induced current from the decrease of the B field in the medium
 - Lorentz current:** conducting medium expanding in longitudinal direction (\vec{u}) experiences Lorentz force result in an electric current perpendicular to \vec{u} and \vec{B} (Hall effect)
 - Coulomb current:** positively charged spectators passing the collision zone exerts an electric force on charged plasma



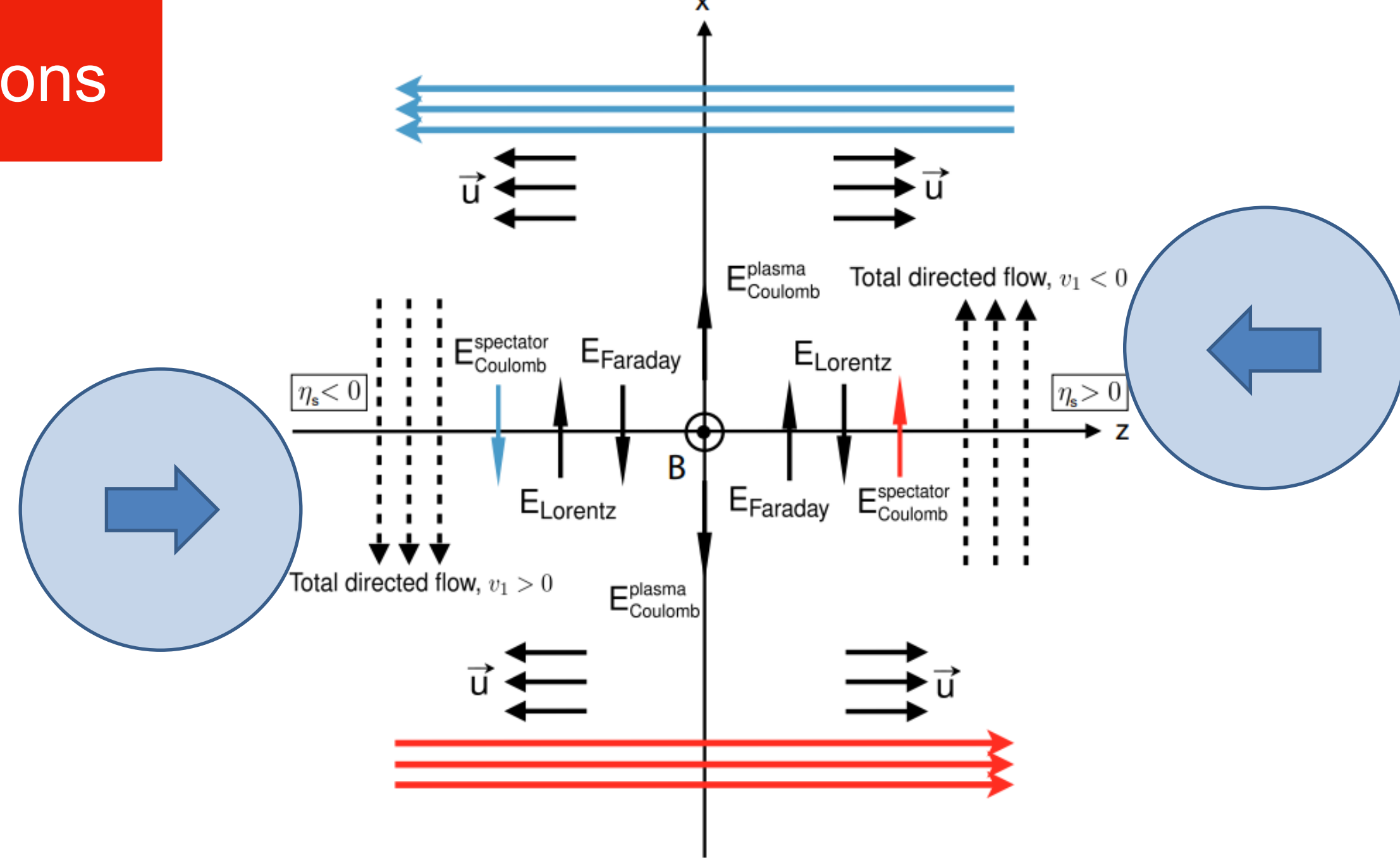
Electromagnetic field in heavy-ion collisions

- Complicated interplay of E and B fields gives rise to four types of currents
 - **Faraday current:** induced current from the decrease of the B field in the medium
 - **Lorentz current:** conducting medium expanding in longitudinal direction (\vec{u}) experiences Lorentz force result in an electric current perpendicular to \vec{u} and \vec{B} (Hall effect)
 - **Coulomb current:** positively charged spectators passing the collision zone exerts an electric force on charged plasma
 - **Plasma current:** the net positive charge in the plasma creates an outward E field

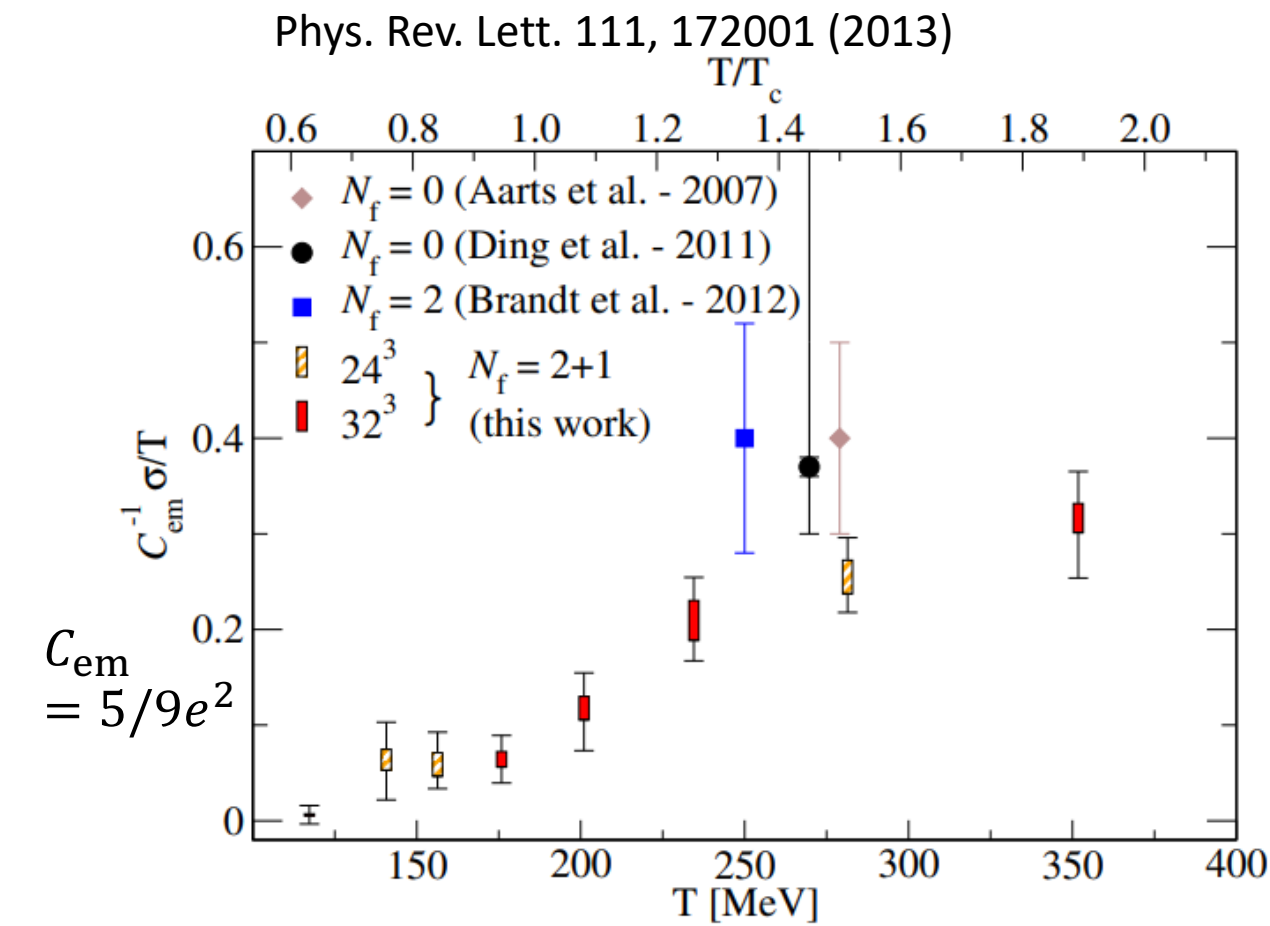


Electromagnetic field in heavy-ion collisions

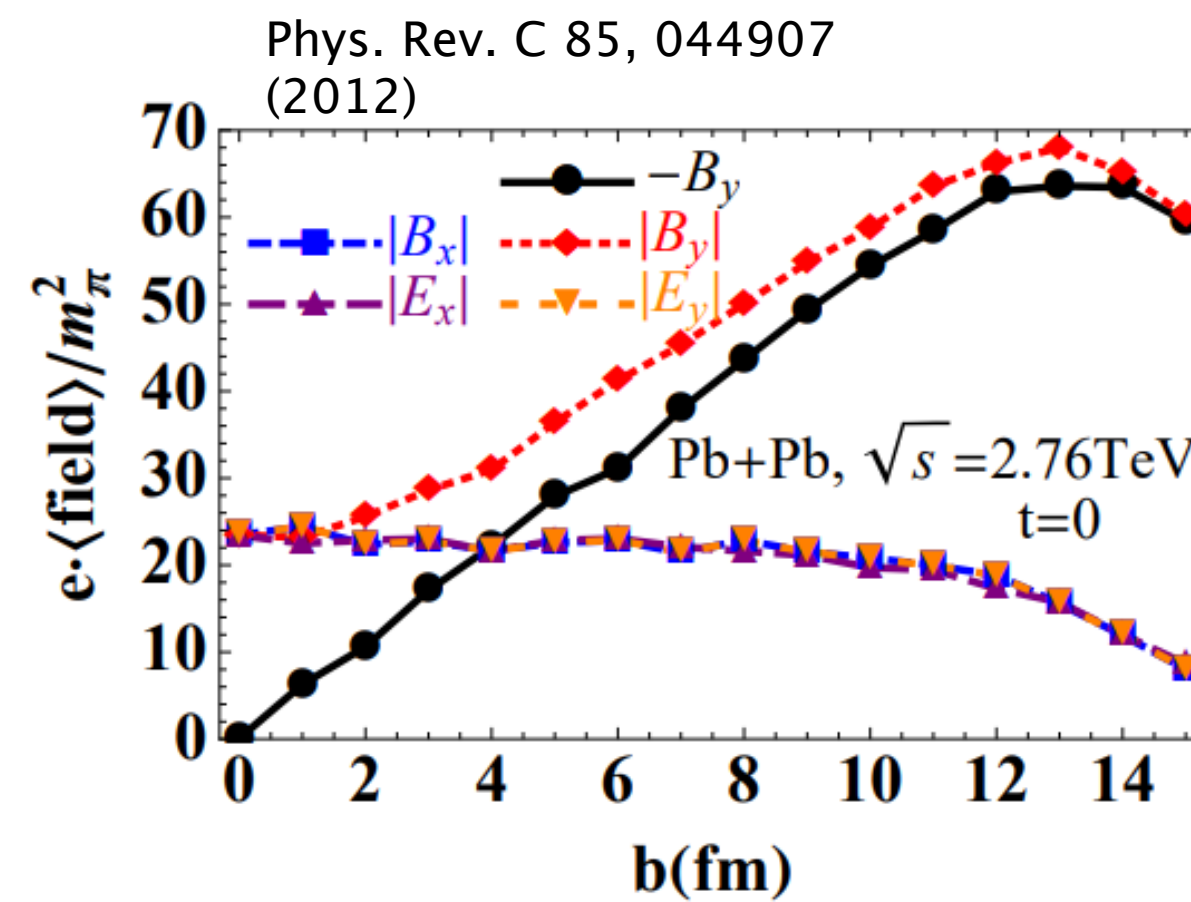
- Complicated interplay of E and B fields gives rise to four types of currents
 - **Faraday current:** induced current from the decrease of the B field in the medium
 - **Lorentz current:** conducting medium expanding in longitudinal direction (\vec{u}) experiences Lorentz force result in an electric current perpendicular to \vec{u} and \vec{B} (Hall effect)
 - **Coulomb current:** positively charged spectators passing the collision zone exerts an electric force on charged plasma
 - **Plasma current:** the net positive charge in the plasma creates an outward E field
- First three effects result in a charge-odd directed flow Δv_1
- Fourth effect results in a non-trivial radial and elliptic flow



Electromagnetic field in heavy-ion collisions

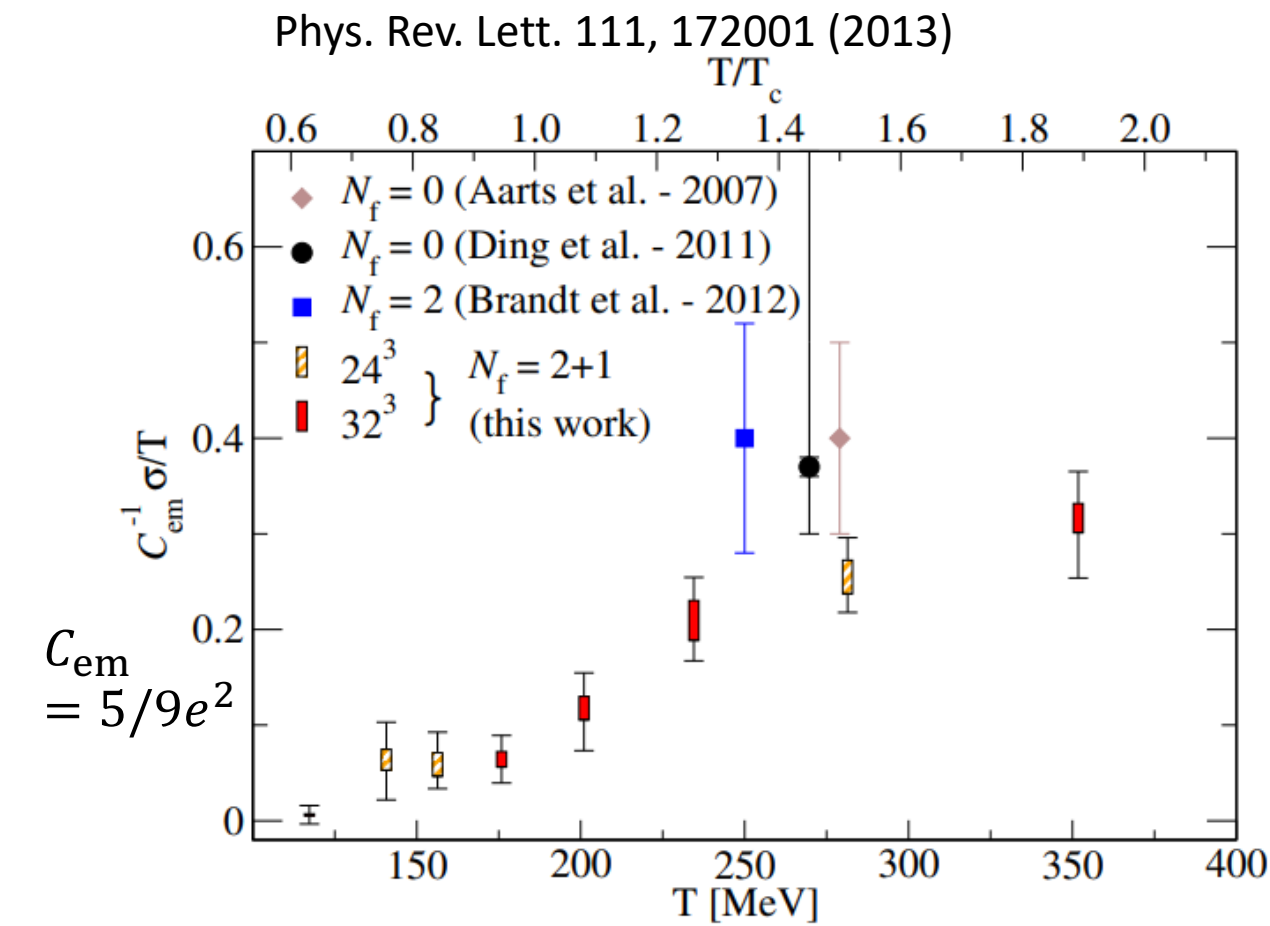


- Lattice QCD calculation suggests that the conductivity (σ) of QGP drops as the plasma cools.
- $\tau_{\text{decay}} \propto \sigma$ (τ_{decay} time constant of the magnetic field)

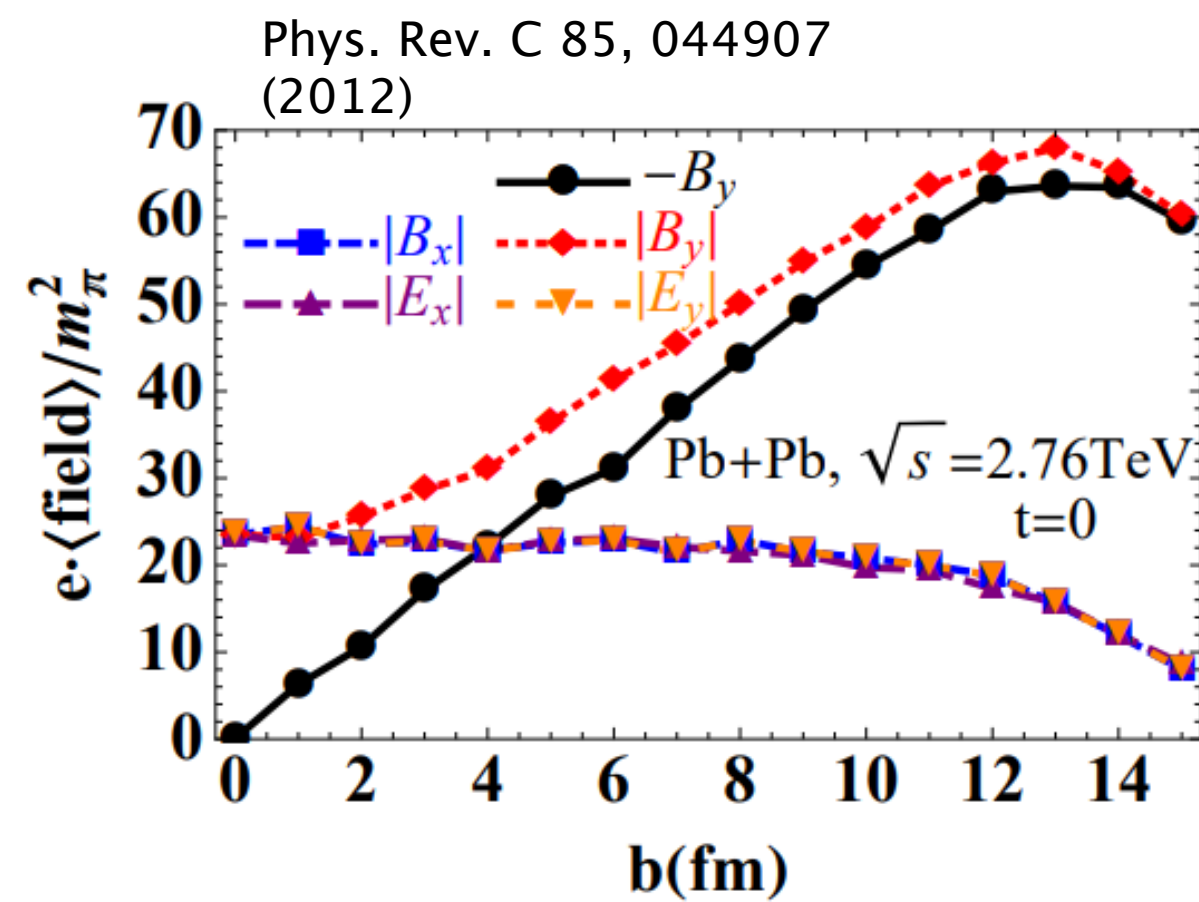


- EM field at $t = 0$ and $r = 0$, the centre of the collision region. B_y increases for more peripheral collisions

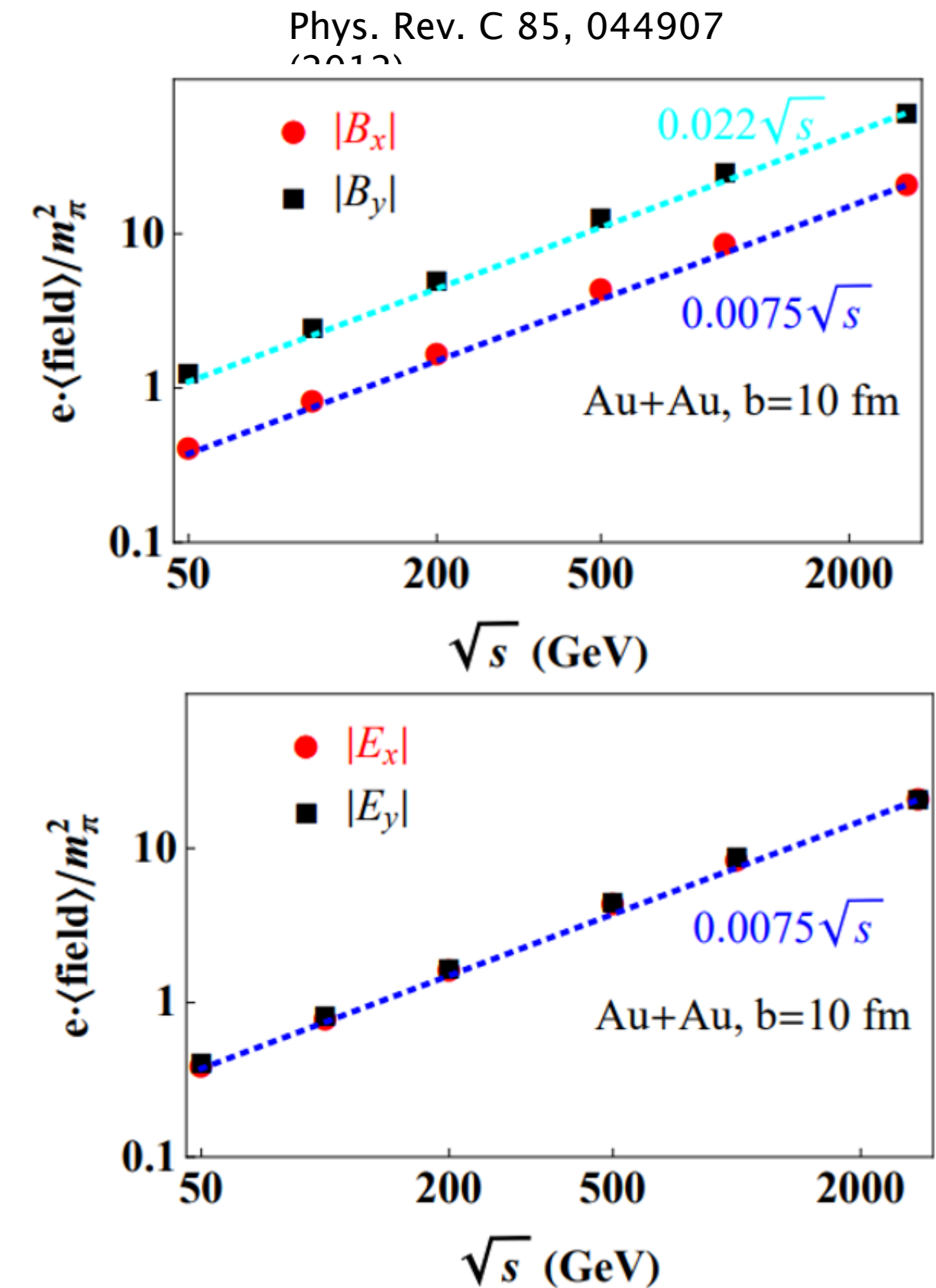
Electromagnetic field in heavy-ion collisions



- Lattice QCD calculation suggests that the conductivity (σ) of QGP drops as the plasma cools.
- $\tau_{\text{decay}} \propto \sigma$ (τ_{decay} time constant of the magnetic field)

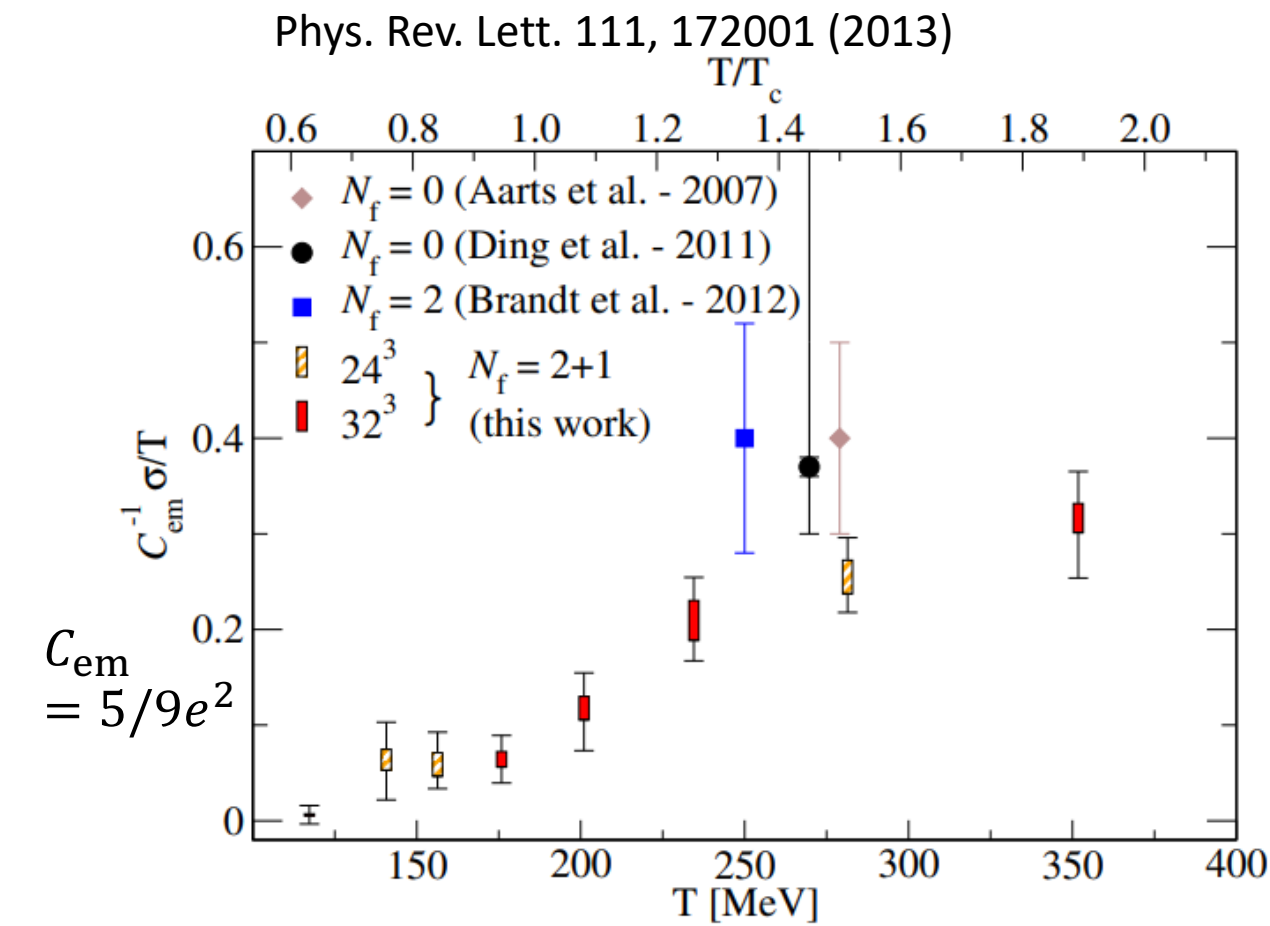


- EM field at $t = 0$ and $r = 0$, the centre of the collision region. B_y increases for more peripheral collisions

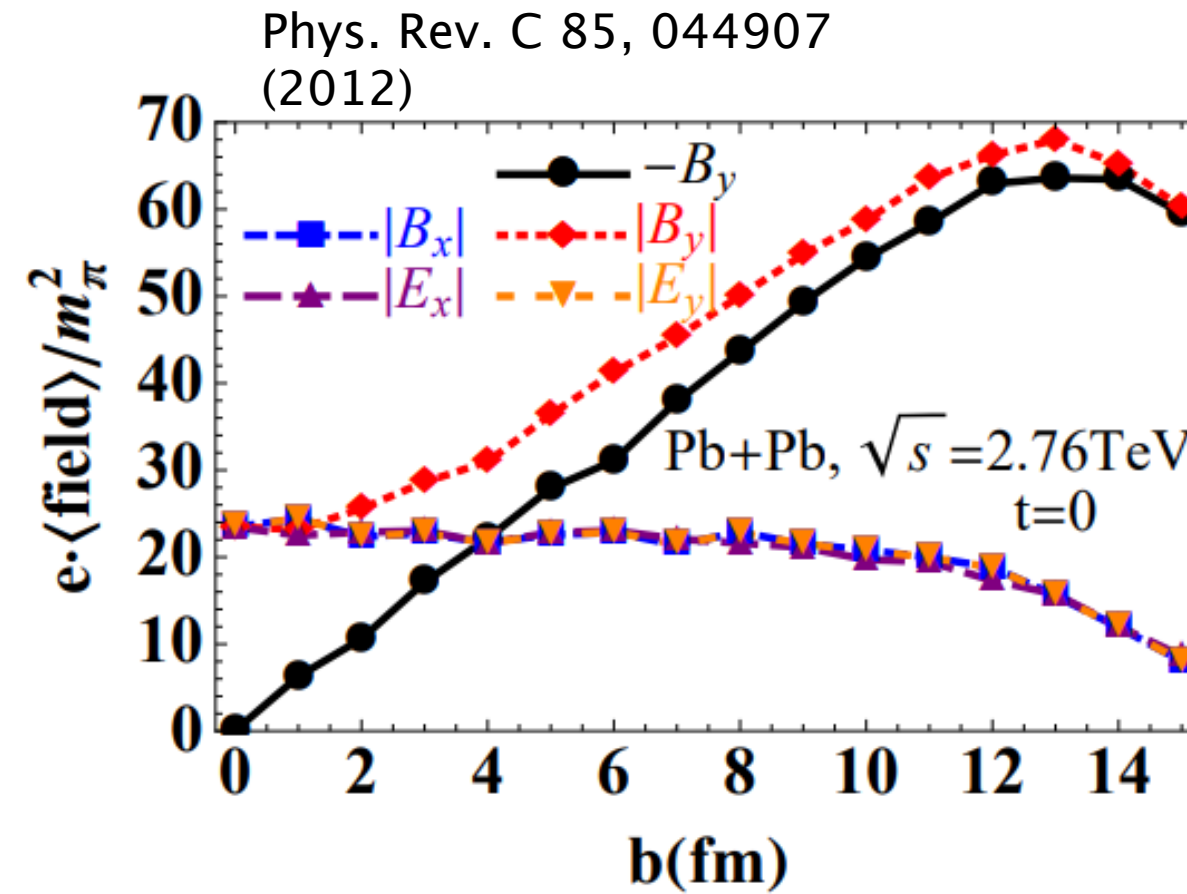


- Higher \sqrt{s} leads to higher E and B field

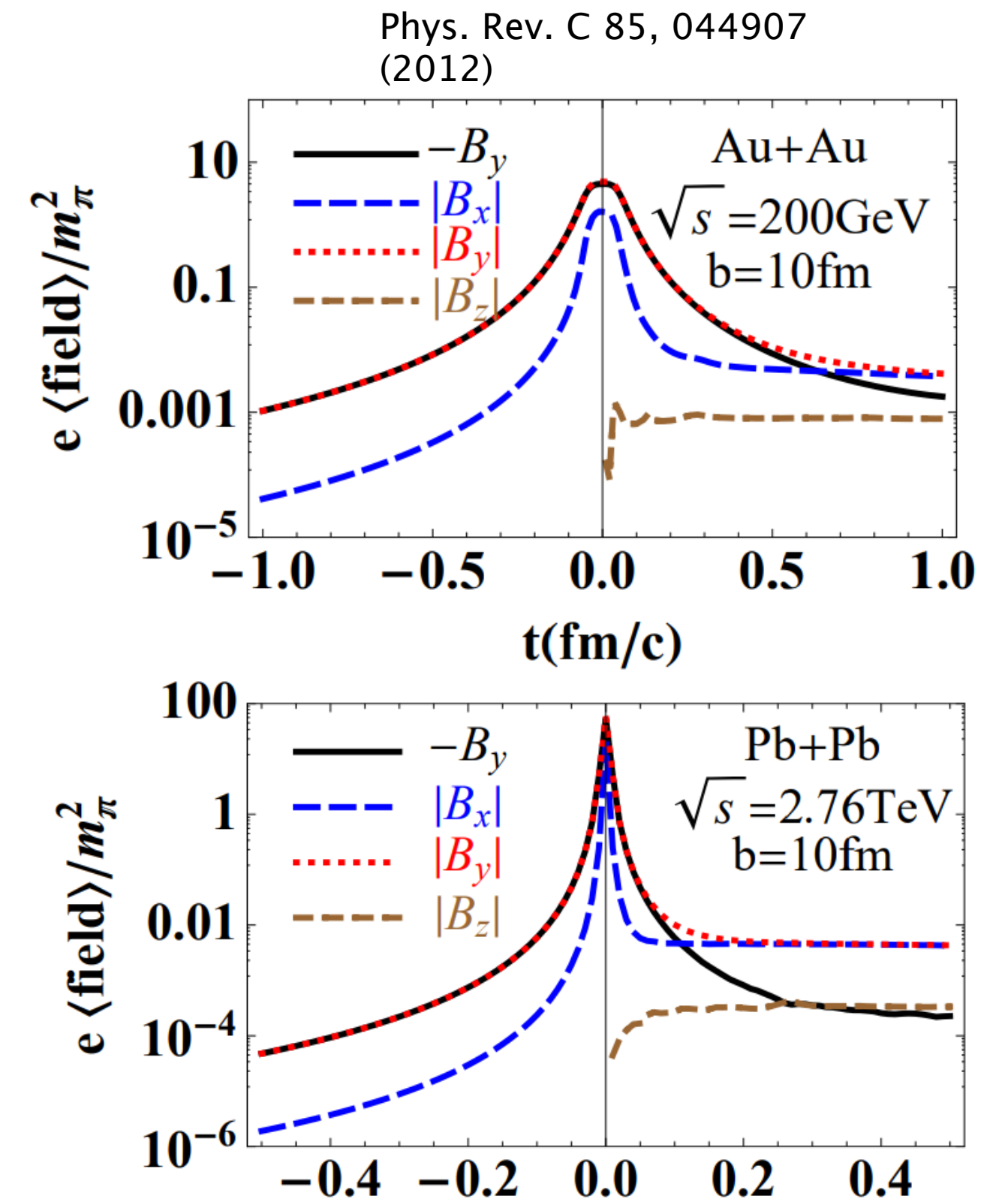
Electromagnetic field in heavy-ion collisions



- Lattice QCD calculation suggests that the conductivity (σ) of QGP drops as the plasma cools.
- $\tau_{\text{decay}} \propto \sigma$ (τ_{decay} time constant of the magnetic field)



- EM field at $t = 0$ and $r = 0$, the centre of the collision region. B_y increases for more peripheral collisions



- Lifetime of the B field is shorter at 2.76 TeV than at 200 GeV.
- Lifetime of QGP at ALICE is $\sim 10 \text{ fm}/c^{[1]}$

[1] Phys. Lett. B 696 (2011) 328

Charge-dependent directed flow

- Produced particle azimuthal distribution relative to the collision-spectator plane is directly sensitive to the presence of the electromagnetic fields
- The particle (α) azimuthal distribution:

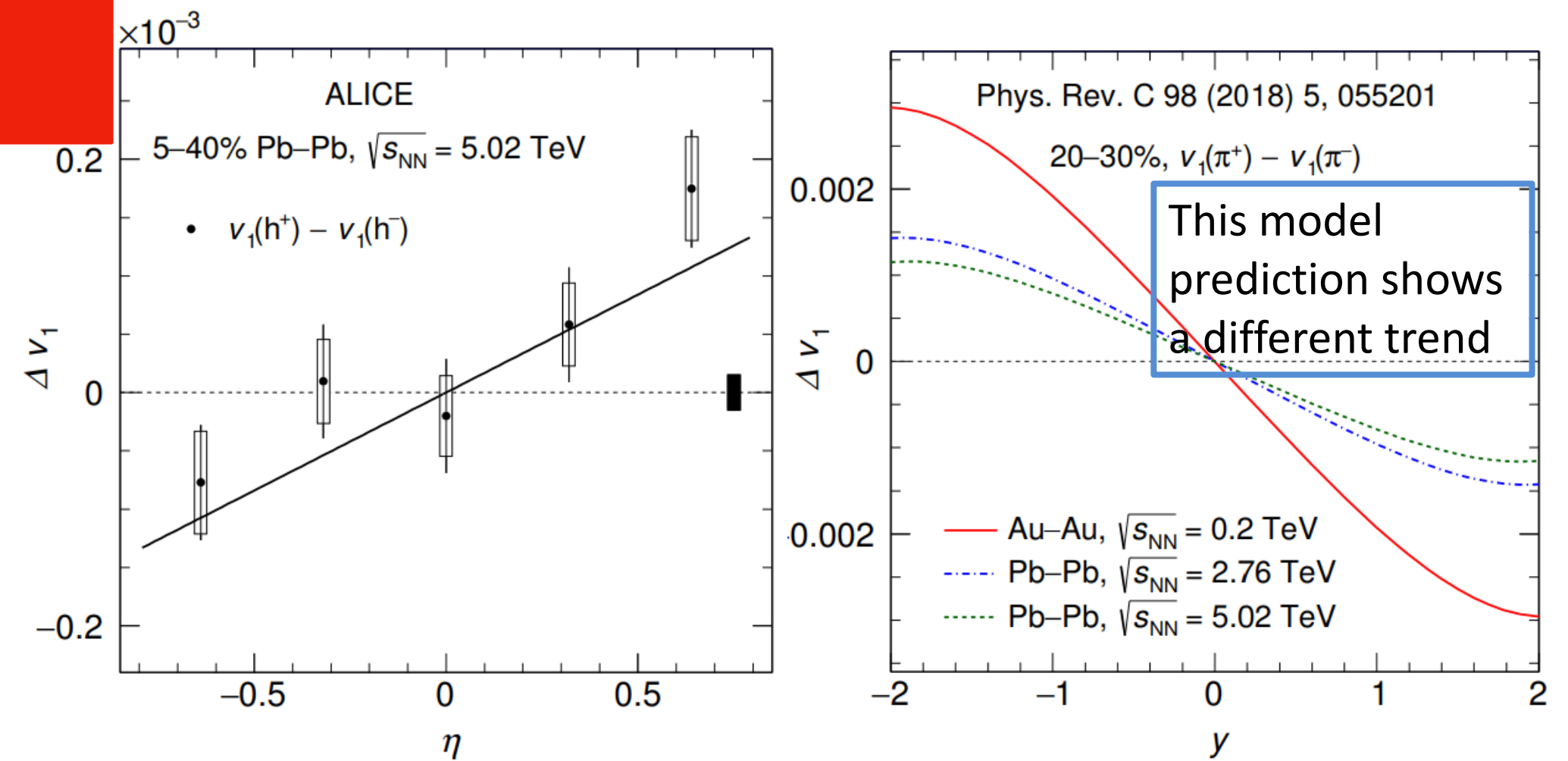
$$\frac{dN}{d\varphi_\alpha} \sim 1 + 2 \sum_n v_{n,\alpha} \cos[n(\varphi_\alpha - \Psi_{RP})]$$

- ALICE uses the scalar-product ^[1] method

$$v_1 = \frac{\langle u_x Q_x^{A,C} + u_y Q_y^{A,C} \rangle}{\sqrt{|\langle Q_x^A Q_x^C + Q_y^A Q_y^C \rangle|}}$$

where $\mathbf{u} = (\cos \varphi, \sin \varphi)$ and $\mathbf{Q}^{A,C}$ is the flow vector for ZDC at A and C side.

- $\Delta v_1 = v_1(h^+) - v_1(h^-)$
- ALICE inclusive hadron measurements



[1] Phys.Rev. C87 (2013) 044907

Charge-dependent directed flow

- Produced particle azimuthal distribution relative to the collision-spectator plane is directly sensitive to the presence of the electromagnetic fields

- The particle (α) azimuthal distribution:

$$\frac{dN}{d\varphi_\alpha} \sim 1 + 2 \sum_n v_{n,\alpha} \cos[n(\varphi_\alpha - \Psi_{RP})]$$

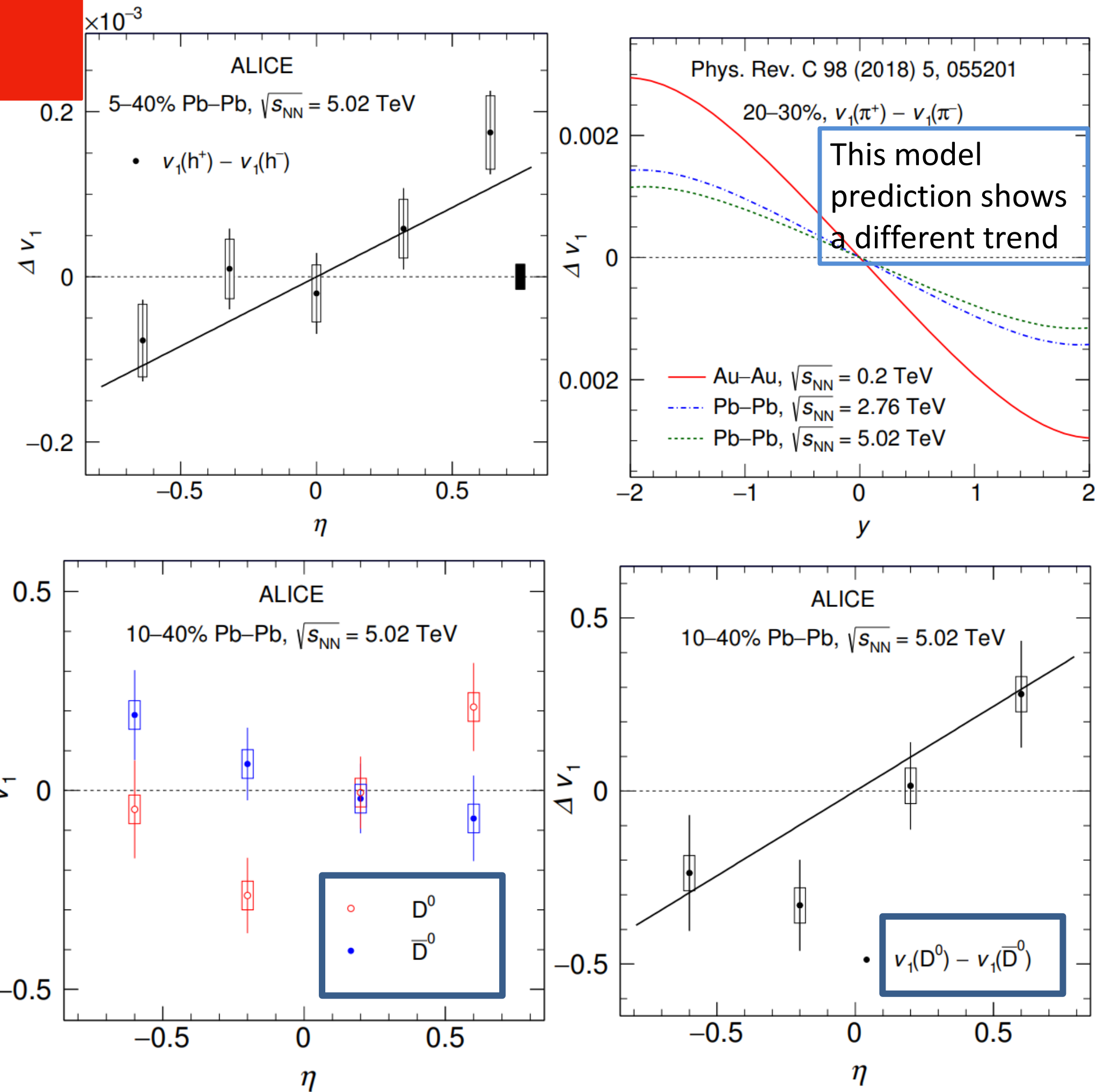
- ALICE uses the scalar-product ^[1] method

$$v_1 = \frac{\langle u_x Q_x^{A,C} + u_y Q_y^{A,C} \rangle}{\sqrt{|\langle Q_x^A Q_x^C + Q_y^A Q_y^C \rangle|}}$$

where $\mathbf{u} = (\cos \varphi, \sin \varphi)$ and $\mathbf{Q}^{A,C}$ is the flow vector for ZDC at A and C side.

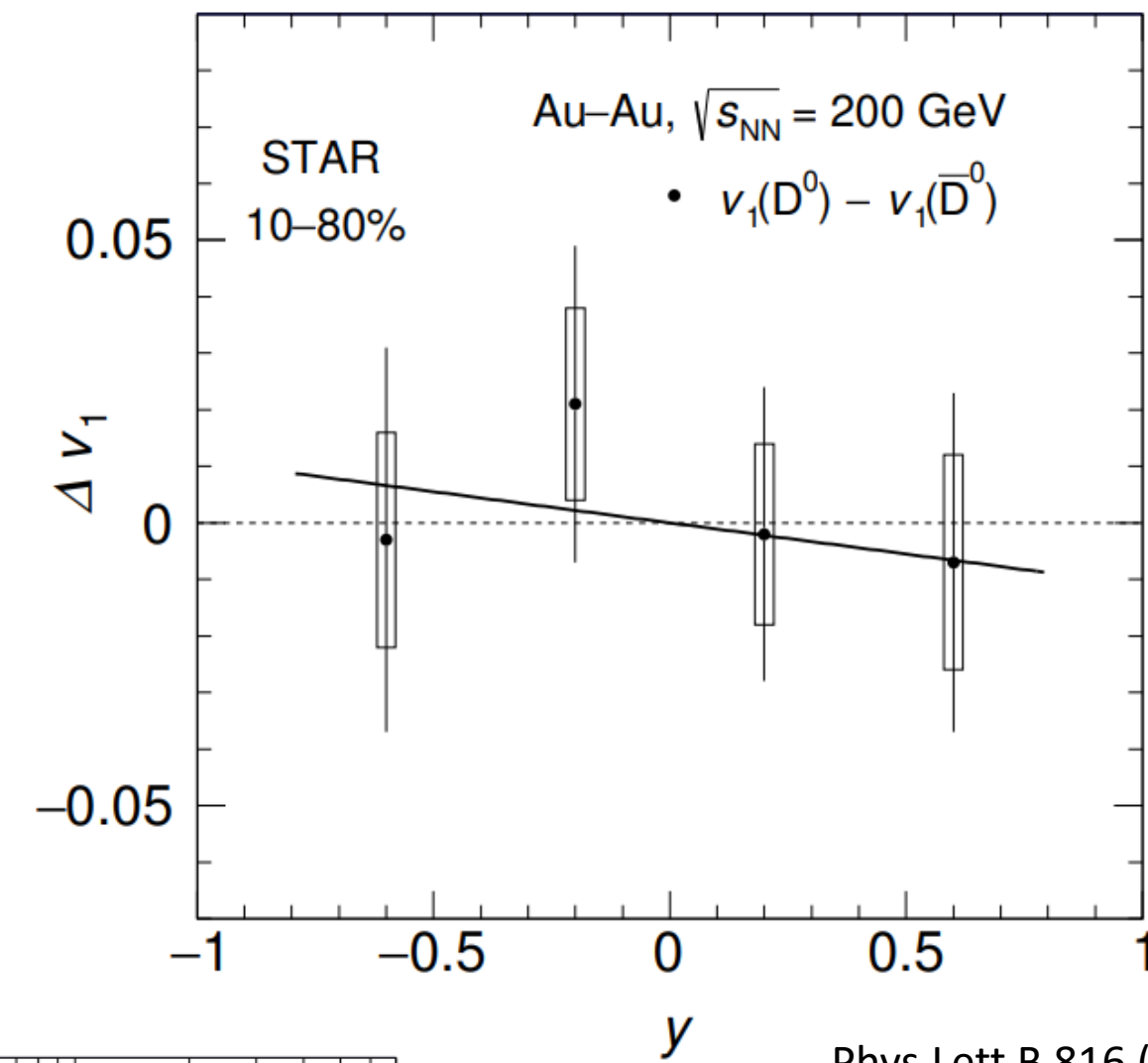
- $\Delta v_1 = v_1(h^+) - v_1(h^-)$
- ALICE inclusive hadron measurements
- ALICE Heavy-flavor measurements on D^0 ($c\bar{u}$) and \bar{D}^0 ($\bar{c}u$)
 - Charm quarks produced early in the hard process experience more effect from EM field leading to a larger gradient ($0.49 \pm 0.17(\text{stats}) \pm 0.06(\text{syst.})$)

[1] Phys.Rev. C87 (2013) 044907

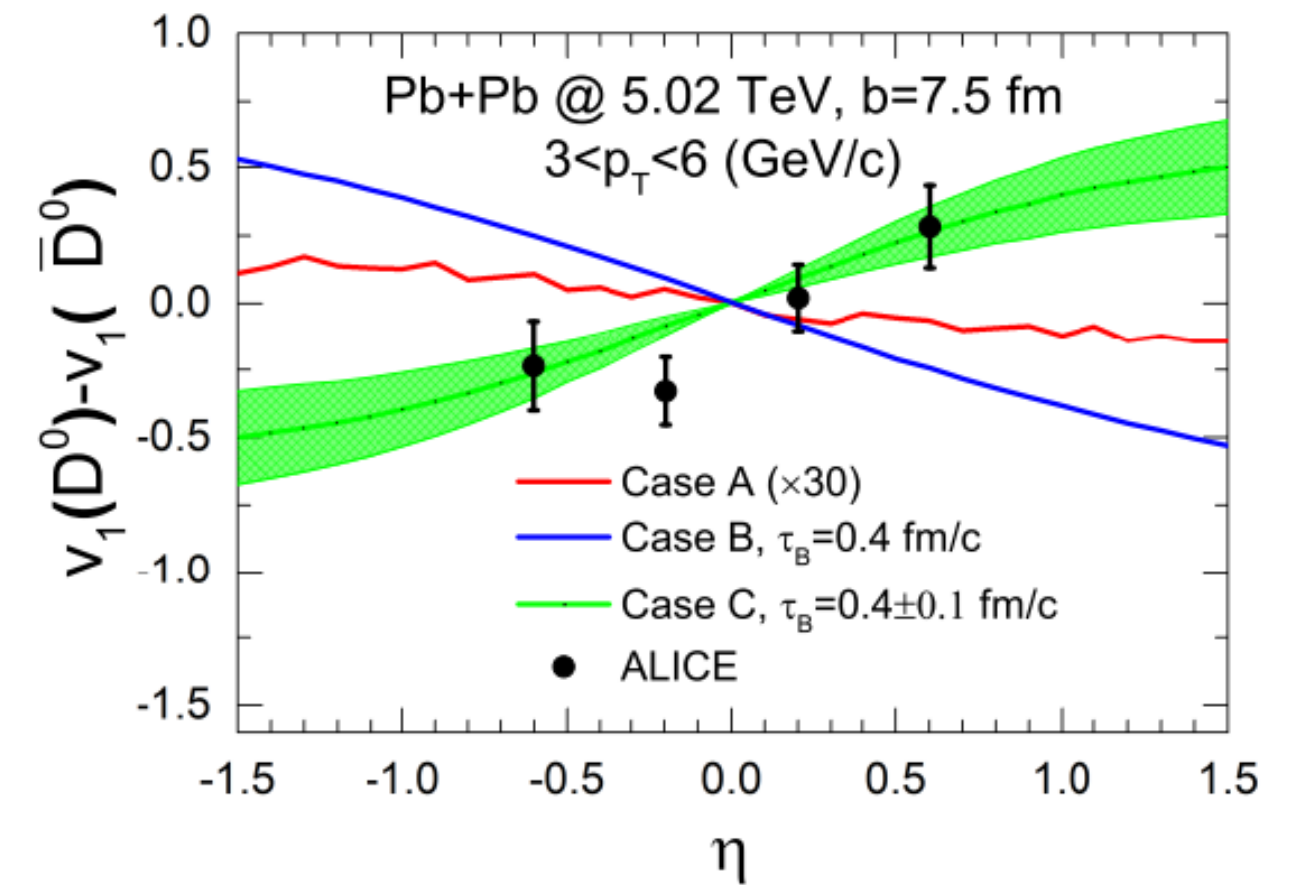
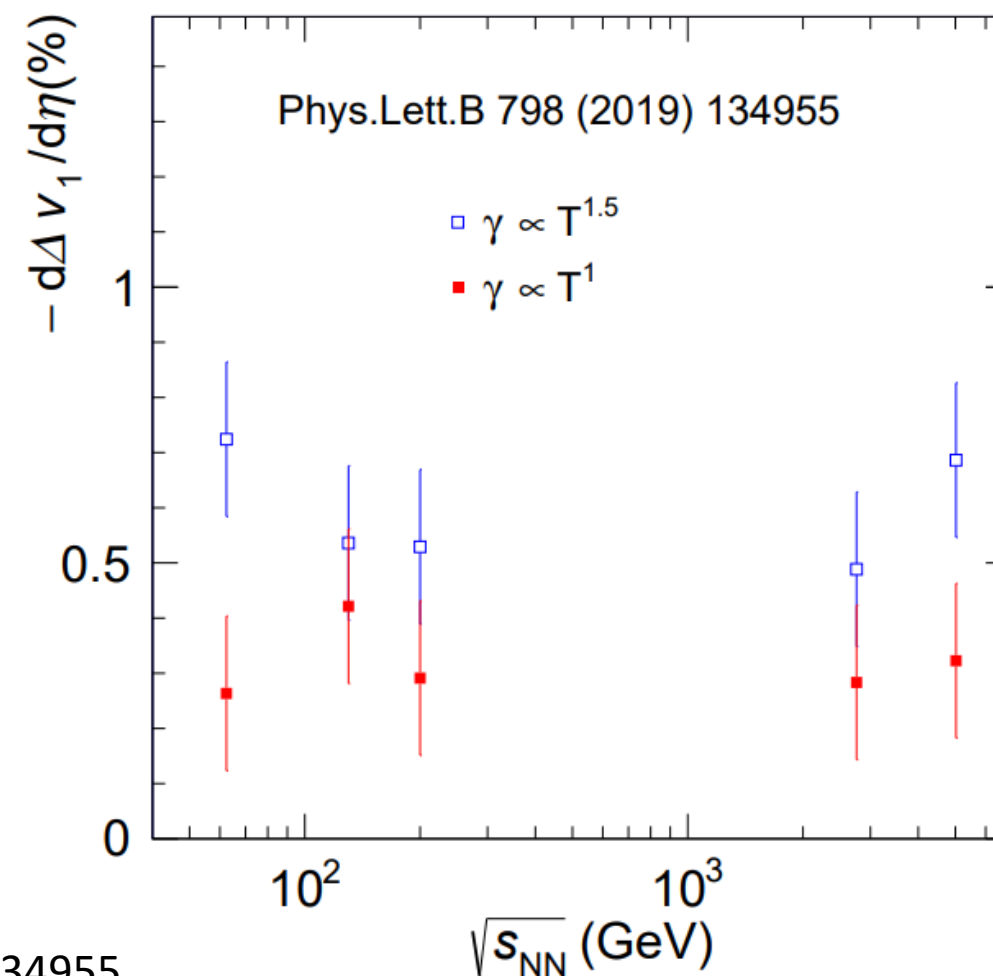


Charge-dependent directed flow

- Measurement of Δv_1 at RHIC energy shows a slope of $-0.011 \pm 0.034(\text{stat.}) \pm 0.020(\text{syst.})$
- The larger slope of Δv_1 of D mesons at LHC might indicate a stronger effect of the Lorentz force relative to the Coulomb one
- Predictions for the dependence of the Δv_1 for the D meson is almost flat for two choice of drag $\gamma^{[1]}$ from the matter
- Recent model calculation using a magnetic field with a slower time evolution and with a lifetime of about 0.4 fm/c can reproduce the measurement



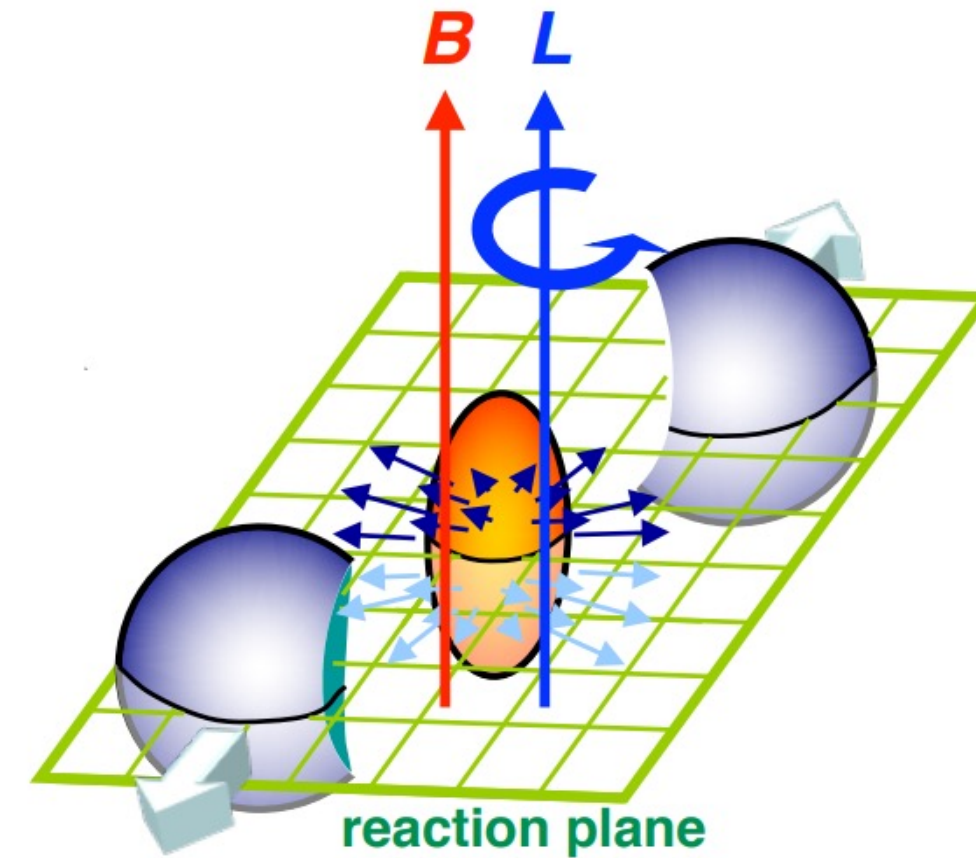
Phys.Lett B 816 (2021) 136271



[1] More discussion on drag vs. Lorentz force, see Phys.Lett.B 798 (2019) 134955

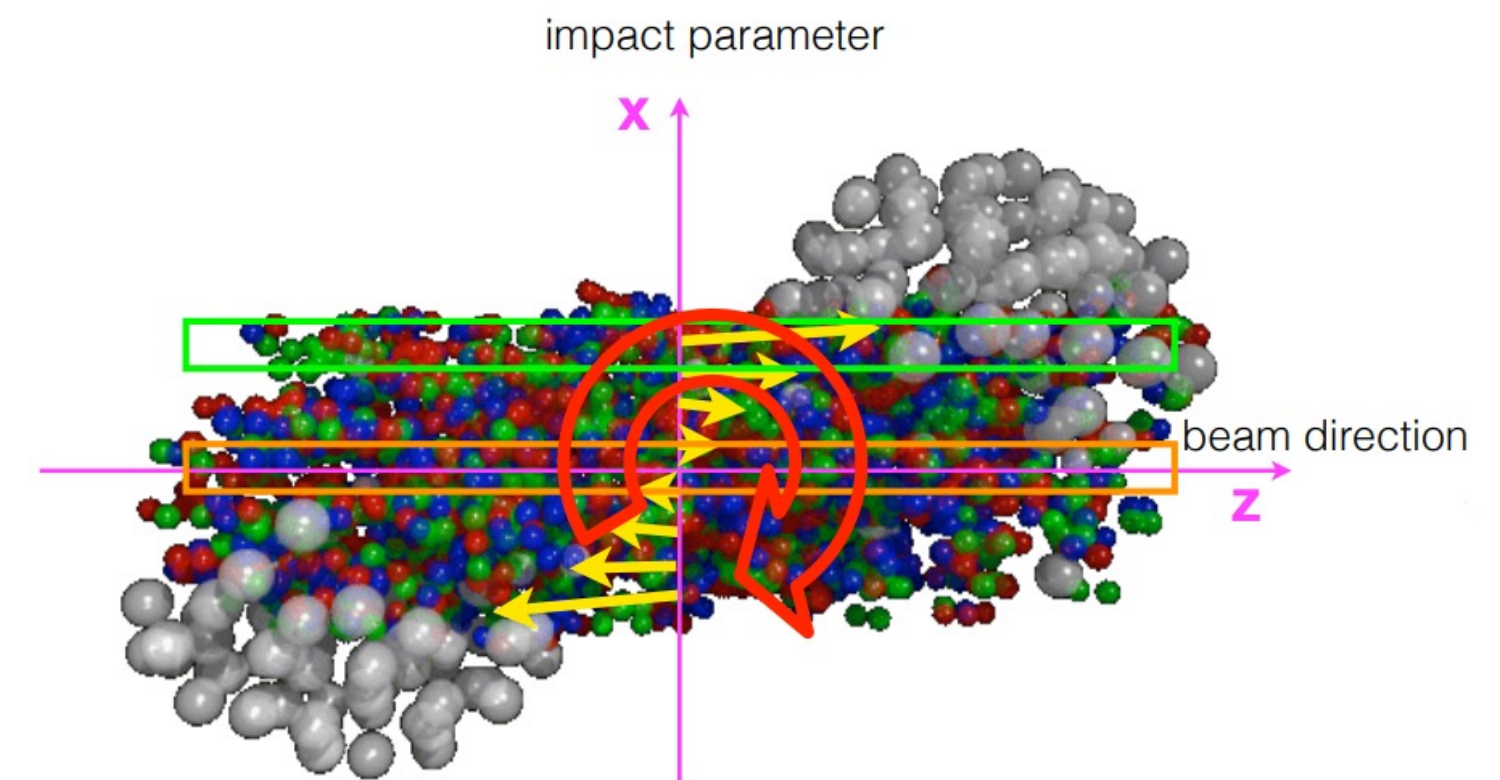
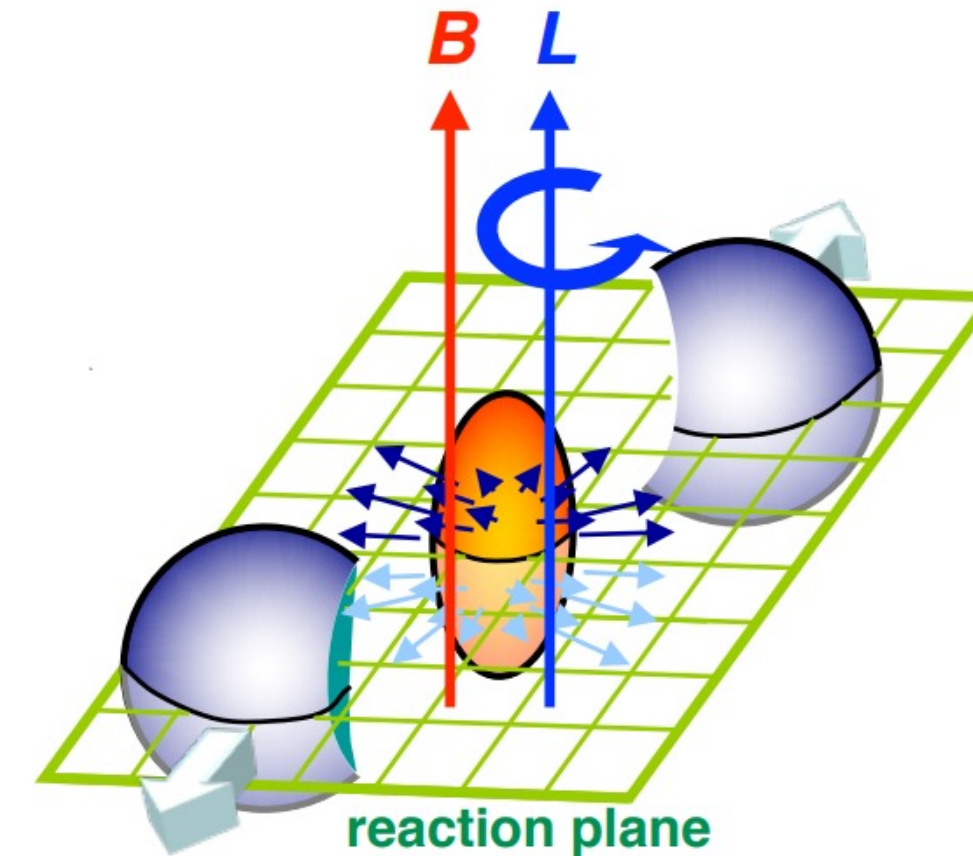
Hyperon polarization

- Global polarization: spin alignment along the initial angular momentum
- Two contributions to the global polarization:
 - Magnetic field align particles' and antiparticles' spin oppositely due to the opposite sign of magnetic moment



Hyperon polarization

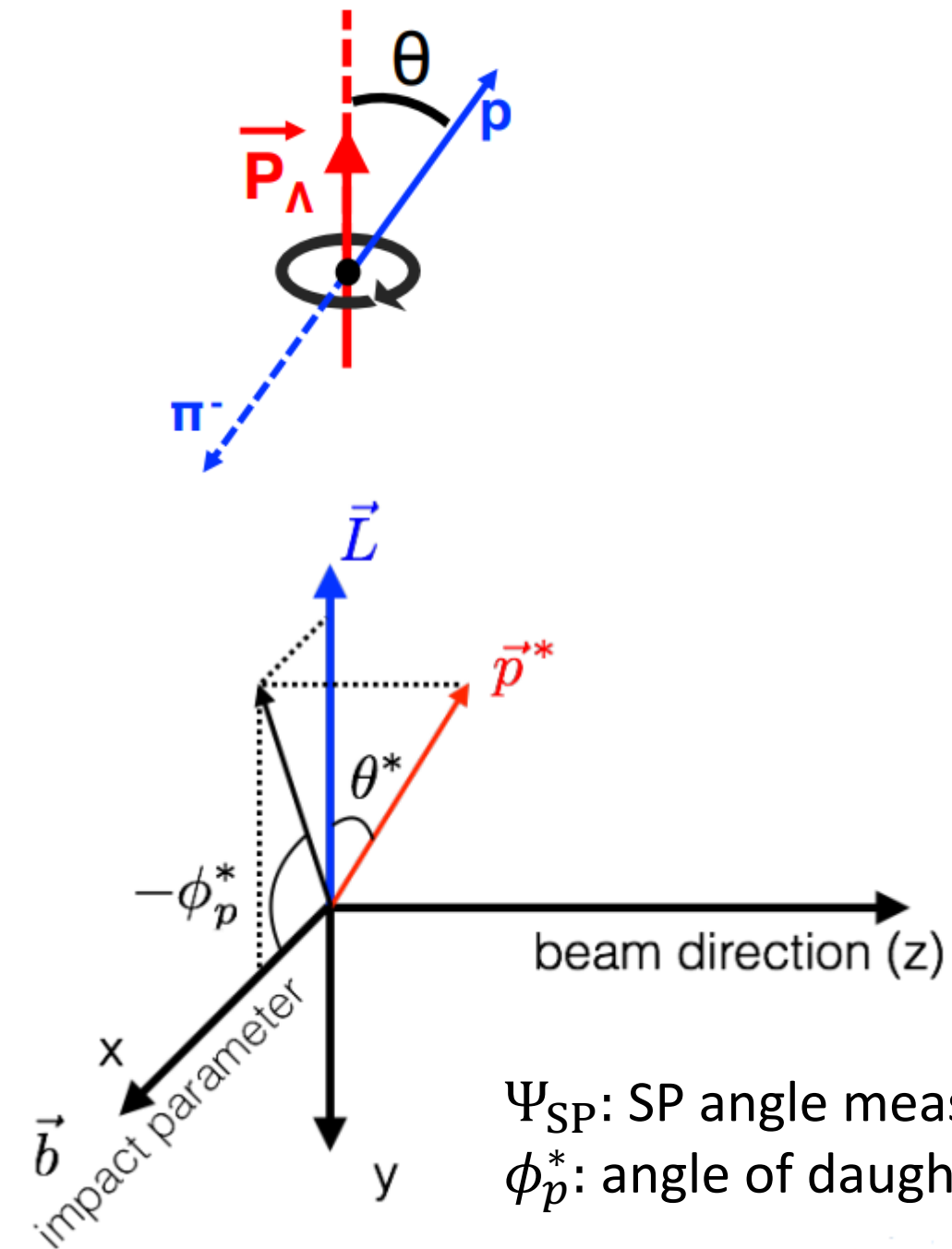
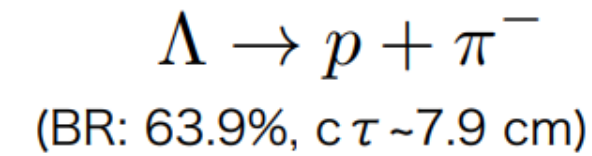
- Global polarization: spin alignment along the initial angular momentum
- Two contributions to the global polarization:
 - Magnetic field align particles' and antiparticles' spin oppositely due to the opposite sign of magnetic moment
 - Large vorticity expected due to velocity difference along beam direction. Orbital angular momentum is transferred to particle spin equally for particles and antiparticles



Hyperon polarization

- Global polarization: spin alignment along the initial angular momentum
- Two contributions to the global polarization:
 - Magnetic field align particles' and antiparticles' spin oppositely due to the opposite sign of magnetic moment
 - Large vorticity expected due to velocity difference along beam direction. Orbital angular momentum is transferred to particle spin equally for particles and antiparticles
- How to measure?
 - Λ ($\bar{\Lambda}$) decays dominantly into $p + \pi$. Protons are emitted preferentially in the direction of the Λ ($\bar{\Lambda}$) spin
- The global polarization is given by

$$P_H = -\frac{8}{\pi\alpha_H} \frac{\langle \sin(\phi_p^* - \Psi_{SP}) \rangle}{R_{SP}^{(1)}}$$



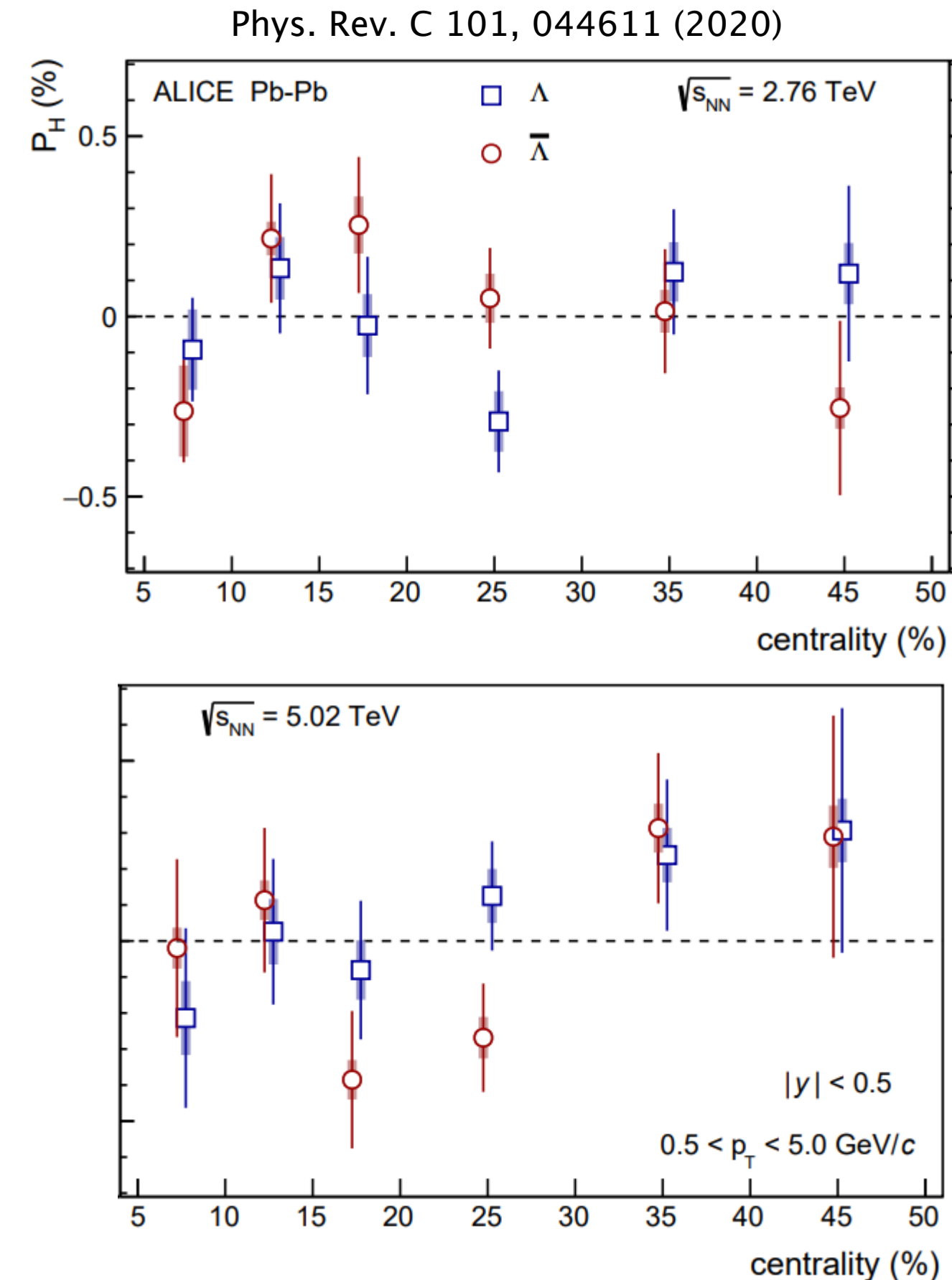
Ψ_{SP} : SP angle measured by ZDC
 ϕ_p^* : angle of daughter proton in Λ rest frame

arXiv:1710.08934v1

Hyperon polarization

- Global polarization: spin alignment along the initial angular momentum
- Two contributions to the global polarization:
 - Magnetic field align particles' and antiparticles' spin oppositely due to the opposite sign of magnetic moment
 - Large vorticity expected due to velocity difference along beam direction. Orbital angular momentum is transferred to particle spin equally for particles and antiparticles
- How to measure?
 - Λ ($\bar{\Lambda}$) decays dominantly into $p + \pi$. Protons are emitted preferentially in the direction of the Λ ($\bar{\Lambda}$) spin
- The global polarization is given by

$$P_H = -\frac{8}{\pi\alpha_H} \frac{\langle \sin(\phi_p^* - \Psi_{SP}) \rangle}{R_{SP}^{(1)}}$$



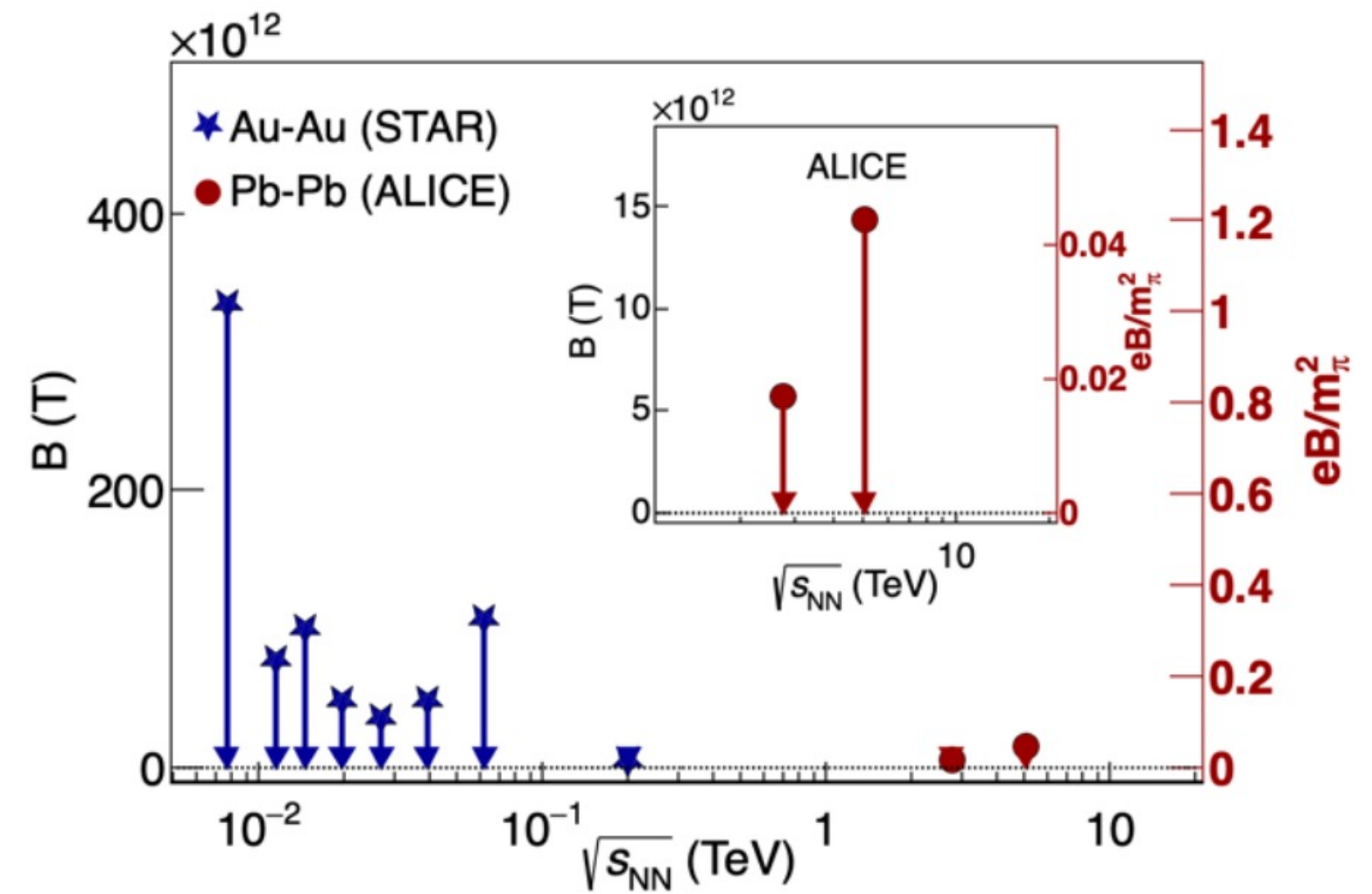
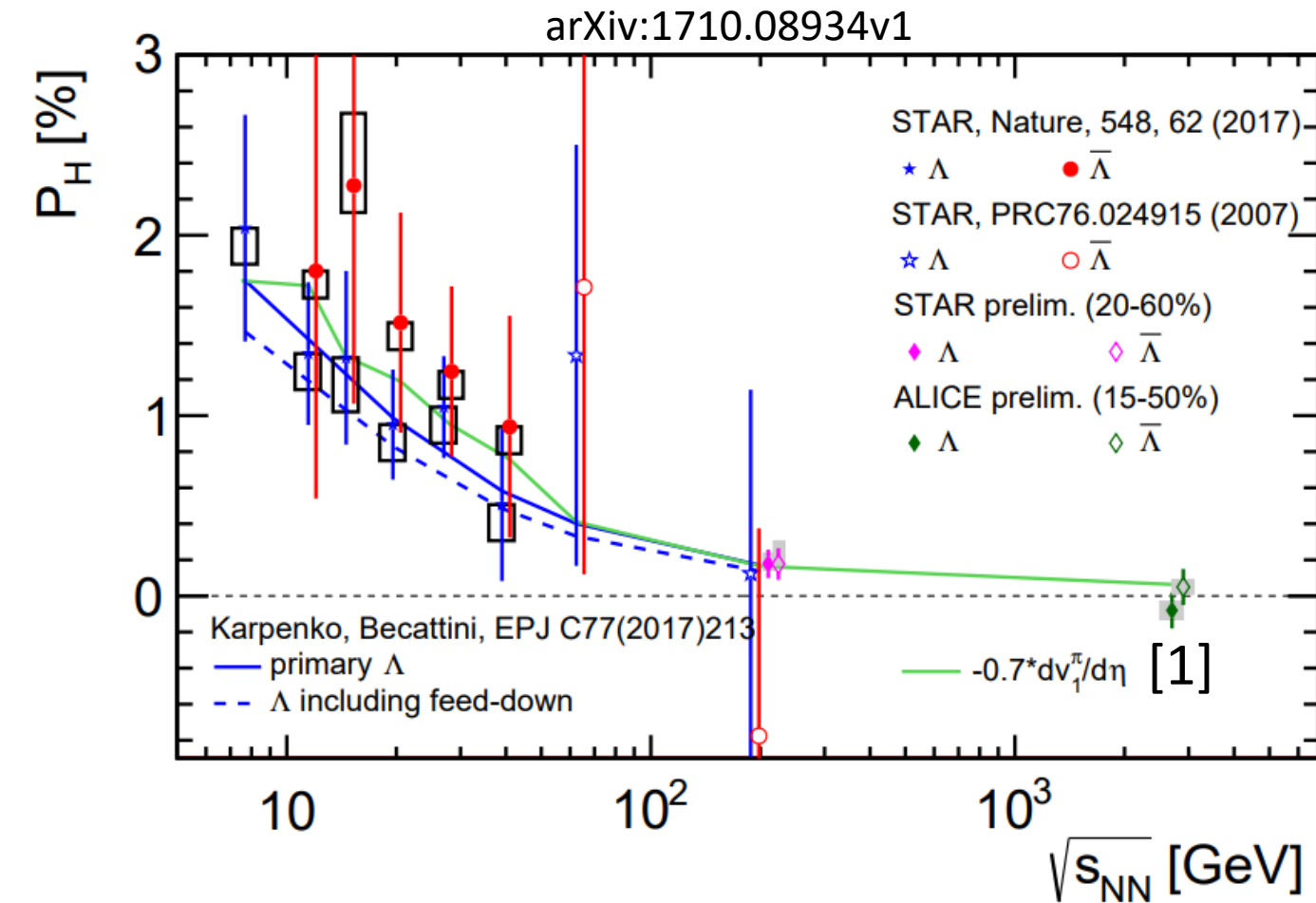
Hyperon polarization

- Confirms the observed earlier trend of the global polarization decrease with increasing $\sqrt{s_{NN}}$
- The magnitude of B field at kinetic freeze-out can be probed by $\Lambda - \bar{\Lambda}$ splitting:

$$P_{\Lambda} \approx \frac{0.5\omega}{T} \left[+ \frac{|\mu_{\Lambda}|B}{T} \right], P_{\bar{\Lambda}} \approx \frac{0.5\omega}{T} \left[- \frac{|\mu_{\Lambda}|B}{T} \right]$$

$$\Rightarrow P_{\Lambda} - P_{\bar{\Lambda}} \approx \frac{2|\mu_{\Lambda}|B}{T}$$

- More data (100x more) in the high luminosity LHC run will help to test the predicted splitting of P_{Λ} and $P_{\bar{\Lambda}}$ better



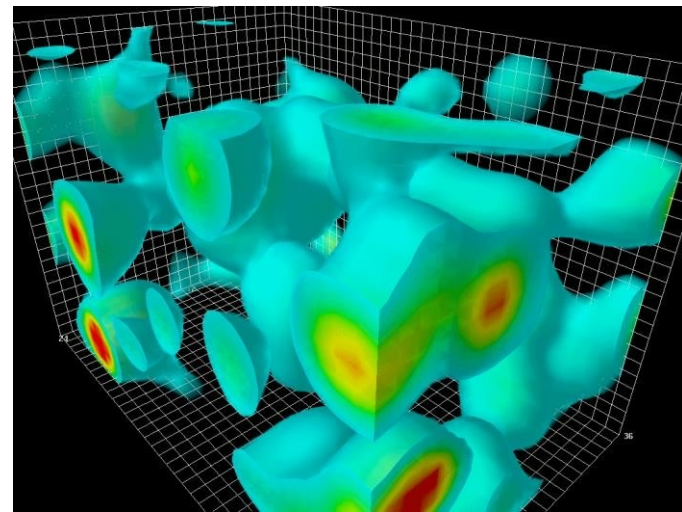
[1] Advances in High Energy Physics, vol. 2016, Article ID 2836989

Chiral magnetic effect

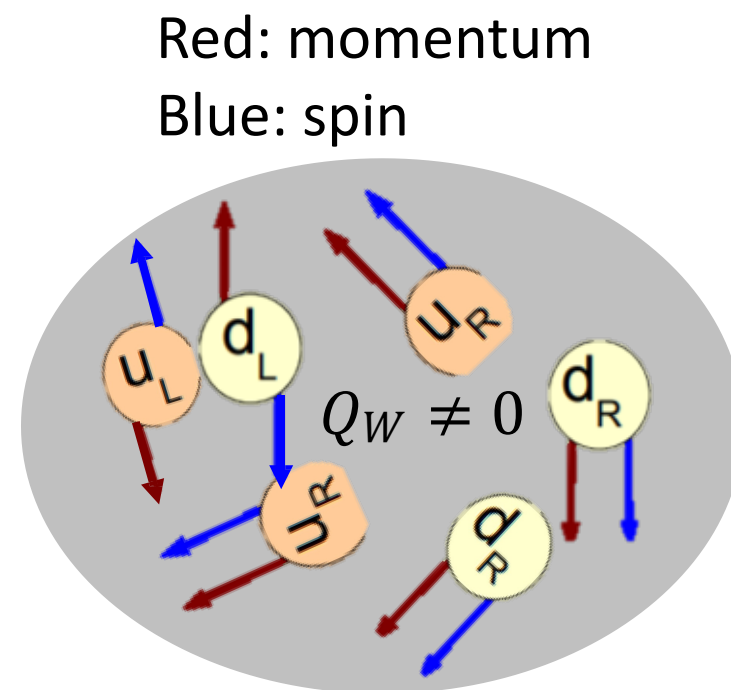
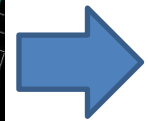
- CME: The generation of electric current along an external magnetic field induced by chirality imbalance.

Chiral magnetic effect

- CME: The generation of electric current along an external magnetic field induced by chirality imbalance.



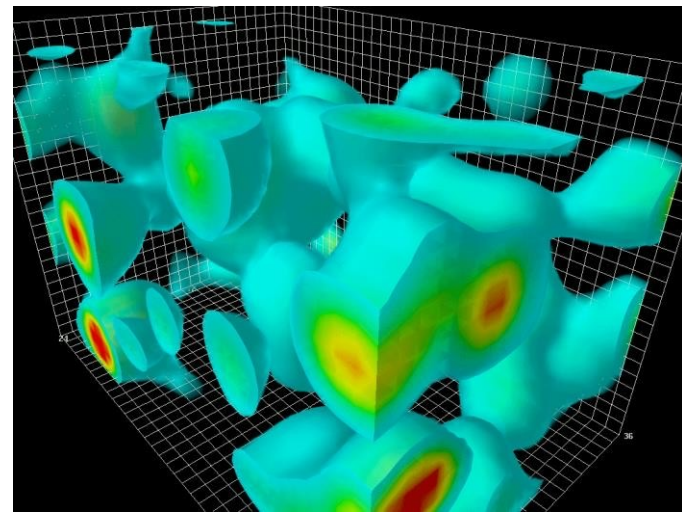
Non-trivial QCD
vacuum structure



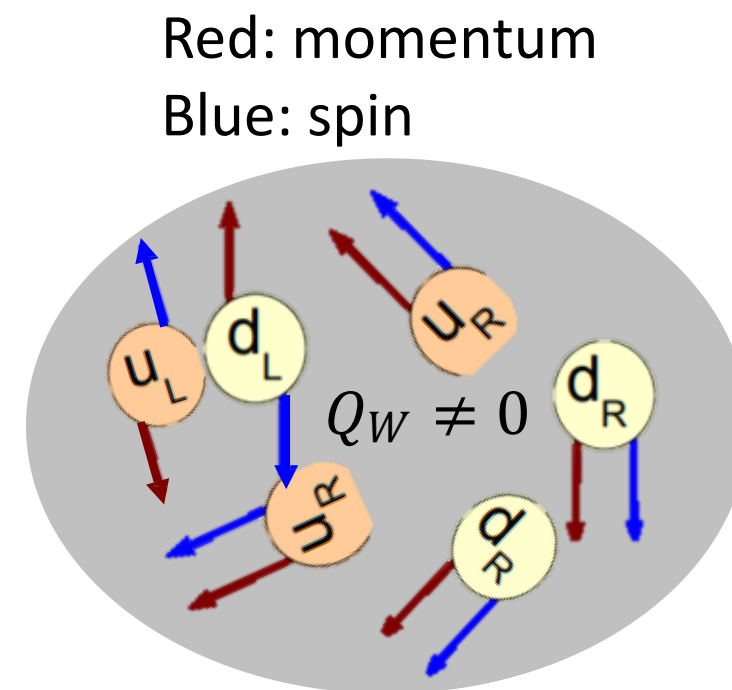
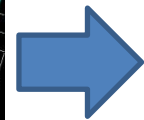
Chirality
imbalance

Chiral magnetic effect

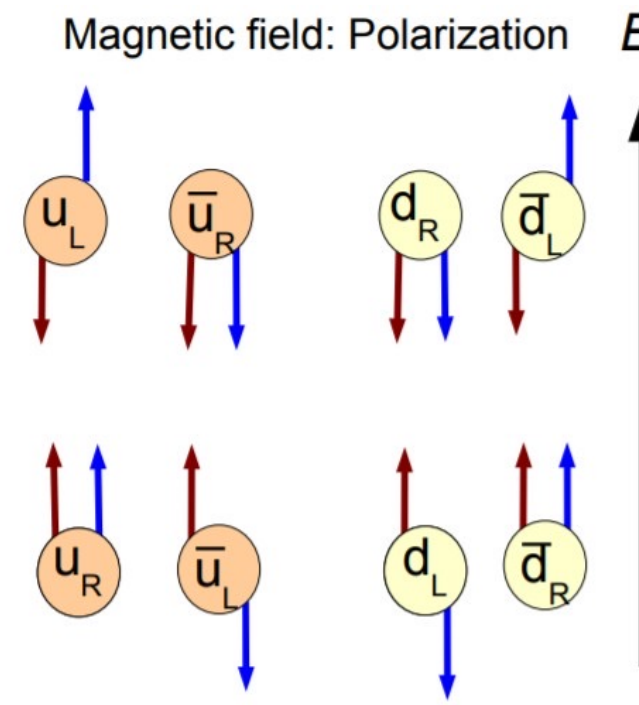
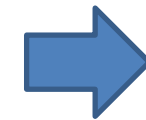
- CME: The generation of electric current along an external magnetic field induced by chirality imbalance.



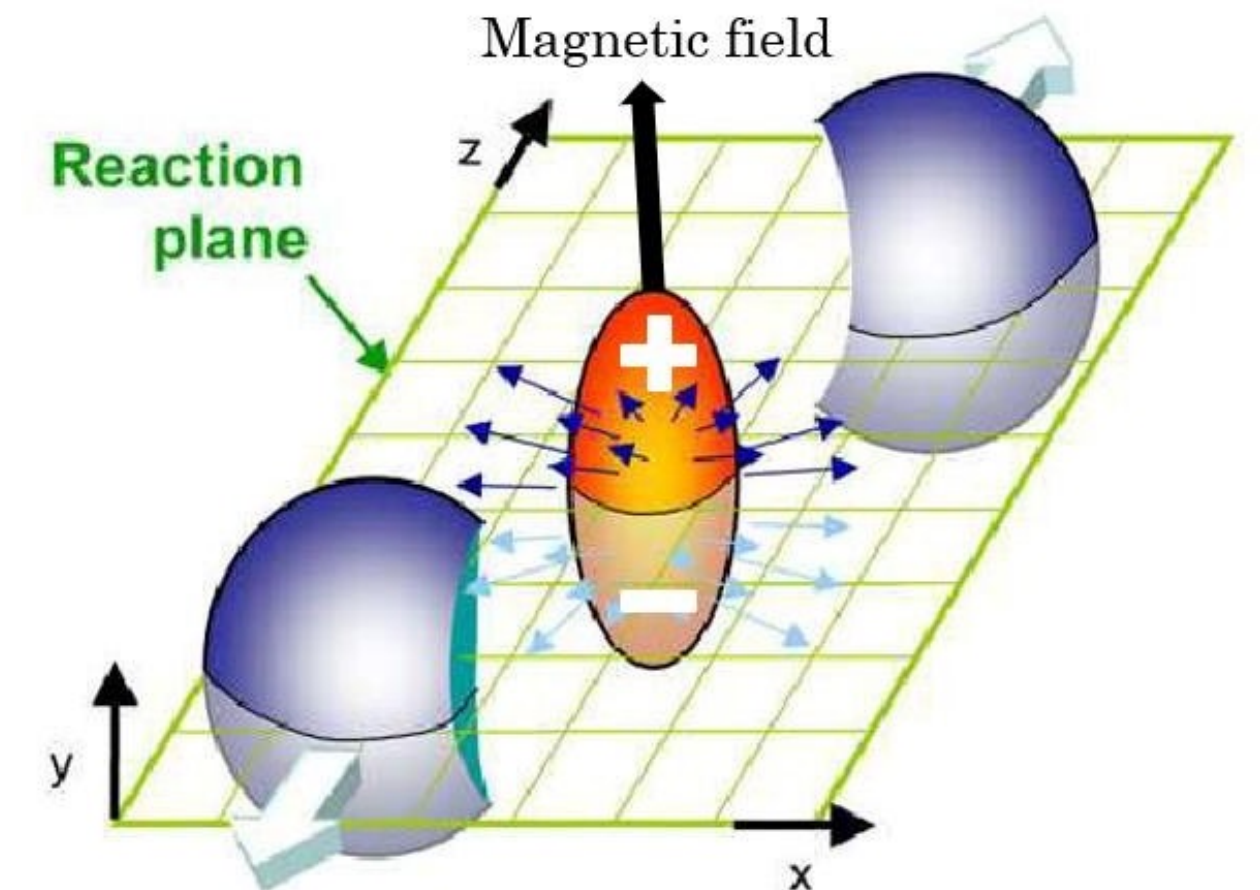
Non-trivial QCD vacuum structure



Chirality imbalance



Under B field, LH quarks have \vec{p} opposite to RH quarks



Current along B field

- Discovery of the CME implies: chiral symmetry restoration, local P/CP violation that may solve the strong CP problem

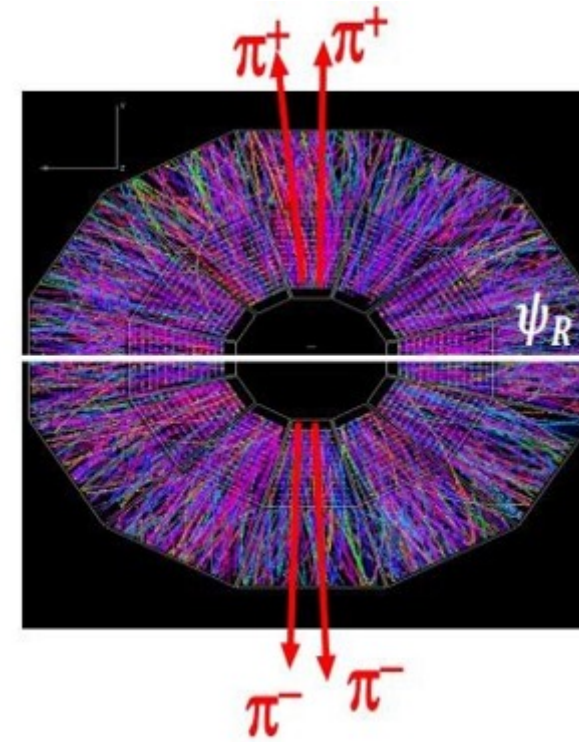
Chiral magnetic effect

- The commonly used $\Delta\gamma$ observable

$$\gamma_{\alpha\beta} = \langle \cos(\varphi_\alpha + \varphi_\beta - 2\Psi_{\text{RP}}) \rangle$$

$$\gamma_{\pm\mp} > 0, \gamma_{\pm\pm} < 0$$

$$\Delta\gamma = \gamma_{\pm\mp} - \gamma_{\pm\pm} > 0$$



Chiral magnetic effect

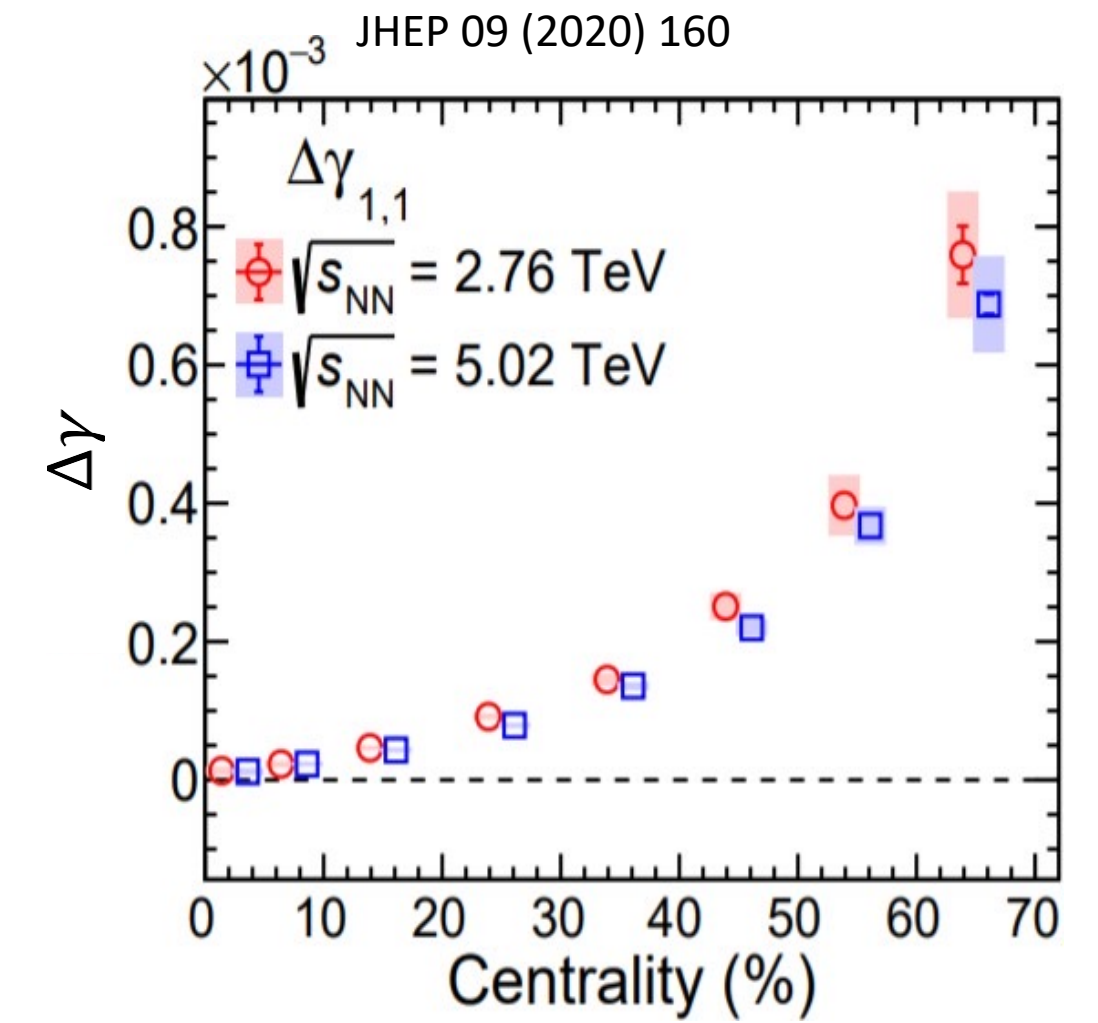
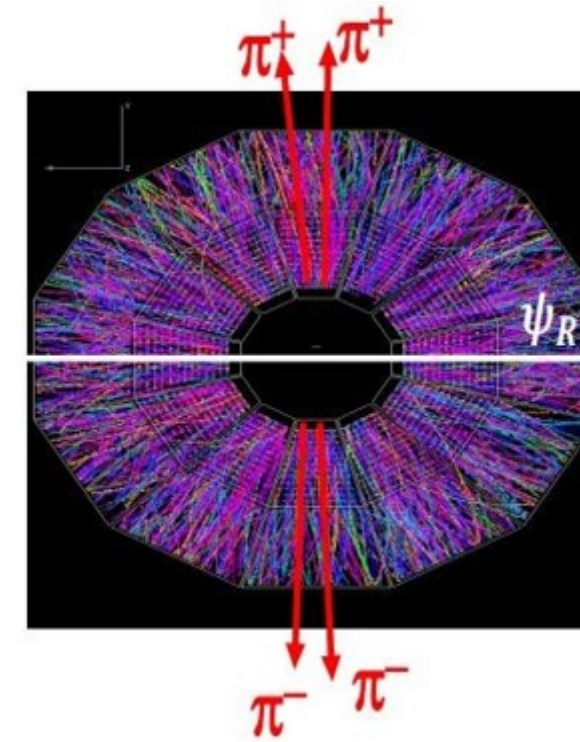
- The commonly used $\Delta\gamma$ observable

$$\gamma_{\alpha\beta} = \langle \cos(\varphi_\alpha + \varphi_\beta - 2\Psi_{\text{RP}}) \rangle$$

$$\gamma_{\pm\mp} > 0, \gamma_{\pm\pm} < 0$$

$$\Delta\gamma = \gamma_{\pm\mp} - \gamma_{\pm\pm} > 0$$

- ALICE previous measurement showed that $\Delta\gamma$ significantly > 0 .



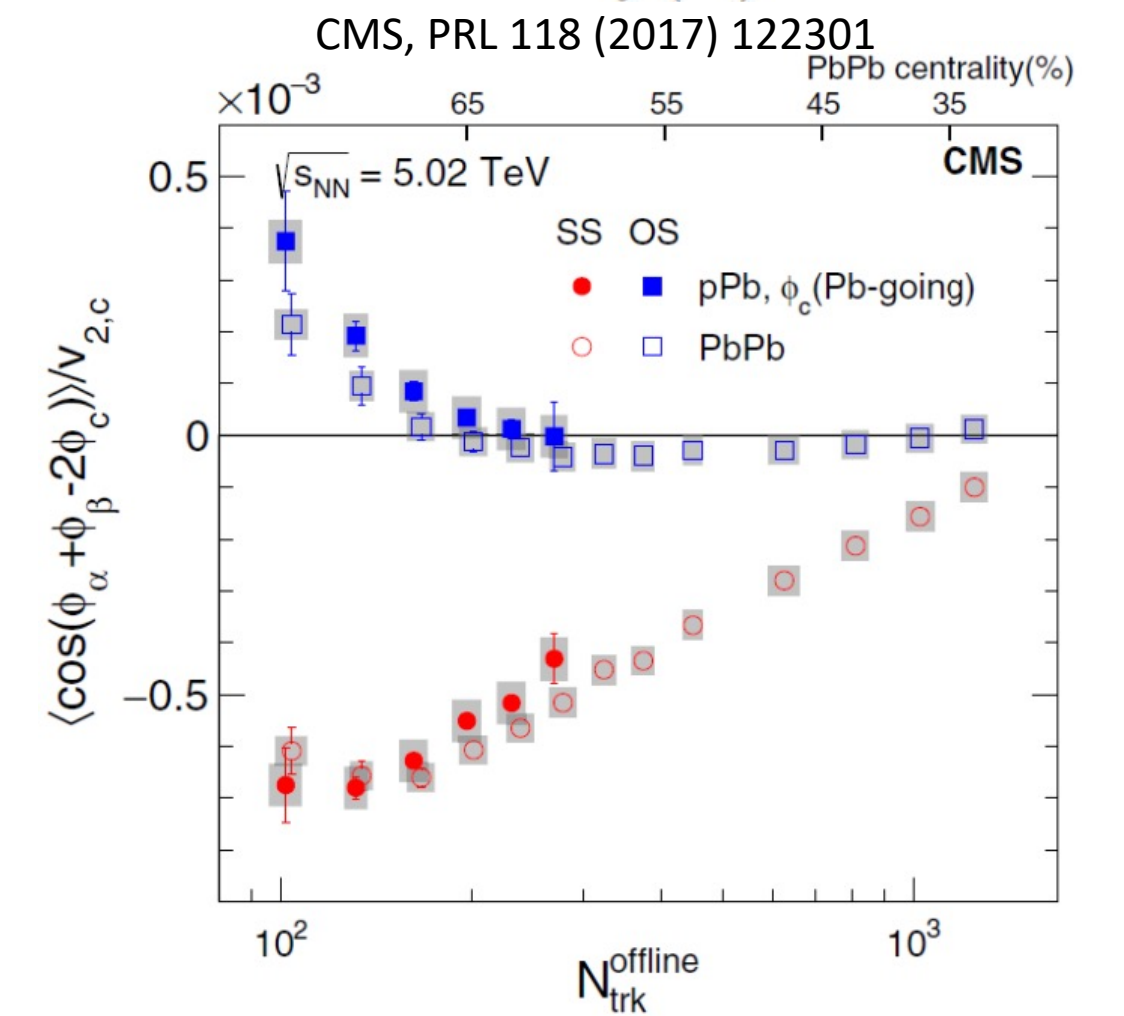
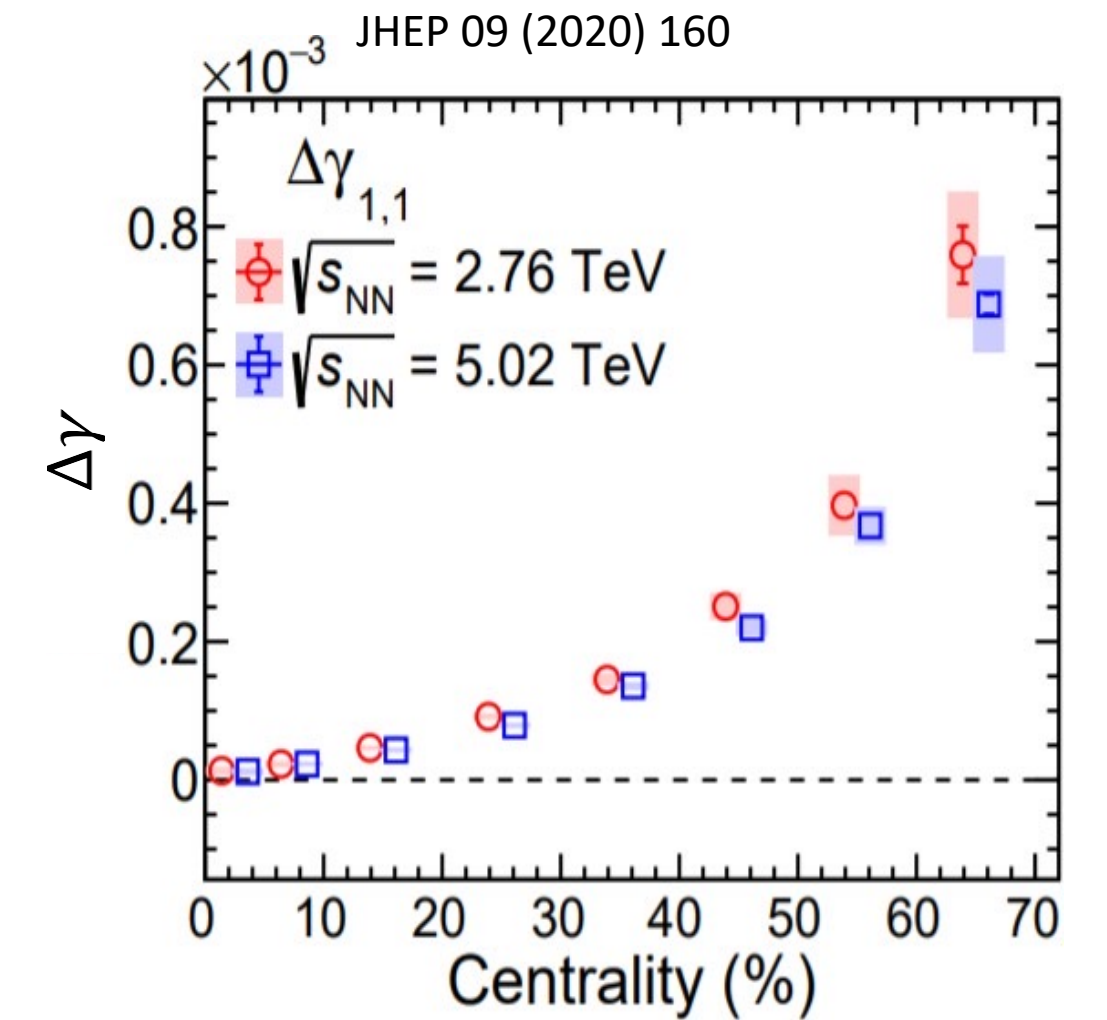
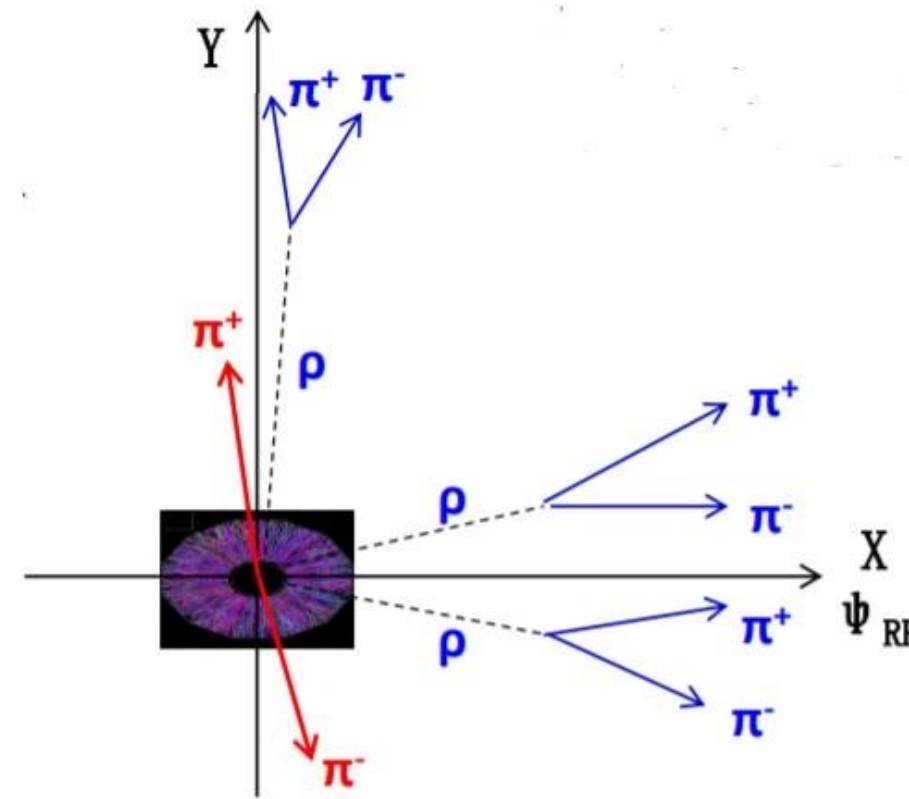
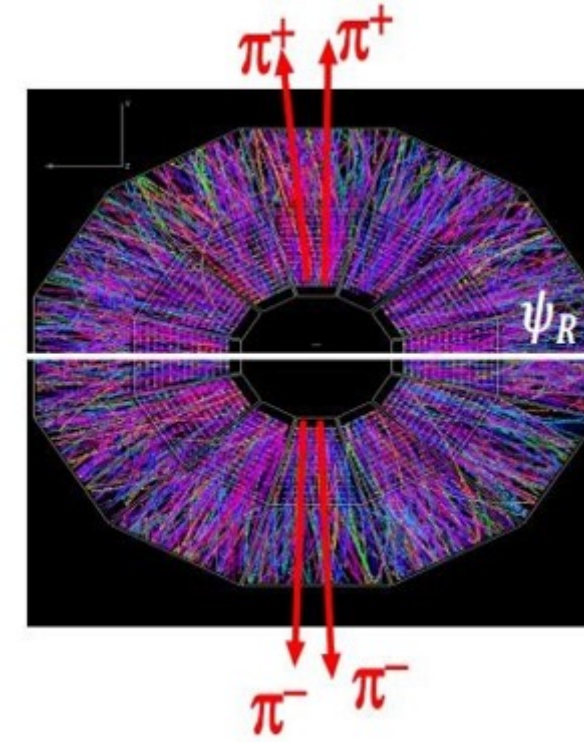
Chiral magnetic effect

- The commonly used $\Delta\gamma$ observable

$$\gamma_{\alpha\beta} = \langle \cos(\varphi_\alpha + \varphi_\beta - 2\Psi_{RP}) \rangle$$

$$\gamma_{\pm\mp} > 0, \gamma_{\pm\pm} < 0$$

$$\Delta\gamma = \gamma_{\pm\mp} - \gamma_{\pm\pm} > 0$$
- ALICE previous measurement showed that $\Delta\gamma$ significantly > 0 .
- However, $\Delta\gamma$ is heavily contaminated by local charge conservation (LCC) and resonance decays, mainly coupled with elliptic flow (noted as v_2)
 - e.g. $\rho^0 \rightarrow \pi^+\pi^-$, more OS pairs align in the Ψ_{RP} than B direction
- Similar value of γ observed in small system (no CME expected, pPb) confirming that the background is huge

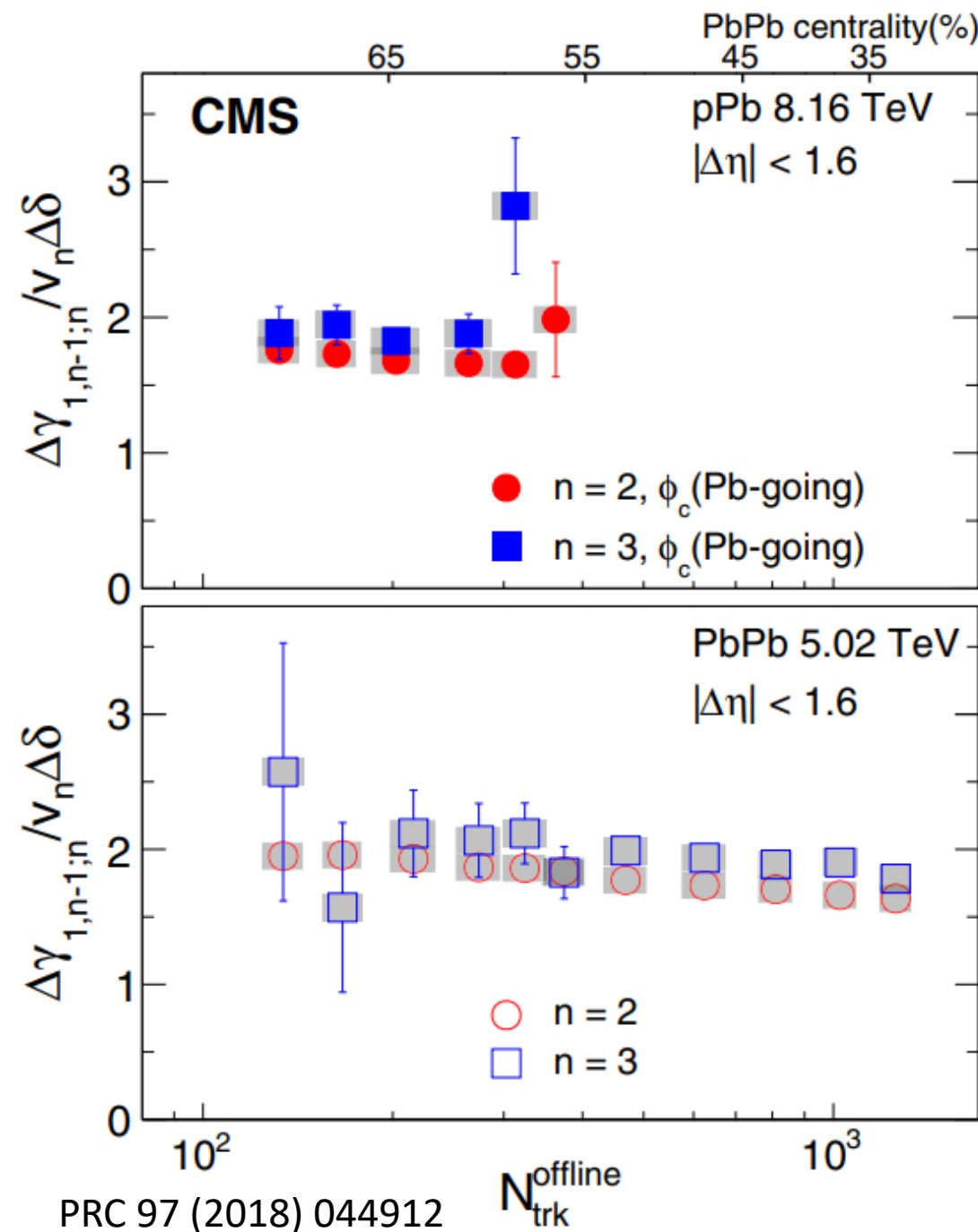


Chiral magnetic effect

- Mix harmonic measurements

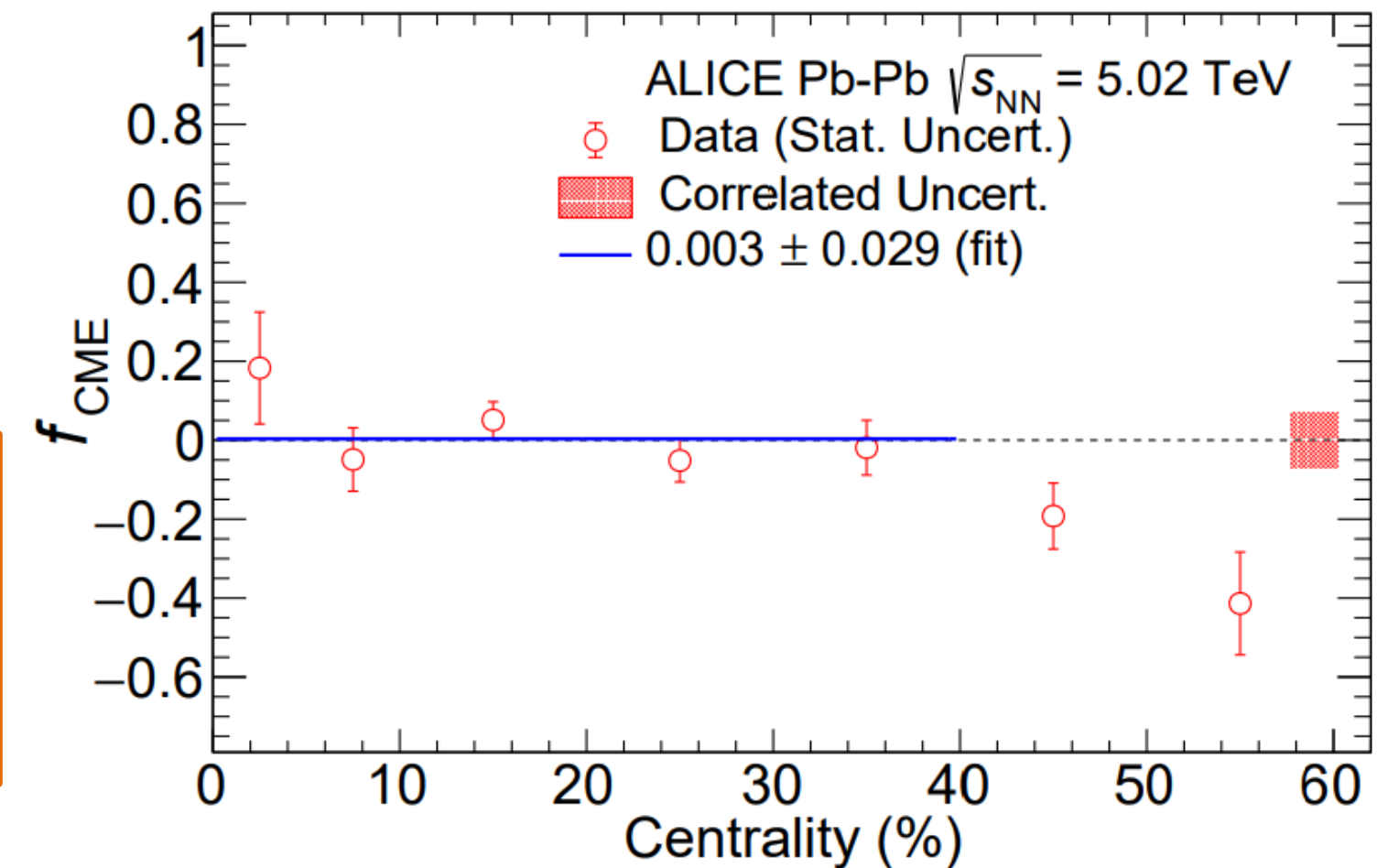
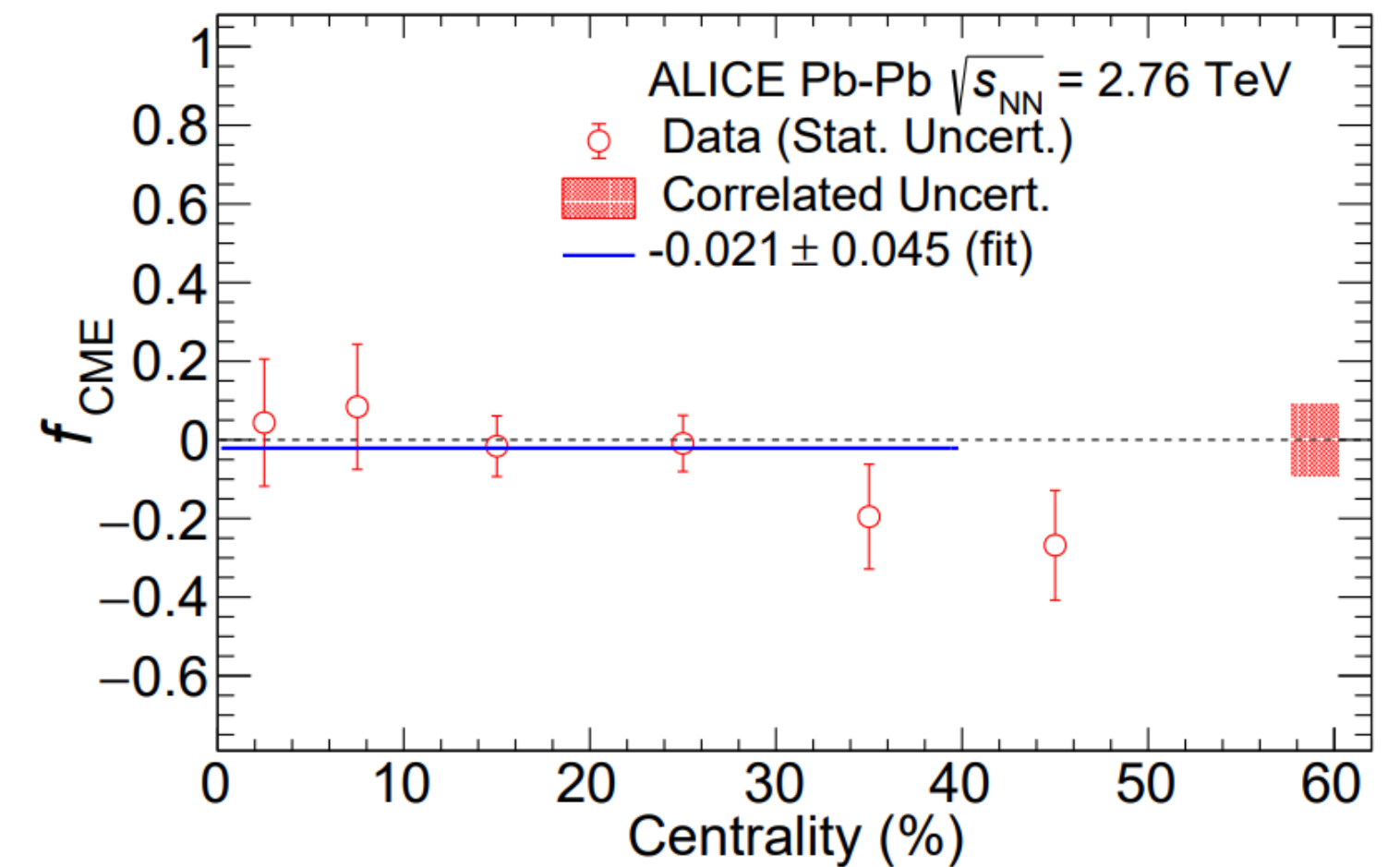
$$\Delta\gamma_{112} \equiv \langle \cos(\phi_\alpha + \phi_\beta - 2\psi_2) \rangle \approx \langle \cos(\phi_\alpha - \phi_\beta) \rangle \langle \cos(2\phi_\beta - 2\psi_2) \rangle \approx \kappa_2 \cdot \Delta\delta \cdot v_2 + \text{CME}$$

$$\Delta\gamma_{123} \equiv \langle \cos(\phi_\alpha + 2\phi_\beta - 3\psi_3) \rangle \approx \langle \cos(\phi_\alpha - \phi_\beta) \rangle \langle \cos(3\phi_\beta - 3\psi_3) \rangle \approx \kappa_3 \cdot \Delta\delta \cdot v_3$$



- $\kappa_2 = \Delta\gamma_{112} / v_2 \Delta\delta$
- $\kappa_3 = \Delta\gamma_{123} / v_3 \Delta\delta$
- Assume $\kappa_2 \approx \kappa_3$
- $\Delta\gamma_{1,1}^{\text{Bkg}} \approx \Delta\gamma_{1,2} \frac{v_2}{v_3} \frac{\kappa_2}{\kappa_3}$
 $\approx \Delta\gamma_{1,2} \frac{v_2}{v_3}$

- Upper limit on fraction of CME for each centrality interval from 0-40%
 - ranging from 15-18% at 95% CL
 - ranging from 20-24% at 99% CL

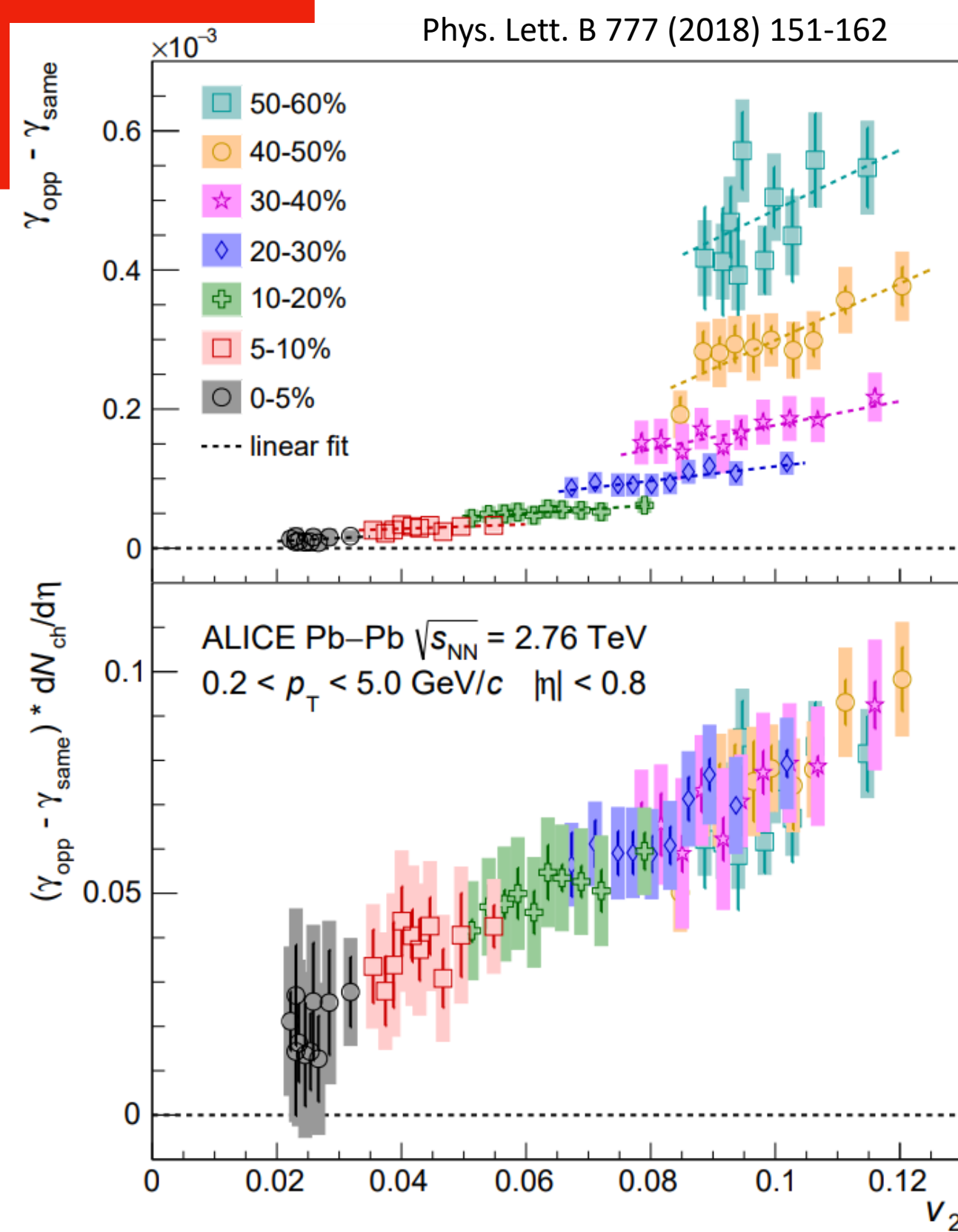
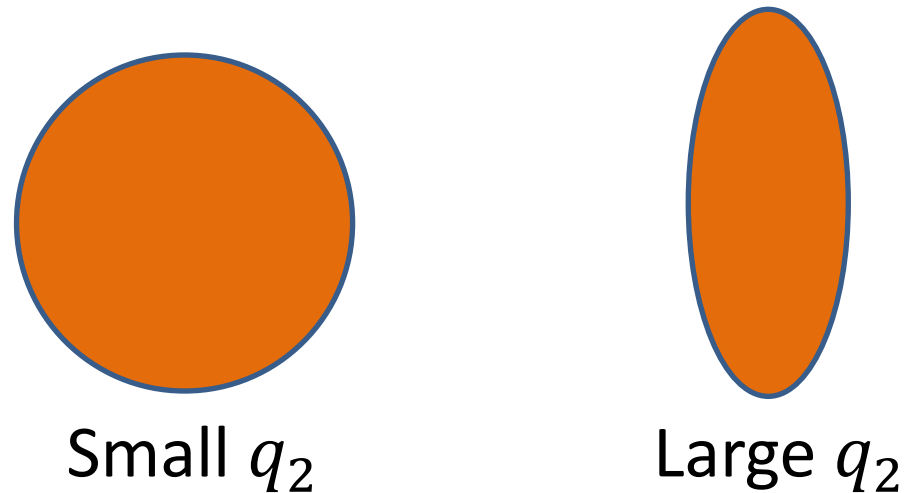


Chiral magnetic effect

- Event-Shape-Engineering
 - Define the second-order reduced flow vector

$$Q_{2,x} = \sum_i^M \cos(2\phi_i), Q_{2,y} = \sum_i^M \sin(2\phi_i)$$

$$q_2 = |\mathbf{Q}_2|/\sqrt{M} = \sqrt{Q_{2,x}^2 + Q_{2,y}^2}/\sqrt{M}$$
 - Use q_2 to select different geometry. Use v_2 to quantify event anisotropy
 - A significant CME contribution \rightarrow non-zero intercepts at $v_2 = 0$ (still the non-zero intercept has background from LCC)



Chiral magnetic effect

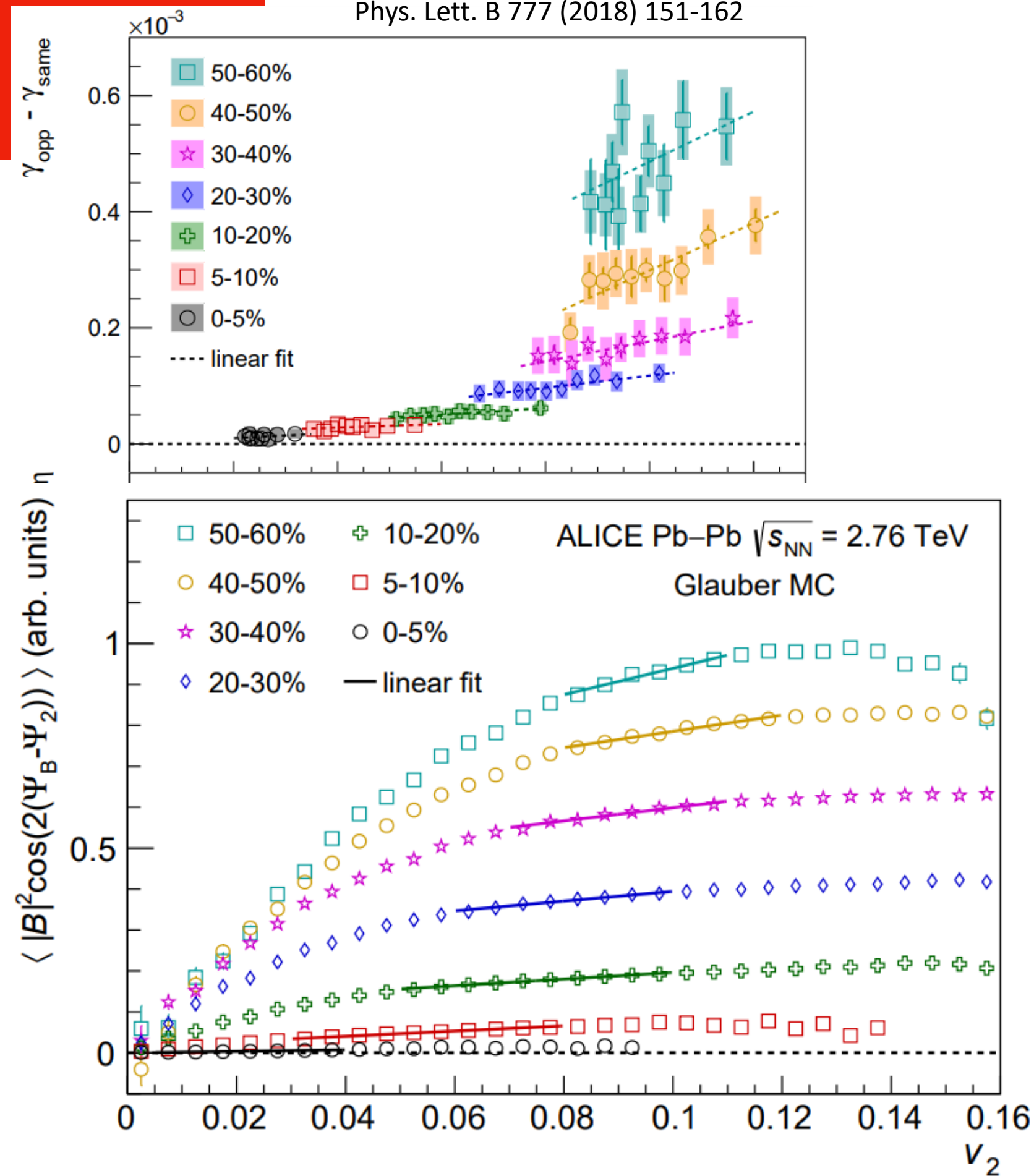
- Event-Shape-Engineering

- Define the second-order reduced flow vector

$$Q_{2,x} = \sum_i^M \cos(2\phi_i), Q_{2,y} = \sum_i^M \sin(2\phi_i)$$

$$q_2 = |\mathbf{Q}_2|/\sqrt{M} = \sqrt{Q_{2,x}^2 + Q_{2,y}^2}/\sqrt{M}$$

- Use q_2 to select different geometry. Use v_2 to quantify event anisotropy
- A significant CME contribution \rightarrow non-zero intercepts at $v_2 = 0$ (still the non-zero intercept has background from LCC)
- Using three different initial-state models to estimate the dependence of CME signal to v_2



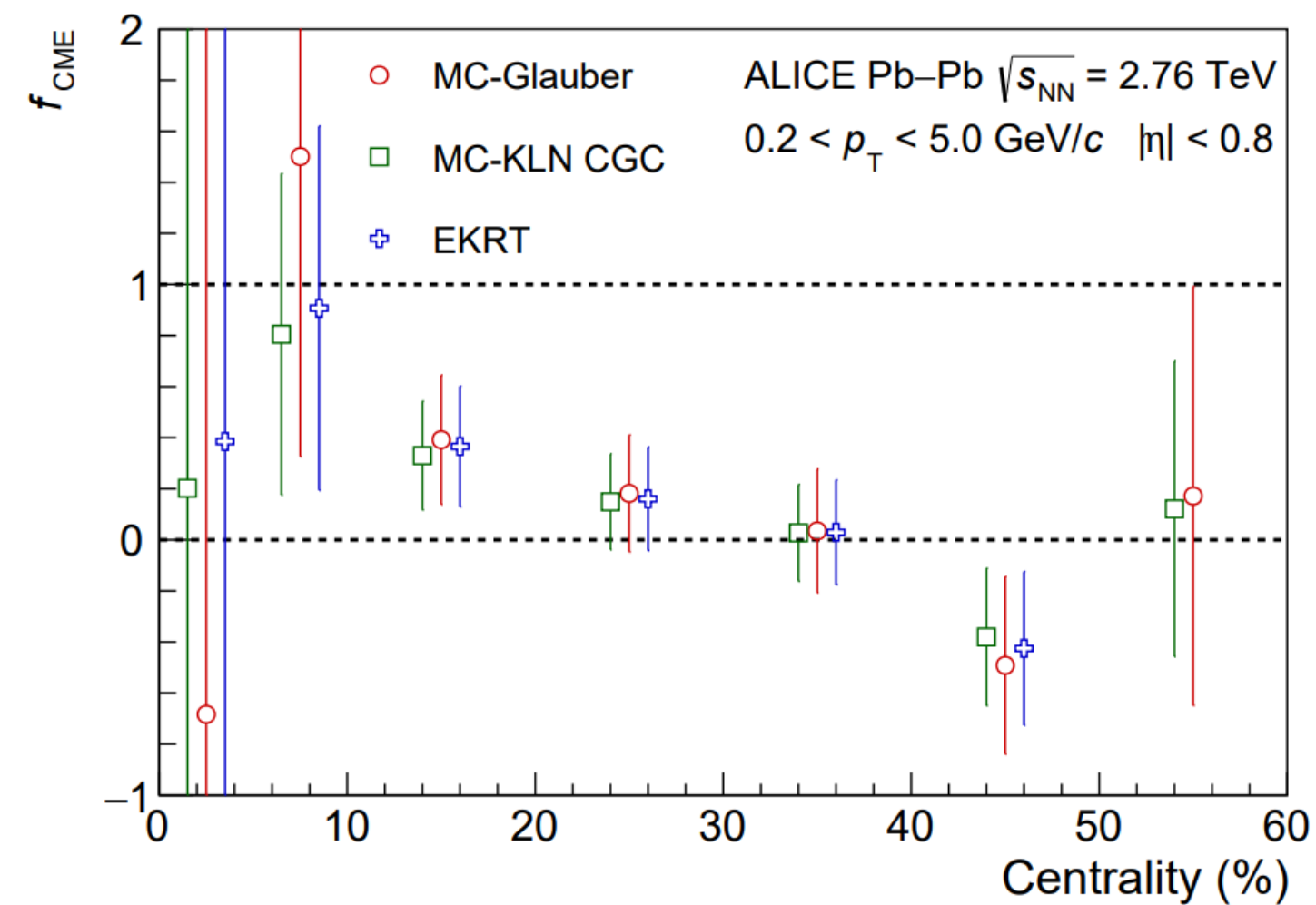
Chiral magnetic effect

- Event-Shape-Engineering
 - Define the second-order reduced flow vector

$$Q_{2,x} = \sum_i^M \cos(2\phi_i), Q_{2,y} = \sum_i^M \sin(2\phi_i)$$

$$q_2 = |\mathbf{Q}_2|/\sqrt{M} = \sqrt{Q_{2,x}^2 + Q_{2,y}^2}/\sqrt{M}$$

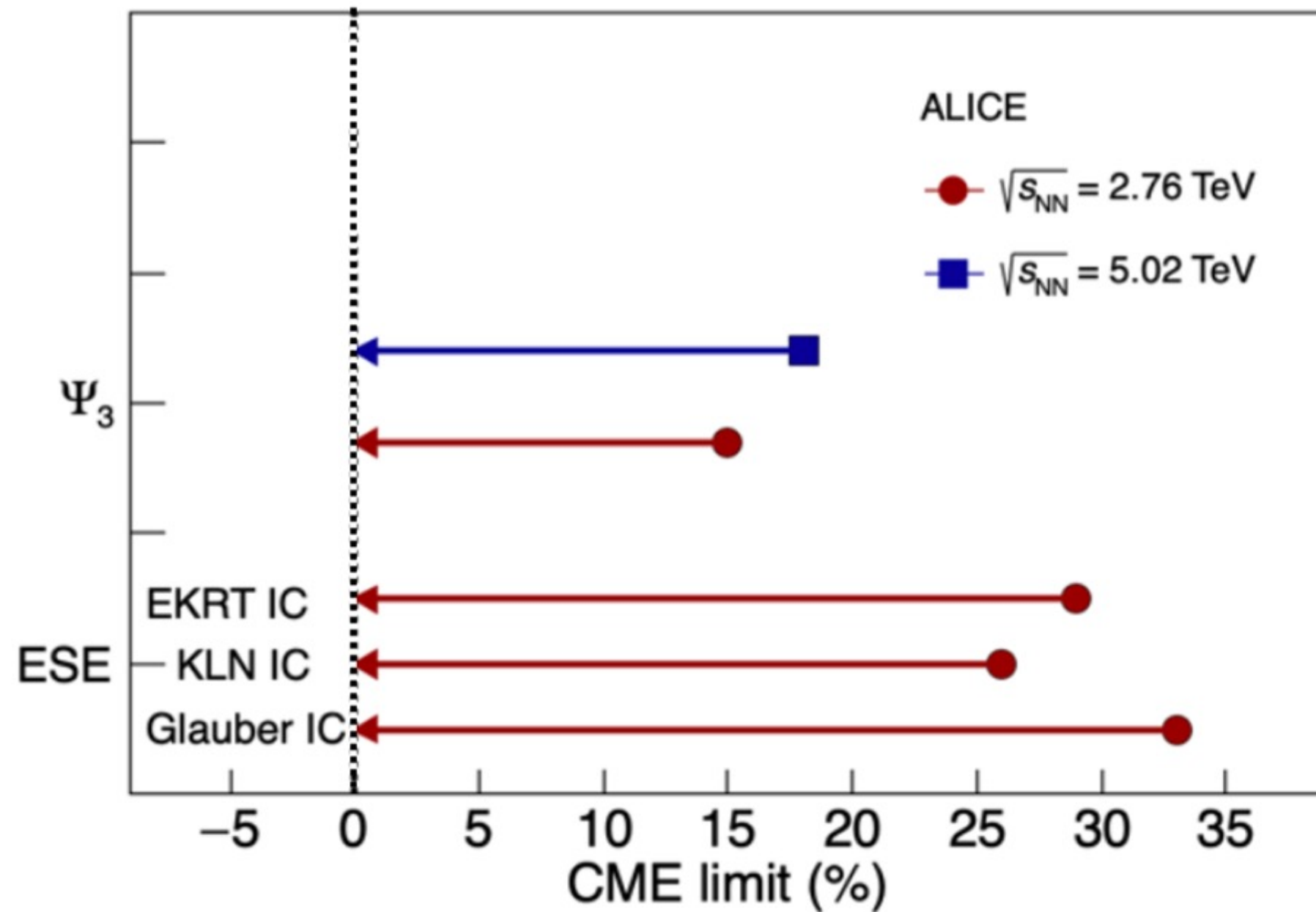
- Use q_2 to select different geometry. Use v_2 to quantify event anisotropy
- A significant CME contribution \rightarrow non-zero intercepts at $v_2 = 0$ (still the non-zero intercept has background from LCC)
- Using three different initial-state models to estimate the dependence of CME signal to v_2
- The extracted CME fraction for three models used in the study suggests an upper limit on f_{CME} of 26 to 33% (depending on initial-state models) at 95% CL for the 10-50% centrality interval



Weak dependence on initial states

Chiral magnetic effect

- A summary of the results for the upper limits for the CME signal at ALICE



- It is clear that background is dominant
- CME is a very important physics. We have to quantify it
- Currently, the work on using the Spectator/participant planes method is ongoing

PhysRevC.105.024913

Summary

- The early-stage EM field can affect the motion of the final state particles via measurements of directed flow for particles and antiparticles.
- The differences in the measured global polarisation of Λ and $\bar{\Lambda}$ provide an upper limit for the magnitude of the magnetic field
- Direct studies for the existence of the CME in Pb-Pb collisions yielded upper limits of 26 to 33% for ESE and 15-18% at 95% CL for mix harmonic for the 0–40% centrality interval

THANKS

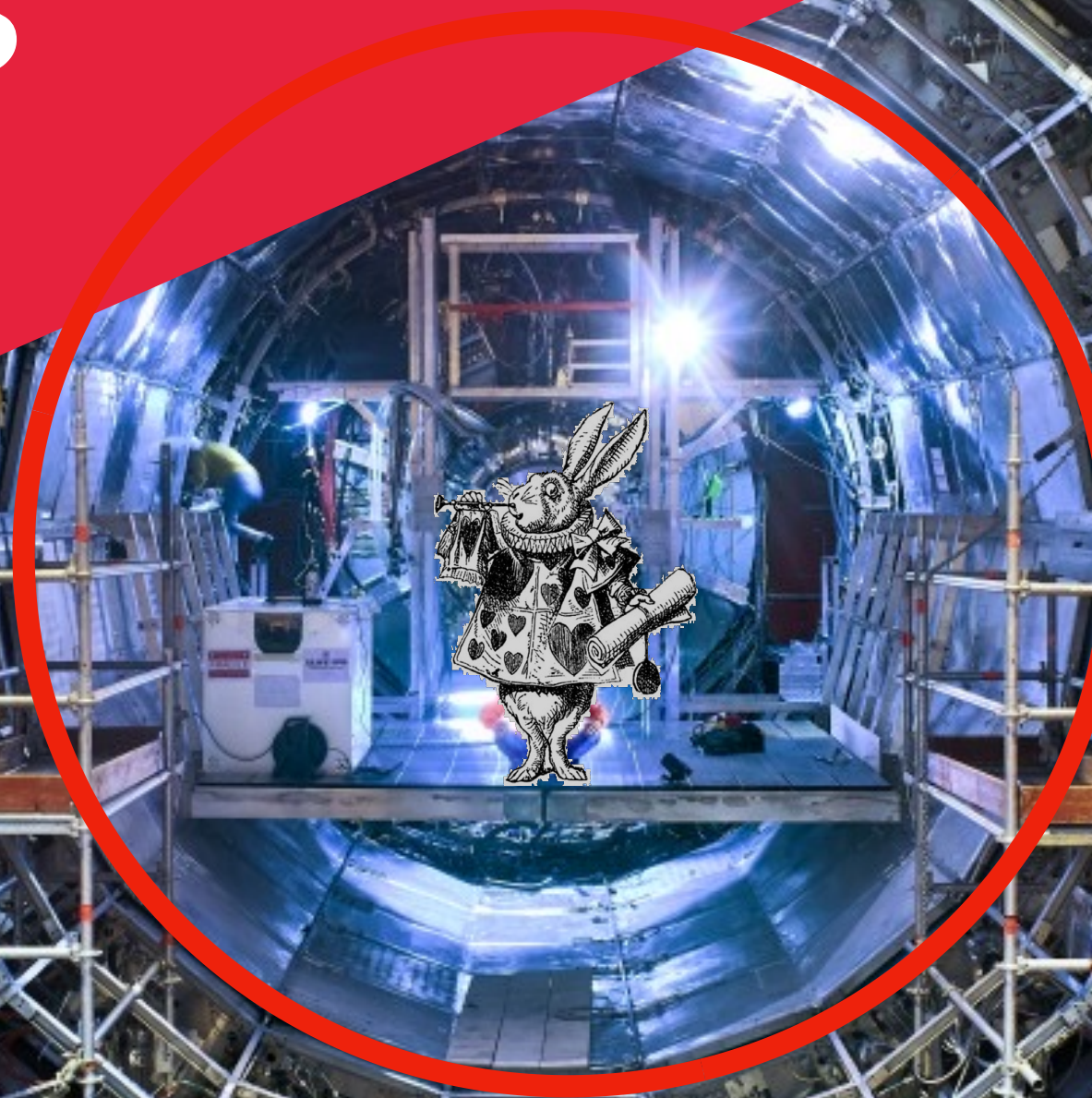


ALICE

Extra Slides



ALICE



AREA
B

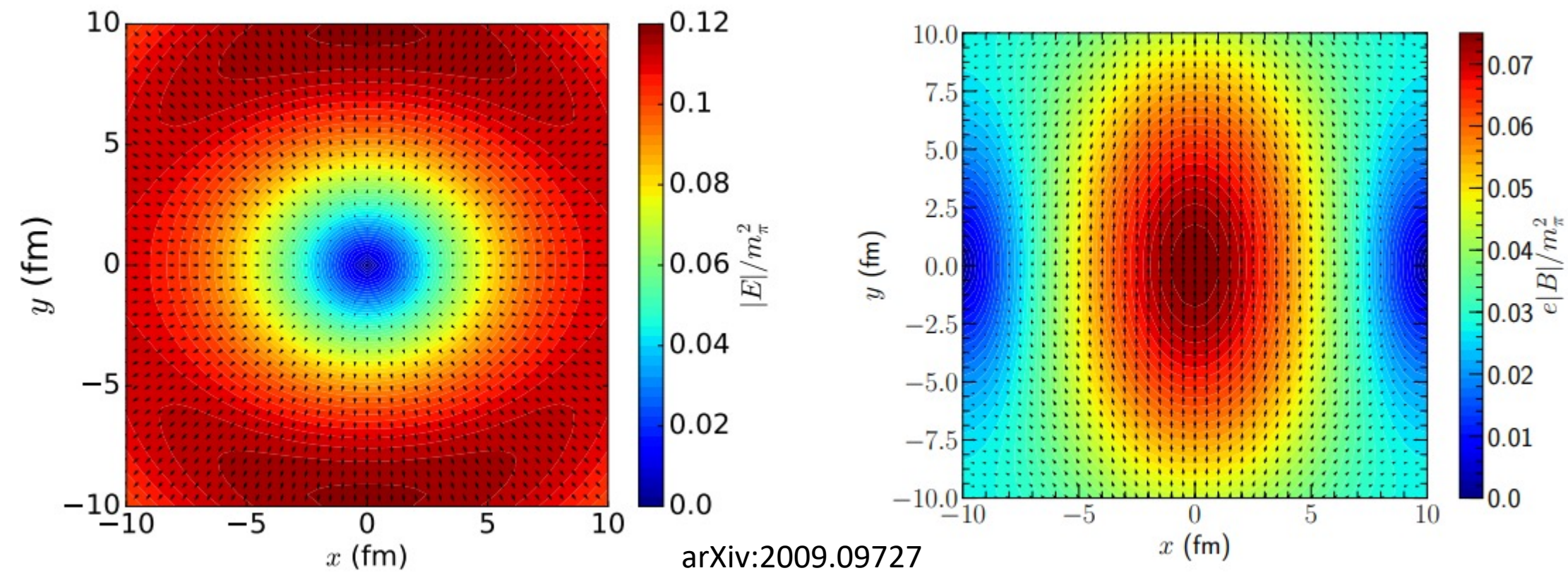
AREA
B

38 Nikhef
Spi Qiu



Nikhef Jamboree, 10th 5/2022

EB field simulation



Simulation of the electric (left) and magnetic (right) fields in the transverse plane after a Pb–Pb collision at 2.76TeV with 20–30% centrality