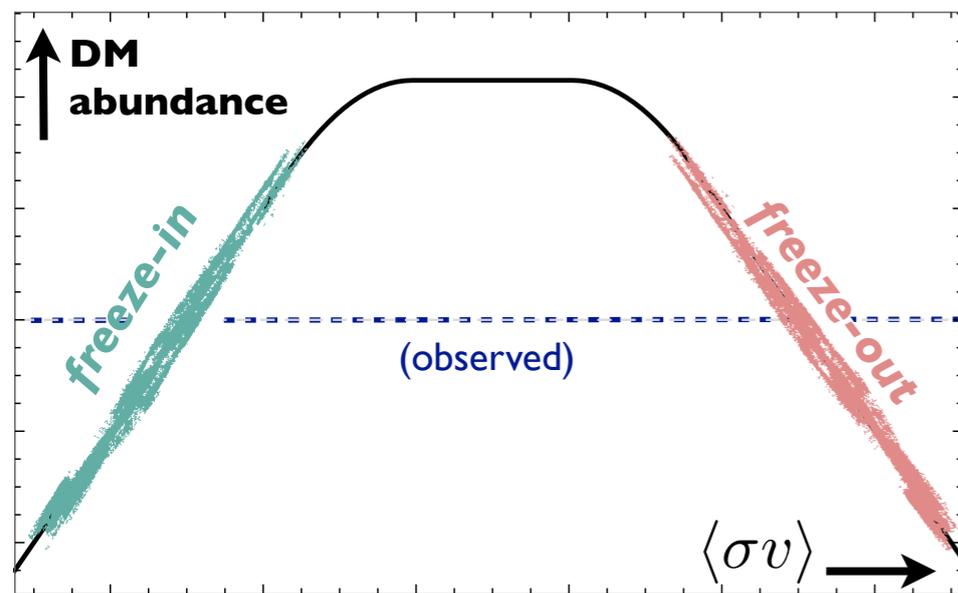


Dark Matter from the Primordial Plasma

From freeze-in to pandemic production

Torsten Bringmann



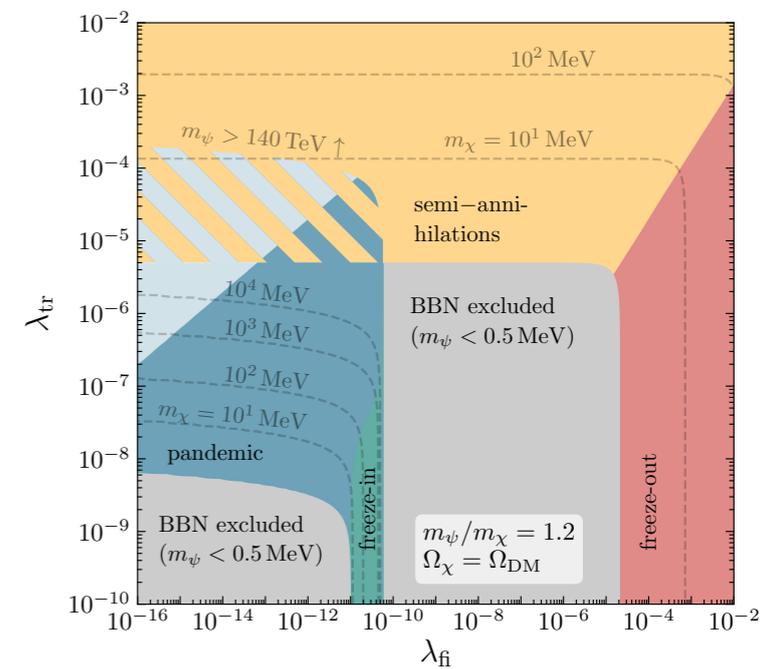
standard production mechanisms

Based on

TB, Heeba, Kahlhoefer & Vangsnes, JHEP '22

TB, Depta, Hufnagel, Ruderman & Schmidt-Hoberg, PRL '21

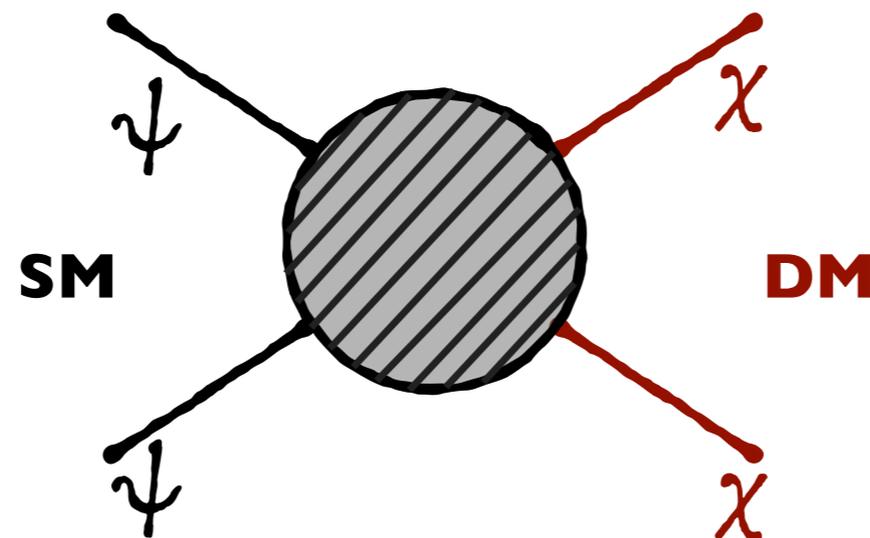
TB, Depta, Hufnagel, Kersten, Ruderman & Schmidt-Hoberg, arXiv:2206.10630



'pandemic' production

The origin of dark matter

- Existence of (particle) DM = **evidence** for BSM physics
- Guiding principle: any convincing model for dark matter must include a **production mechanism**
- Let's postulate some interaction with the primordial heat bath:
 - [Z_2 symmetry not strictly necessary, but automatically guarantees stability of DM]



Boltzmann equation

- Evolution of DM phase-space density:

$$L[f_\chi] = C[f_\chi]$$

“ $\frac{df_\chi}{dt}$ ” = $E (\partial_t - H\mathbf{p} \cdot \nabla_{\mathbf{p}}) f_\chi$ for FRW metric

- Collision term:

$$C[f_\chi] = \frac{1}{2g_\chi} \int \frac{d^3\tilde{p}}{(2\pi)^3 2\tilde{E}} \int \frac{d^3k}{(2\pi)^3 2\omega} \int \frac{d^3\tilde{k}}{(2\pi)^3 2\tilde{\omega}} (2\pi)^4 \delta^{(4)}(\tilde{p} + p - \tilde{k} - k)$$

$$\times \left[\underbrace{|\mathcal{M}|_{\bar{\chi}\chi \leftarrow \bar{\psi}\psi}^2}_{\text{‘production’}} f_\psi(\omega) f_\psi(\tilde{\omega}) - \underbrace{|\mathcal{M}|_{\bar{\chi}\chi \rightarrow \bar{\psi}\psi}^2}_{\text{‘annihilation’}} f_\chi(E) f_\chi(\tilde{E}) \{1 \pm f_\psi(\omega)\} \{1 \pm f_\psi(\tilde{\omega})\} \right]$$

- Detailed balance:** ‘production’ = ‘annihilation’ in equilibrium

↪ allows to re-write everything in terms of a *would-be* equilibrium population of DM

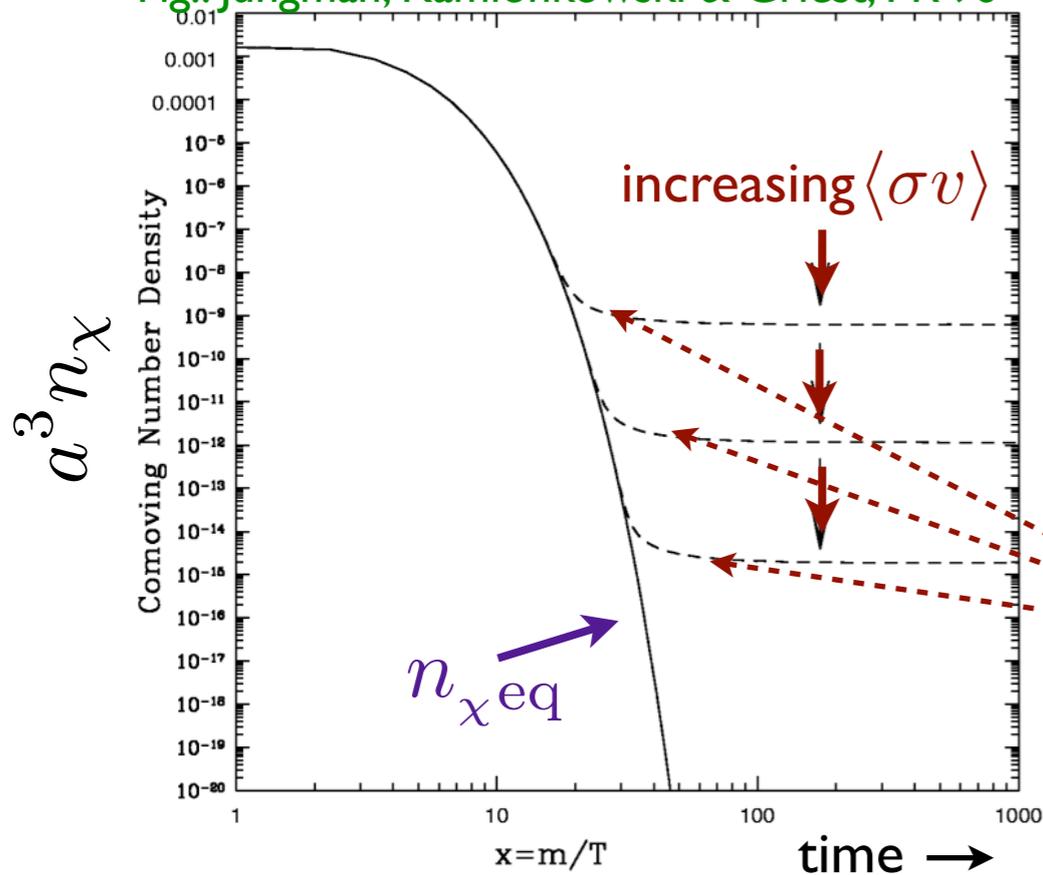
- Pauli suppression / Bose enhancement**

↪ not relevant for production of **non-relativistic DM** [energy conservation!]

Weakly Interacting Massive Particles

- well-motivated from particle physics [SUSY, EDs, ...]
- thermal production in early universe:

Fig.: Jungman, Kamionkowski & Griest, PR'96



$$\frac{dn_\chi}{dt} + 3Hn_\chi = -\langle\sigma v\rangle (n_\chi^2 - n_{\chi\text{eq}}^2)$$

$\langle\sigma v\rangle$: $\chi\chi \rightarrow \text{SM SM}$
(thermal average)



“Freeze-out” when annihilation rate falls behind expansion rate

→ Relic density (today): $\Omega_\chi h^2 \sim \frac{3 \cdot 10^{-27} \text{ cm}^3/\text{s}}{\langle\sigma v\rangle}$

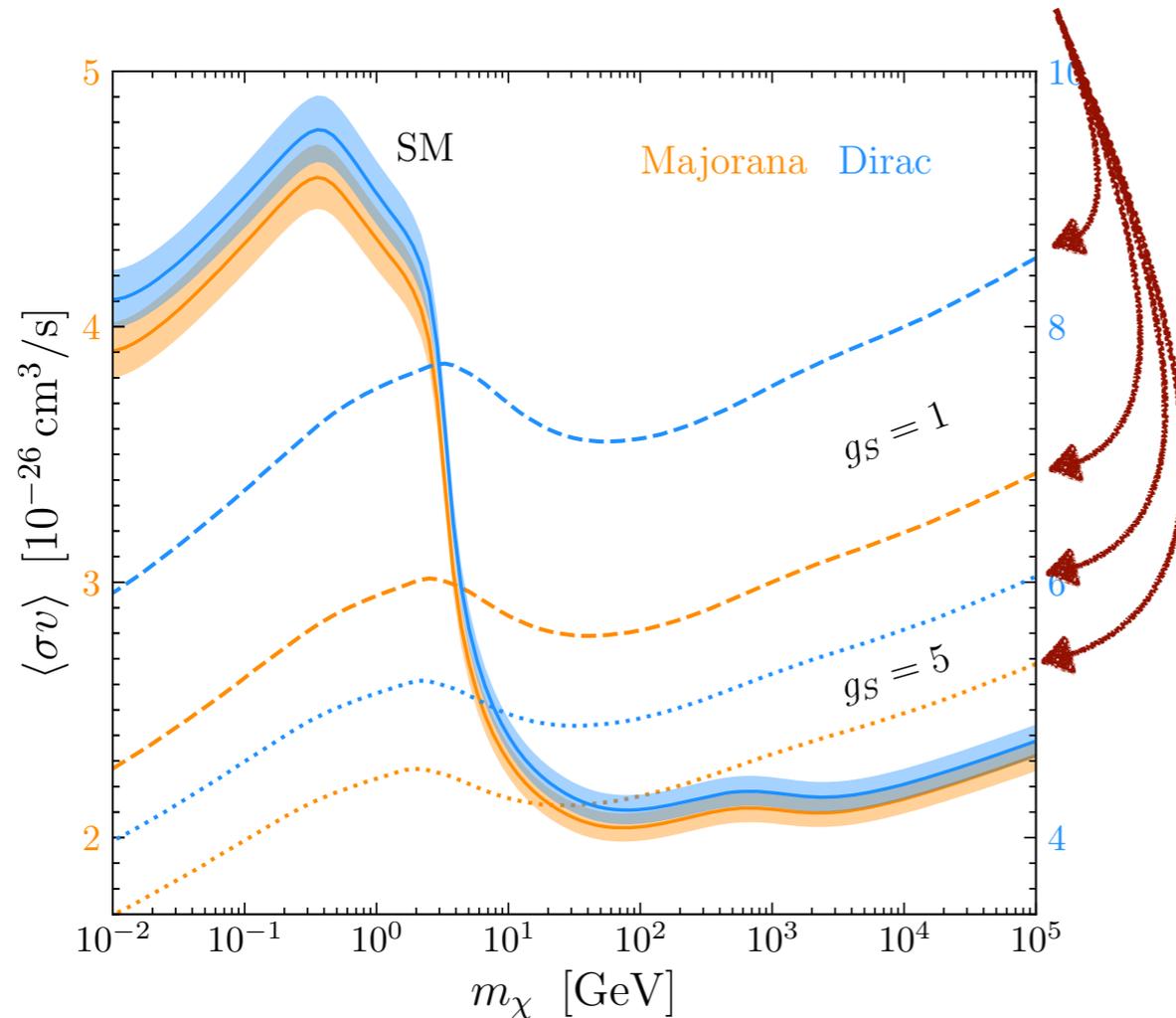
's-wave' vs. 'p-wave'

- Precision calculations of 'thermal cross section':

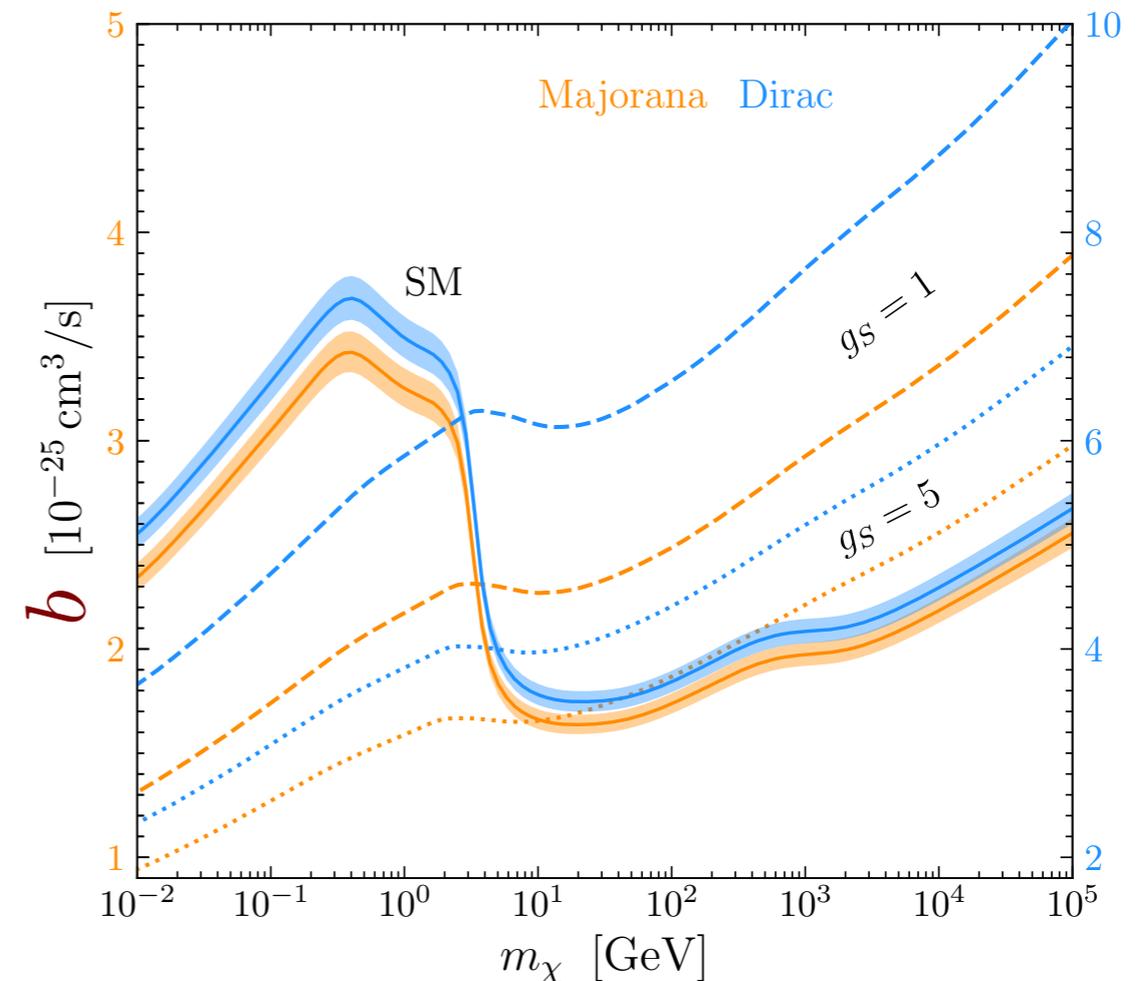
TB, Depta, Hufnagel & Schmidt-Hoberg, PLB '21



- works also in secluded dark sectors (with g_S heat bath d.o.f.)



$$\langle\sigma v\rangle = \text{const.}$$



$$\langle\sigma v\rangle = b v^2$$



- $v \sim 0.3$ around freeze-out
- $v \sim 10^{-3}$ today

Feably Interacting Massive Particles

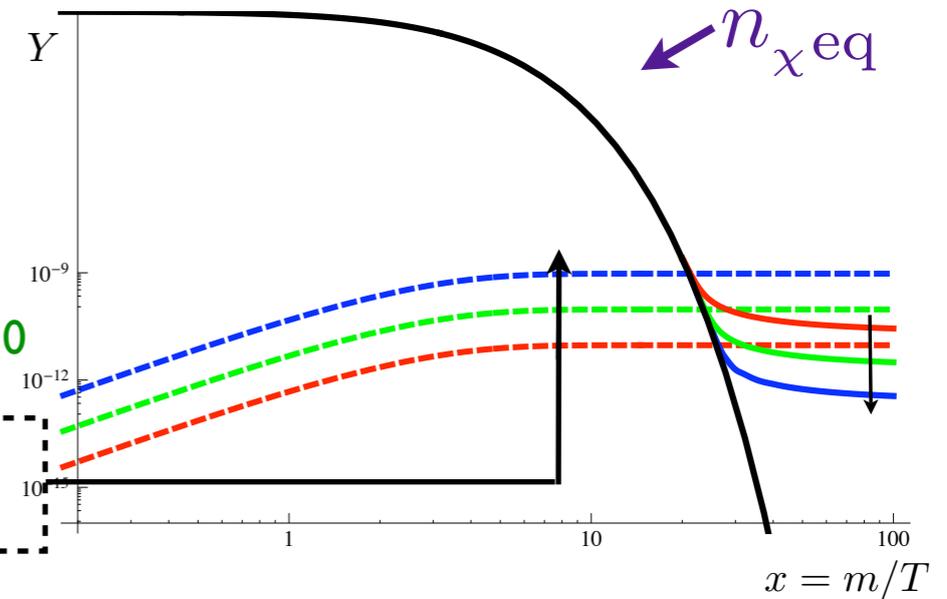
- WIMP DM is seriously **pressured** (but not dead!)

- alternative **production** from **thermal bath** with much smaller interactions

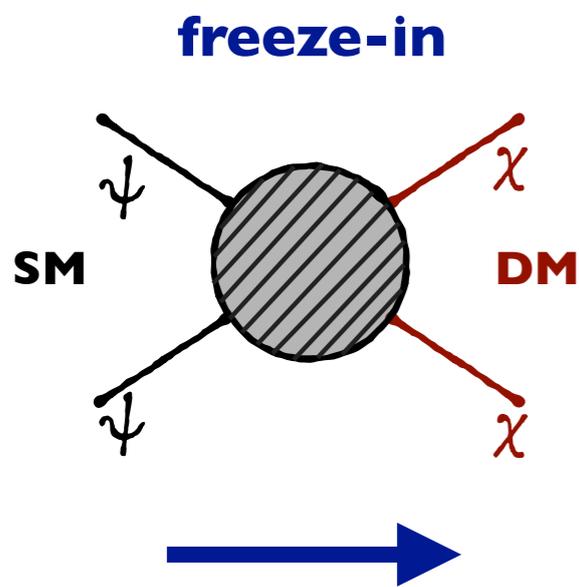
- DM never equilibrates in this case

Hall+, JHEP '10

Increasing interaction rate for **freeze-in**



- Smooth transition between two regimes:



depends on **initial conditions**

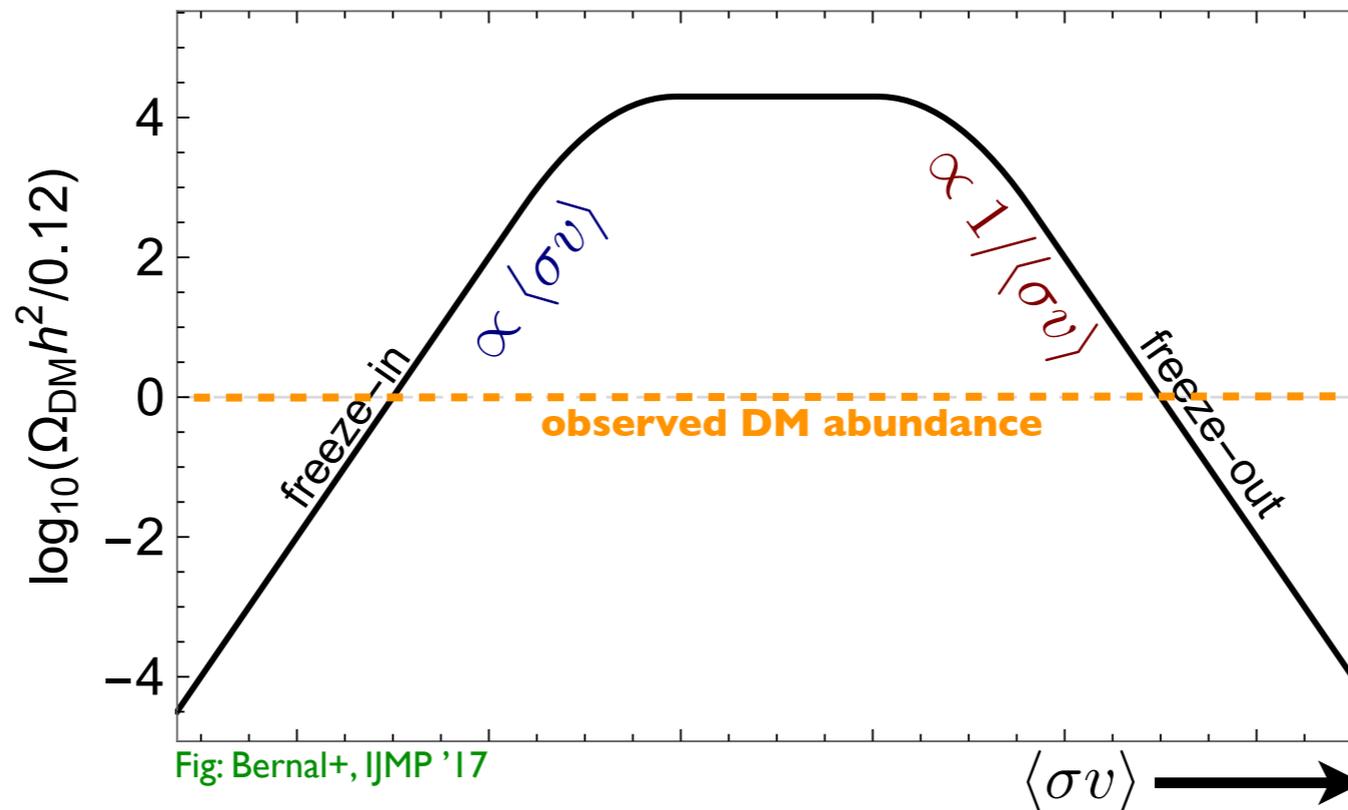
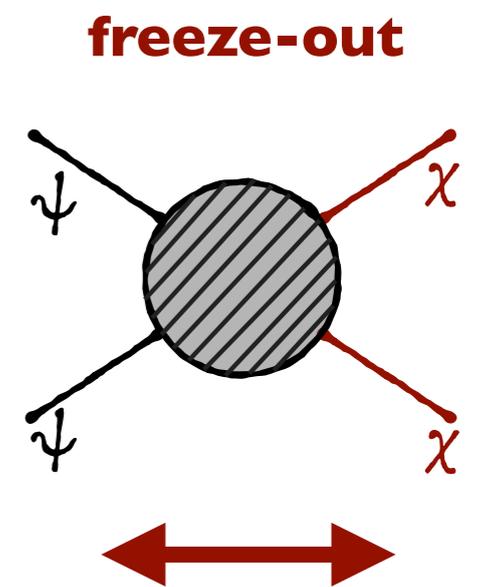


Fig: Bernal+, IJMP '17



insensitive to initial conditions



Collision term for FIMPs

$$C[f_\chi] = \frac{1}{2g_\chi} \int \frac{d^3\tilde{p}}{(2\pi)^3 2\tilde{E}} \int \frac{d^3k}{(2\pi)^3 2\omega} \int \frac{d^3\tilde{k}}{(2\pi)^3 2\tilde{\omega}} (2\pi)^4 \delta^{(4)}(\tilde{p} + p - \tilde{k} - k) \\ \times \left[\underbrace{|\mathcal{M}|_{\chi\chi \leftarrow \psi\psi}^2}_{\text{'production'}} f_\psi(\omega) f_\psi(\tilde{\omega}) - \underbrace{|\mathcal{M}|_{\chi\chi \rightarrow \psi\psi}^2}_{\text{'annihilation'}} f_\chi(E) f_\chi(\tilde{E}) \{1 \pm f_\psi(\omega)\} \{1 \pm f_\psi(\tilde{\omega})\} \right]$$


 $C[f_\chi] = \langle \sigma v \rangle_{\chi\chi \rightarrow \psi\psi} (n_\chi^{\text{MB}})^2$
annihilation of *would-be* MB population
→ **actual** (in eq)

- Only 2 integrals after exploiting spherical symmetry:
[γ : Lorentz boost to CMS; $\tilde{s} = s/(4m_\chi^2)$]

$$\langle \sigma v \rangle_{\chi\chi \rightarrow \psi\psi} = \frac{8x^2}{K_2^2(x)} \int_1^\infty d\tilde{s} \tilde{s} (\tilde{s} - 1) \int_1^\infty d\gamma \sqrt{\gamma^2 - 1} e^{-2\sqrt{\tilde{s}}x\gamma} \sigma_{\chi\chi \rightarrow \psi\psi}(s, \gamma)$$

→ $K_1(2\sqrt{\tilde{s}}x)/(2\sqrt{\tilde{s}}x)$ ✓

TB, Heeba, Kahlhoefer & Vangsnes, JHEP '22
 (see also Lebedev & Toma, PLB '19 Arcadi+, JHEP '19)

$$\sigma_{\chi\chi \rightarrow \psi\psi}(p, \tilde{p}) = \frac{(2\pi)^4}{4N_\psi E \tilde{E} v_{\text{Møll}}} \int \frac{d^3k}{(2\pi)^3 2\omega} \int \frac{d^3\tilde{k}}{(2\pi)^3 2\tilde{\omega}} \delta^{(4)}(\tilde{p} + p - \tilde{k} - k) |\overline{\mathcal{M}}|^2 \{1 \pm f_\psi(\omega)\} \{1 \pm f_\psi(\tilde{\omega})\}$$

→ 1

- In *this* formulation, direct analogy with **WIMP** case!

- 
 - Can recycle sophisticated numerical tools for thermal averages
 - Easier to estimate higher-order corrections



(Further) Finite temperature effects

- FIMPs are dominantly produced at **higher temperatures** than WIMPs, and over a **larger range**

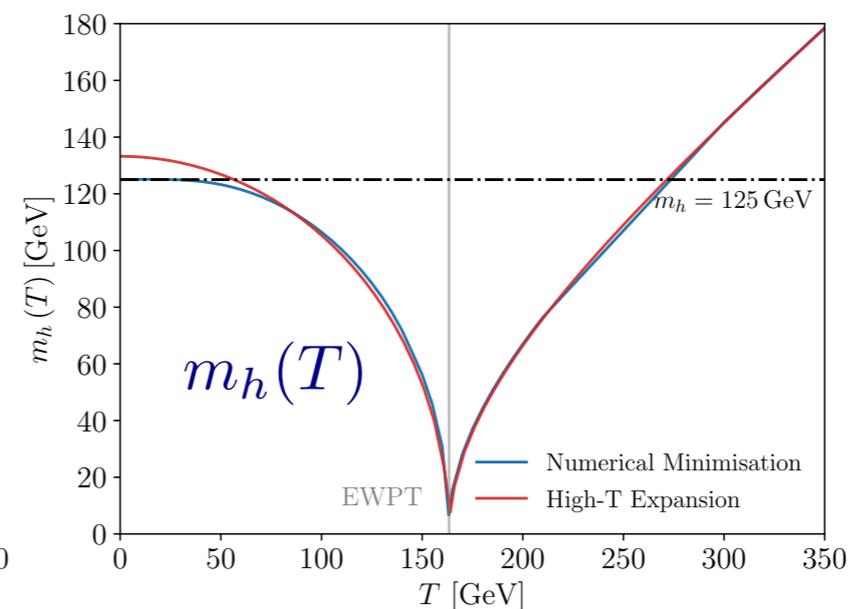
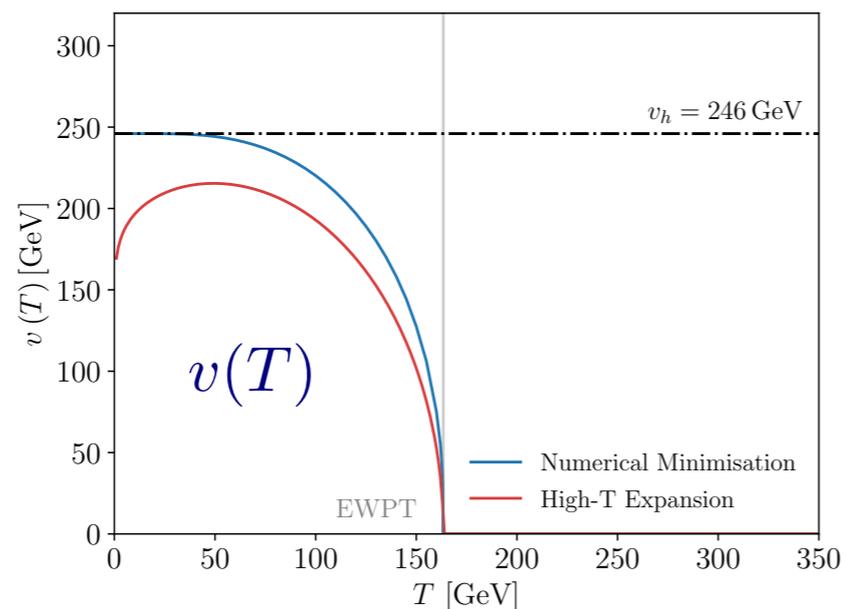
➔ Not only quantum statistics matters, but also

- Thermal masses** of heat bath particles $\delta m^2 \sim g^2 T^2$

- Phase transitions:**

- Potentially *additional* $m(T)$
- Spectrum of states changes

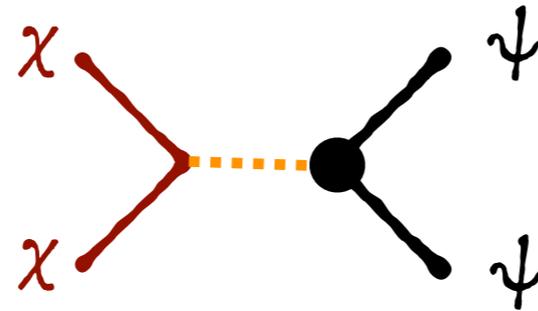
- For example the EWPT:



DM from decays

• 2 → 2 or 1 → 2 ?

• Expect difference only at NLO !



• For a **fully thermalized** mediator **A**, this requires

$$\Gamma_{\text{BW}} = \frac{1}{1 + f_A(\omega')} \sum_{\psi_1 \psi_2} \Gamma_{\psi_1 \psi_2} \bar{G}_{\psi_1 \psi_2}(\gamma, m_A^2)$$

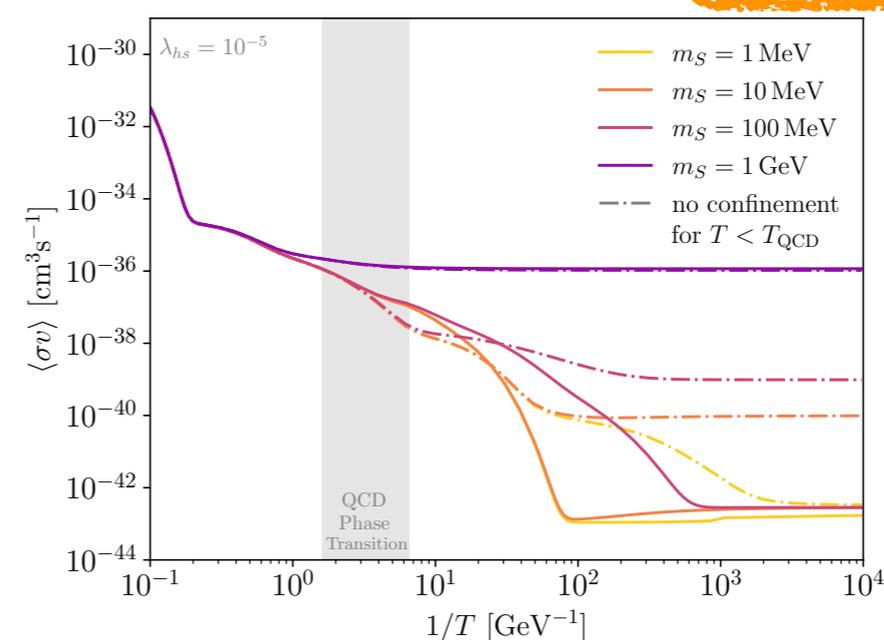
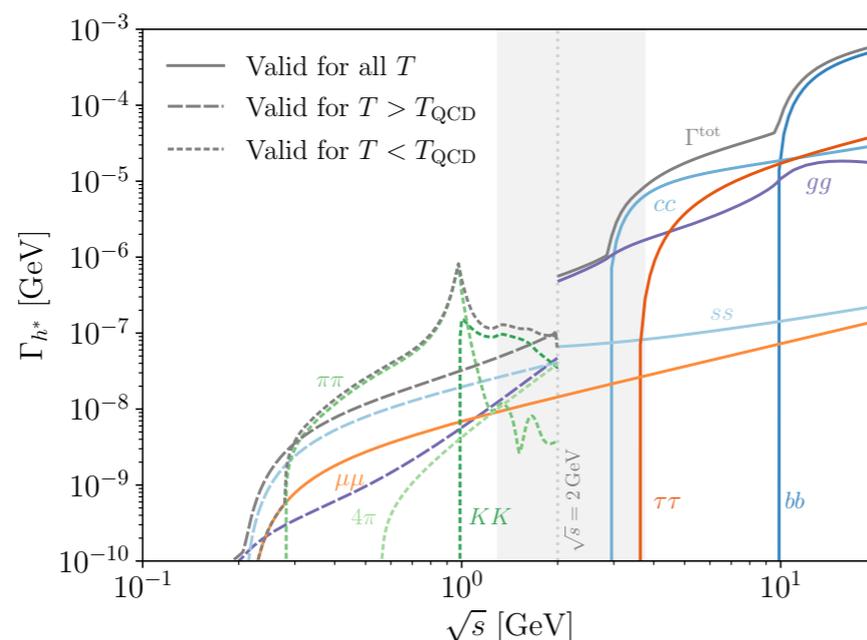
TB, Heeba, Kahlhoefer & Vangsnes, JHEP '22

width in propagator

width in vacuum

final state quantum enhancement/suppression

• For a **SM Higgs** mediator, e.g., we can directly use tabulated partial widths, both below and above QCD PT **for “→”, but not for “←”!**



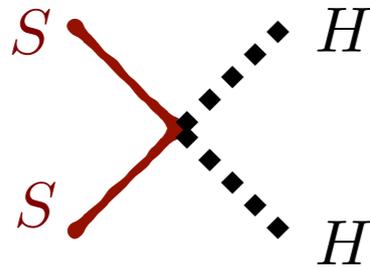
Freeze-in: Case study

Scalar Singlet DM

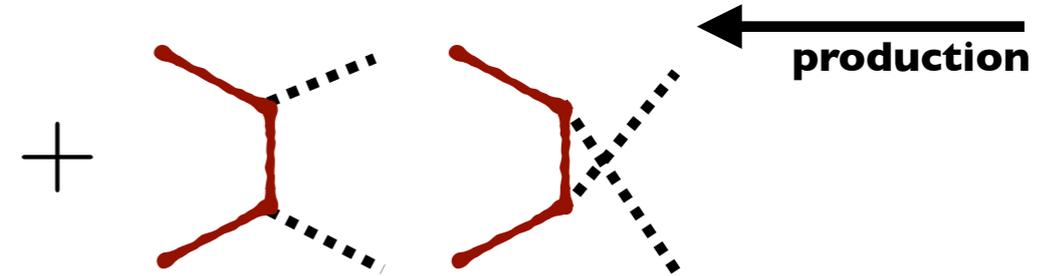
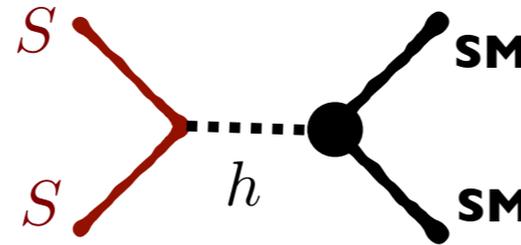
$$\mathcal{L} = \frac{1}{2} \partial_\mu S \partial^\mu S + \frac{1}{2} \mu_S^2 S^2 + \frac{1}{2} \lambda_{hs} S^2 |H|^2 + \frac{1}{4} \lambda_s S^4$$

Silveira & Zee, PLB '85

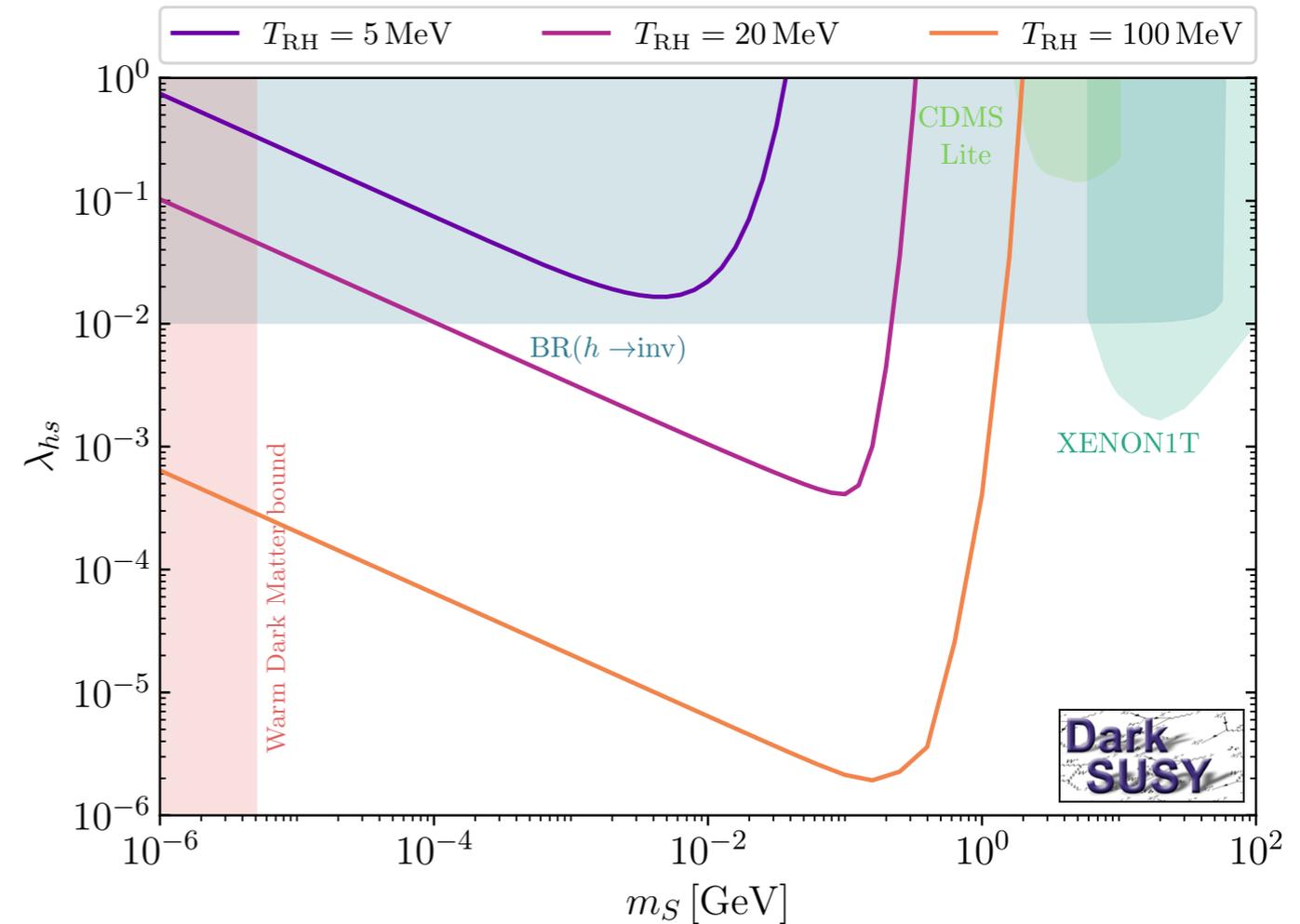
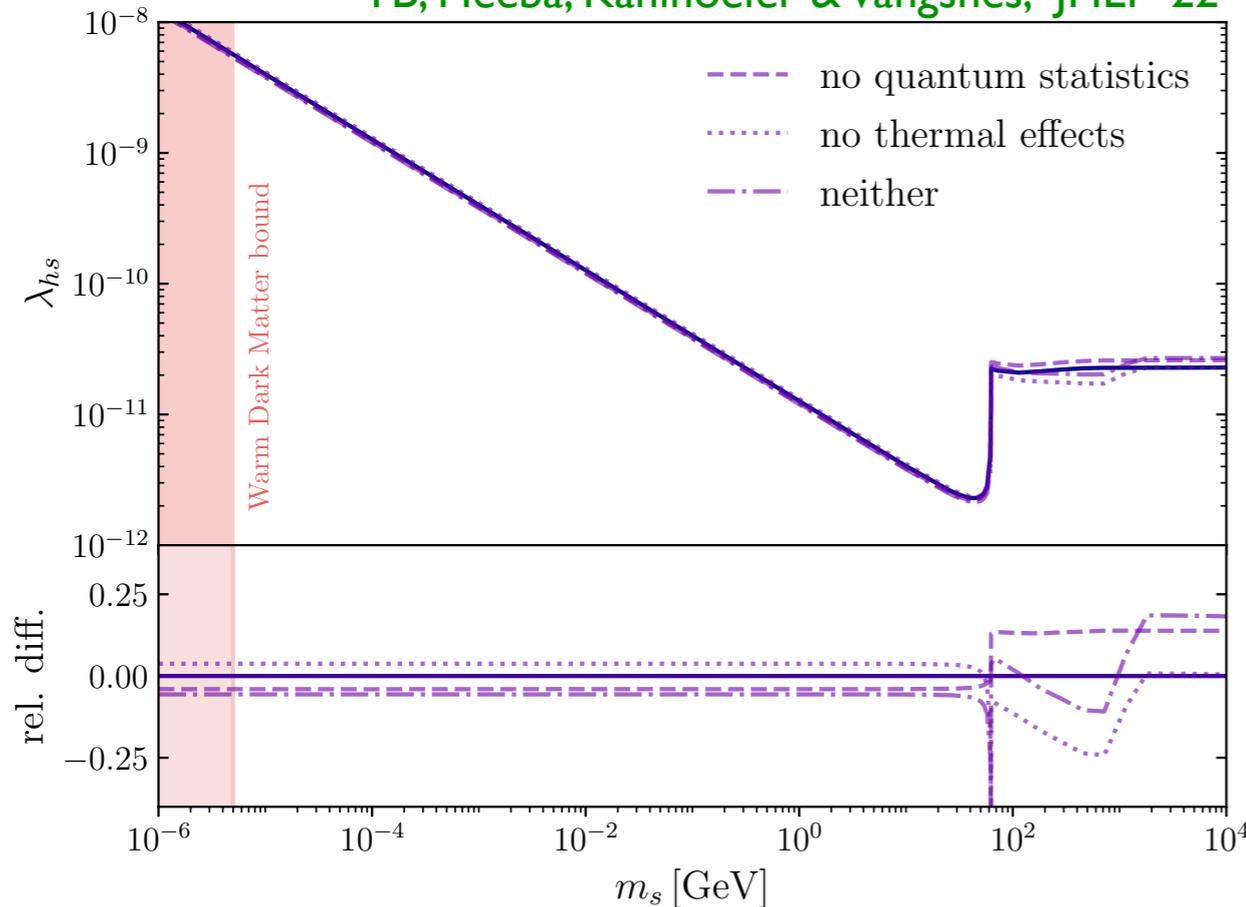
before EWSB:



after EWSB:



TB, Heeba, Kahlhoefer & Vangsnes, JHEP '22



High reheating temperature

Low reheating temperature

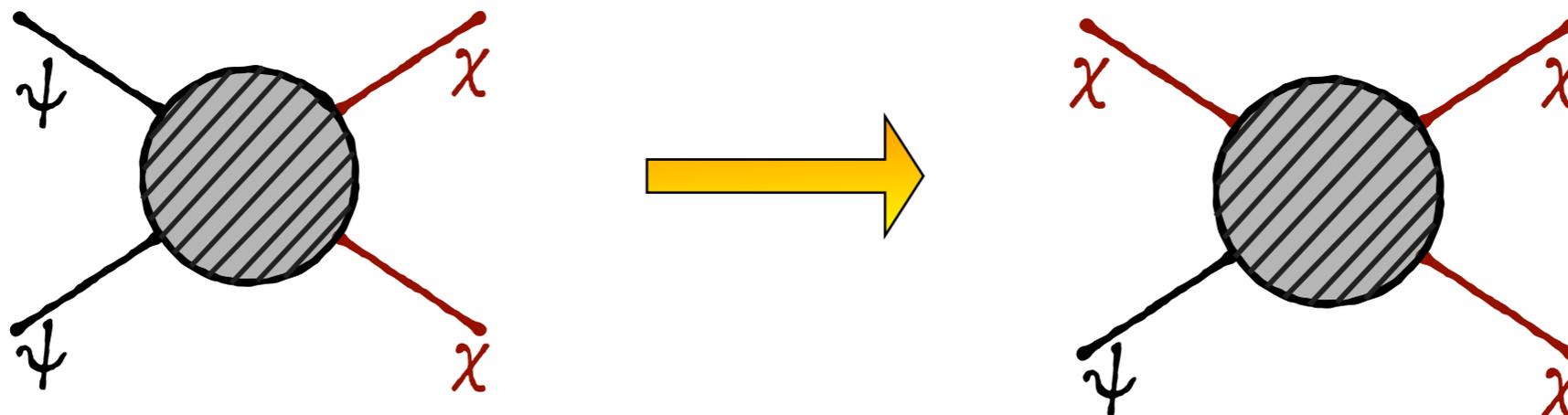


A new production mechanism

- In the past, **many variants** of freeze-out scenario have been explored:

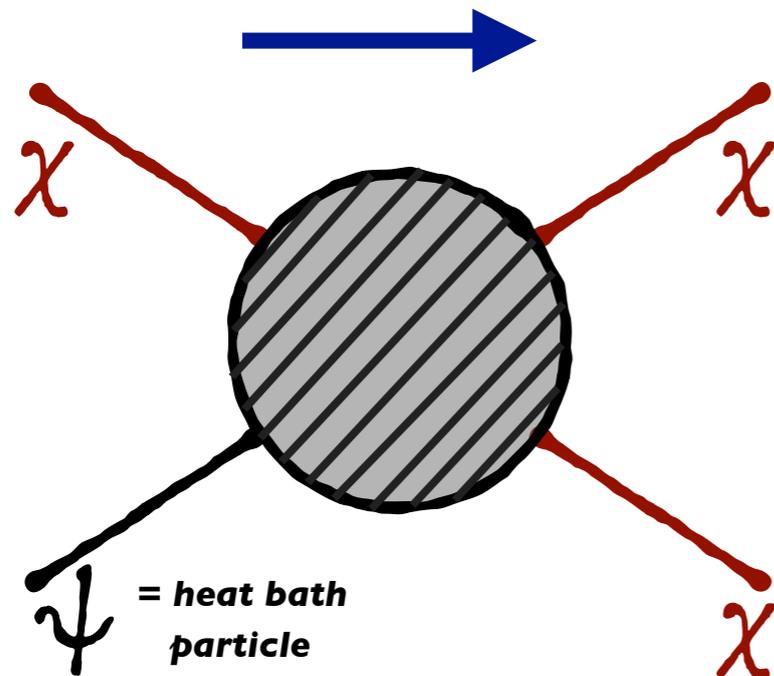


- Can we get a **qualitatively different** freeze-in scenario?



A new production mechanism

• 'Pandemic' dark matter



TB, Depta, Hufnagel, Rudermann
& Schmidt-Hoberg, 2103.16572
Hryczuk & Laletin, 2104.05684

$$\dot{n}_\chi + 3H n_\chi = n_\chi n_\psi^{\text{eq}} \langle \sigma v \rangle$$

[for $n_\chi \ll n_\psi^{\text{eq}}$]

• The 'SIR' compartmental model

A Contribution to the Mathematical Theory of Epidemics.
By W. O. KERMACK and A. G. MCKENDRICK.
(Communicated by Sir Gilbert Walker, F.R.S.—Received May 13, 1927.)

S # susceptible individuals
 I # infected individuals
recovered ($R = \text{tot} - S - I$)

β infection rate
 γ recovery rate

$$\dot{I} = \beta S I - \gamma I$$

→ reproduction number, or 'R-value':

$$R \equiv \frac{\beta S}{\gamma} = \frac{n_\psi^{\text{eq}} \langle \sigma v \rangle}{3H}$$

A new production mechanism

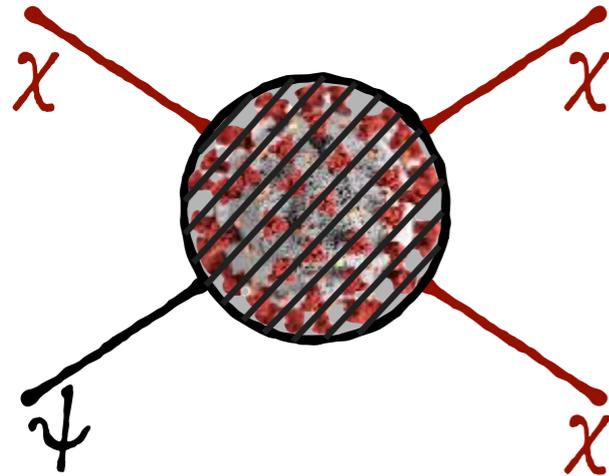


FIG. 1. The transformation process leading to exponential production of DM (χ) from the heat bath (ψ).

Pandemic Dark Matter

Torsten Bringmann,^{1,*} Paul Frederik Depta,^{2,†} Marco Hufnagel,^{3,‡}
Joshua T. Ruderman,^{4,2,5,§} and Kai Schmidt-Hoberg^{2,¶}

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³*Service de Physique Théorique, Université Libre de Bruxelles, Boulevard du Triomphe, CP225, B-1050 Brussels, Belgium*

⁴*Center for Cosmology and Particle Physics, Department of Physics, New York University, New York, NY 10003, USA*

⁵*Kavli Institute for Theoretical Physics, University of California, Santa Barbara, CA 93106, USA*

(Date: April 1st, 2021)

arXiv: 2103.16572

PRL wasn't very happy... 🙄

Exponential DM production

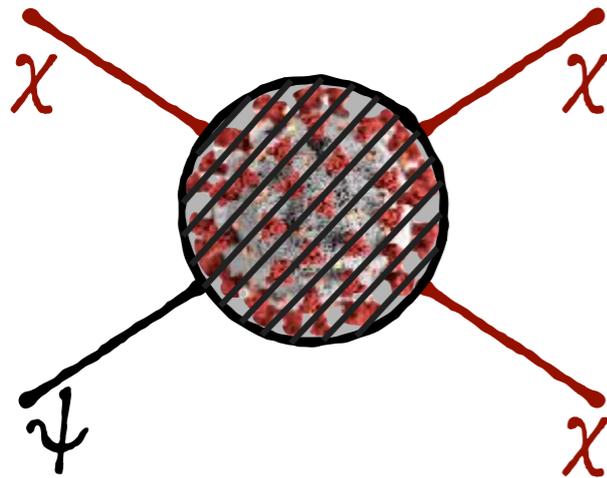


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arXiv:2103.16572



PHYSICAL REVIEW LETTERS **127**, 191802 (2021)

Editors' Suggestion

Featured in Physics

Dark Matter from Exponential Growth

Torsten Bringmann^{1,*} Paul Frederik Depta^{2,†} Marco Hufnagel,^{3,‡}
Joshua T. Ruderman,^{4,2,5,6,§} and Kai Schmidt-Hoberg^{2,||}

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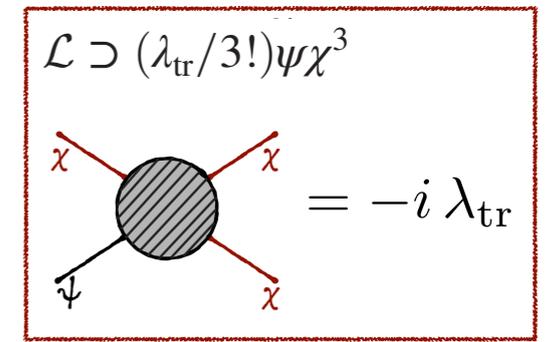
⁵Kavli Institute for Theoretical Physics, University of California, Santa Barbara, California 93106, USA

⁶School of Physics and Astronomy, Tel-Aviv University, Tel-Aviv 69978, Israel

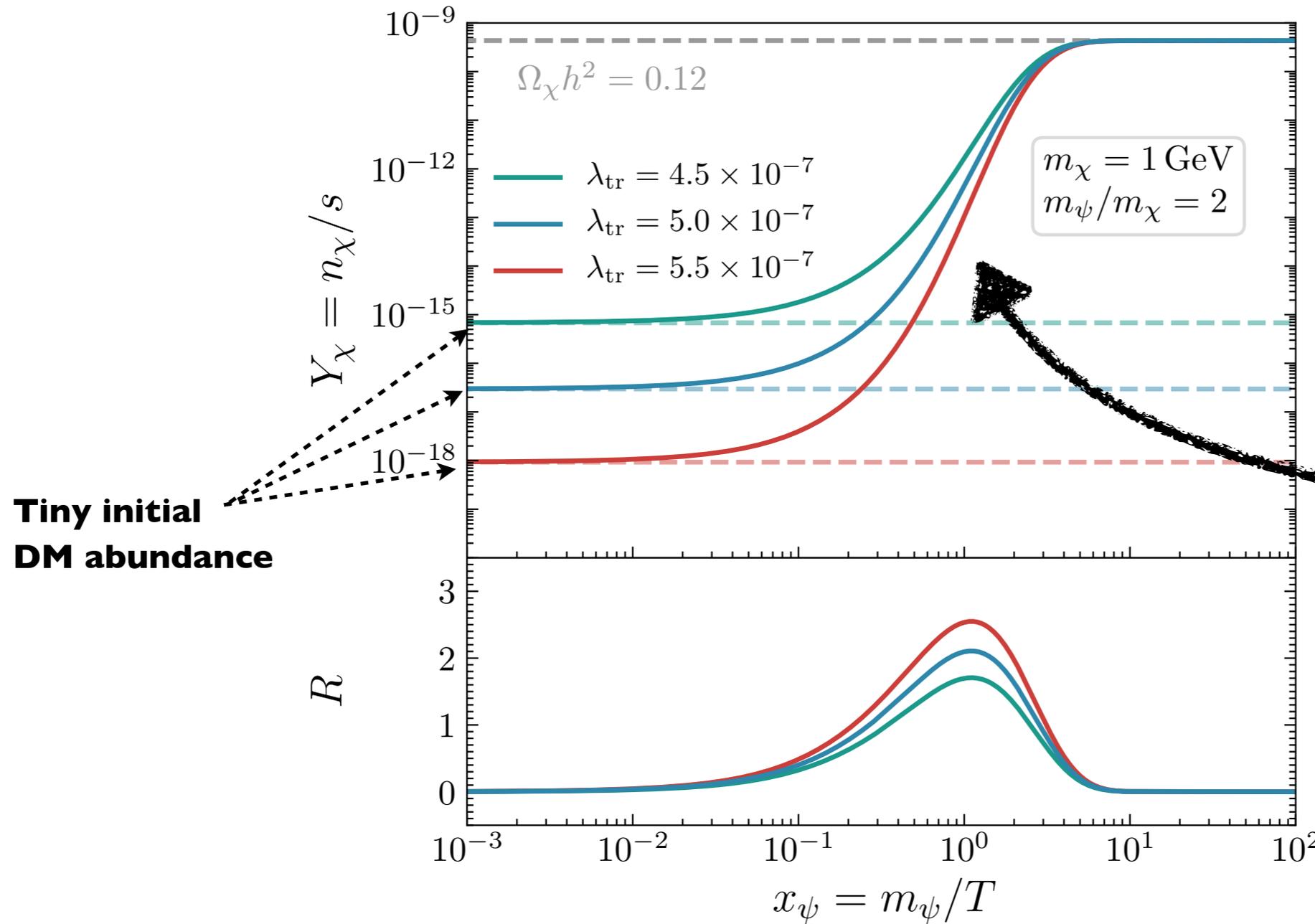
(Received 25 June 2021; revised 20 August 2021; accepted 14 September 2021; published 3 November 2021)



Exponential DM production



toy model



Tiny initial
DM abundance

exponential
growth

$$R \gtrsim 1$$

← large H

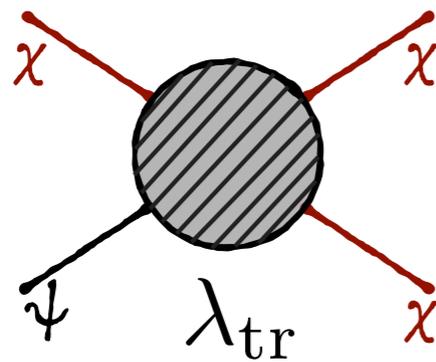
→ small n_{ψ}^{eq}

$$R \equiv \frac{\beta S}{\gamma} = \frac{n_{\psi}^{\text{eq}} \langle \sigma v \rangle}{3H}$$



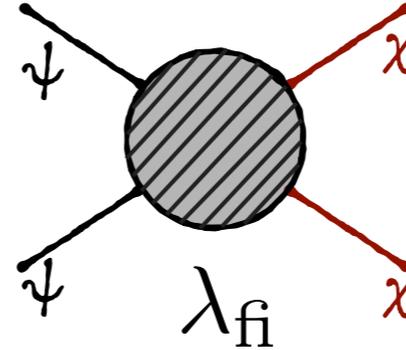
Adding freeze-in production

'transmission'

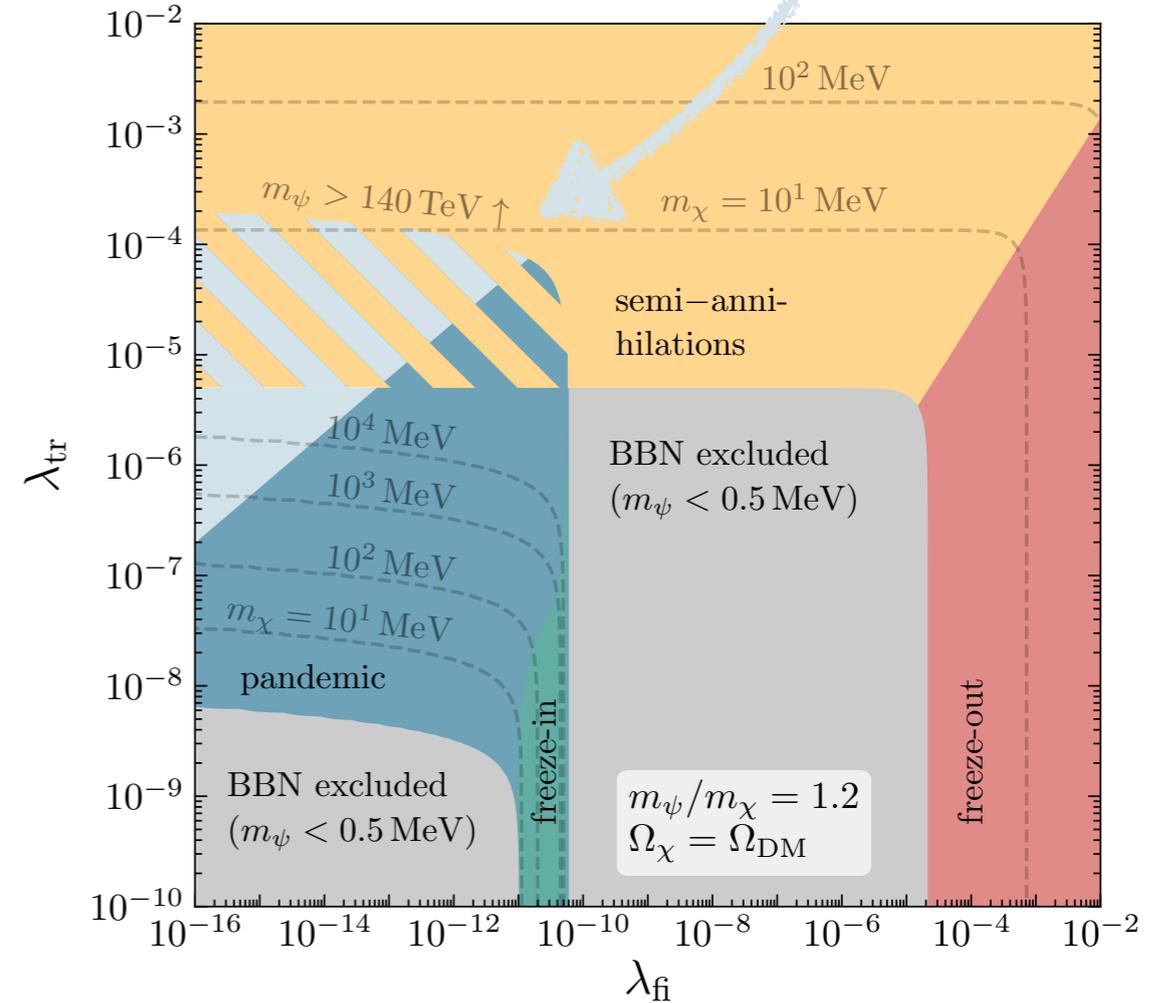
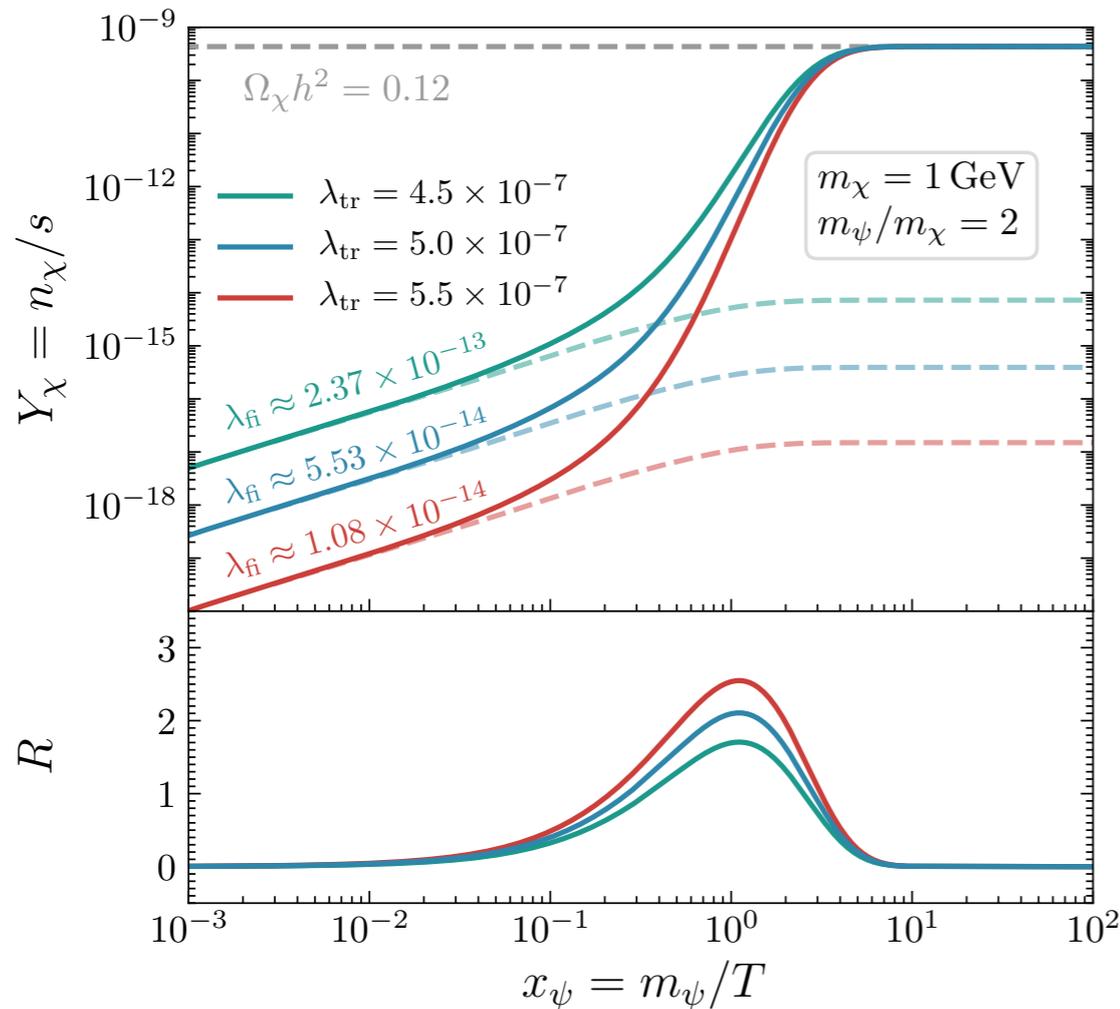


+

'freeze-in'



2→4
freeze-in
dominates



➔ 'Pandemic' production turns out to be a rather generic mechanism for the genesis of DM!

Signals ?

- Necessarily model-dependent

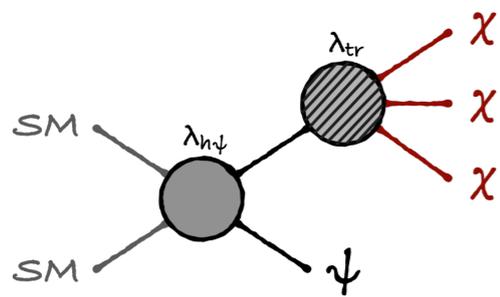
- 'Pandemic DM' describes a **class** of models, just like 'WIMP' does

- Example:** Higgs portal coupling

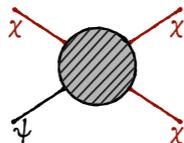
→ ψ is itself a thermal relic (freeze-out!)
= subdominant DM component

$$\mathcal{L} \supset \frac{1}{2} \lambda_{h\psi} |H|^2 \psi^2$$

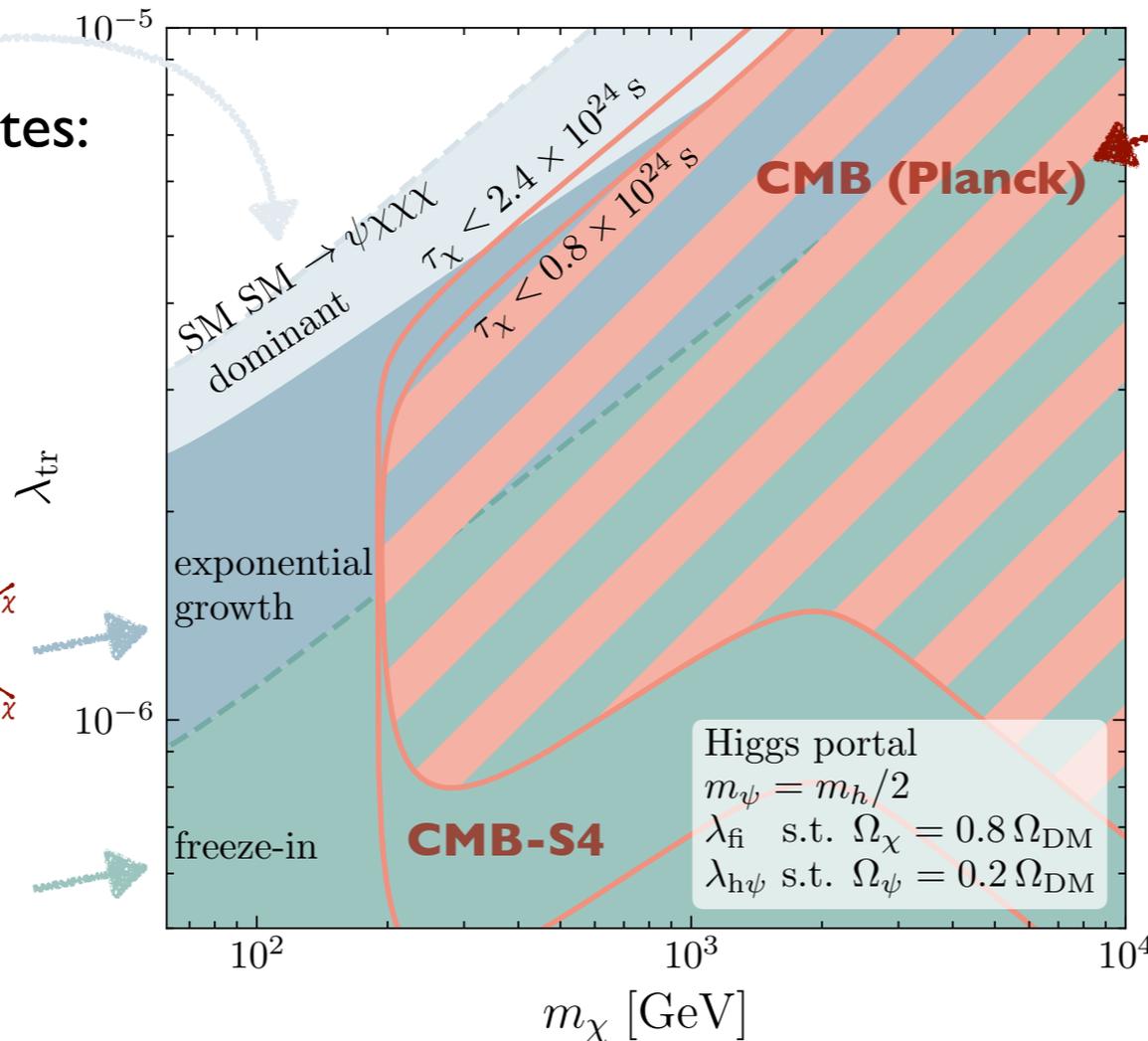
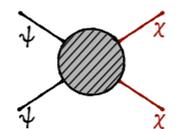
- 2→4 freeze-in production dominates:



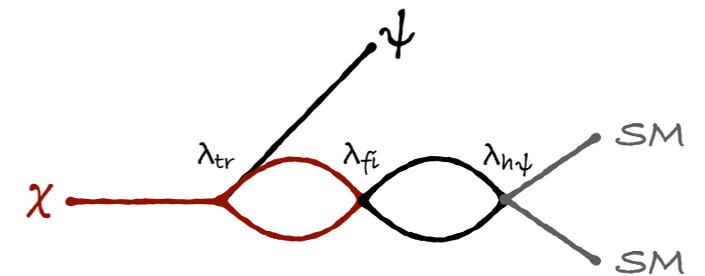
- exponential production



- Freeze-in production



- indirect detection signals from **DM decay**:

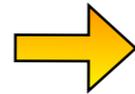


Further model building

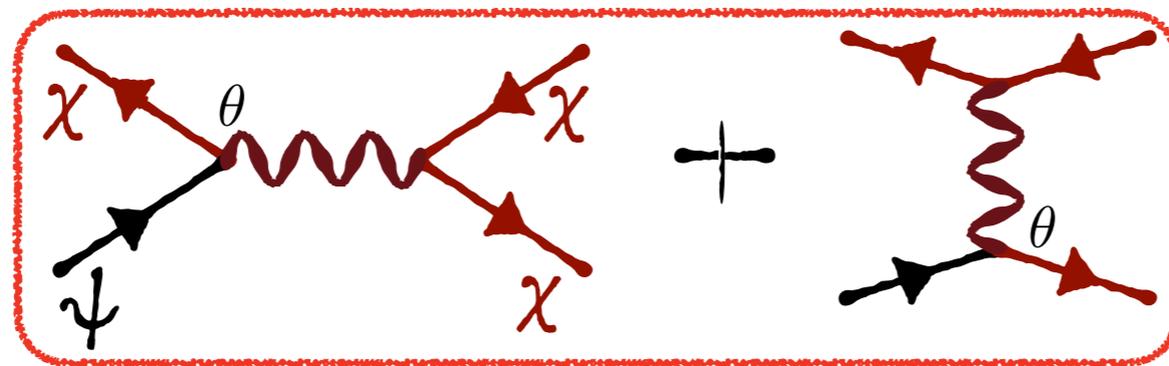
• Q: Is there a generic way to realize $\langle \sigma v \rangle_{\text{fi}} \ll \langle \sigma v \rangle_{\text{tr}}$?

• A: yes — just add a **dark sector mediator** and **mass mixing!**

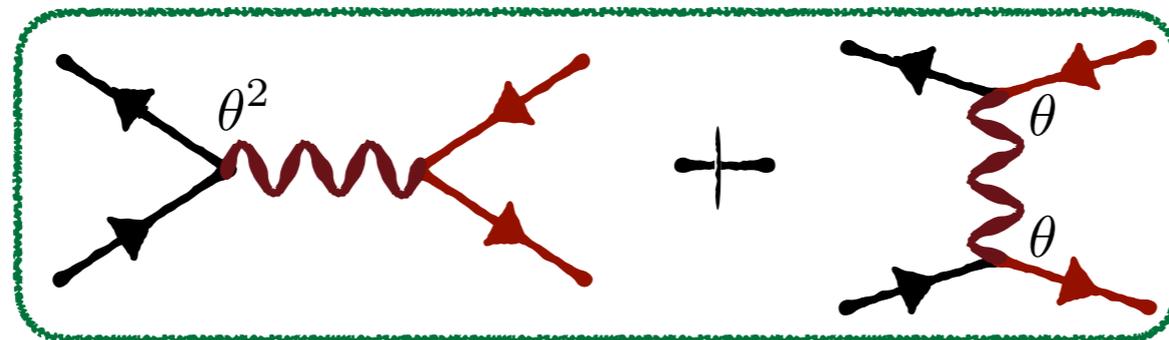
$$\mathcal{L} \supset -\delta m (\bar{\psi}\chi + \bar{\chi}\psi) - g\bar{\chi}V\chi$$



$$\mathcal{L} \supset -g[\bar{\chi}V\chi + \theta (\bar{\psi}V\chi + \bar{\chi}V\psi) + \theta^2\bar{\psi}V\psi]$$



transmission $\propto \theta$



freeze-in $\propto \theta^2$

Sterile neutrinos

- An **excellent**, well-motivated dark matter **candidate**
- **Production by SM processes: oscillations** with active neutrinos, combined with CC and NC scatterings

Dodelson &
Widrow, PRL '94

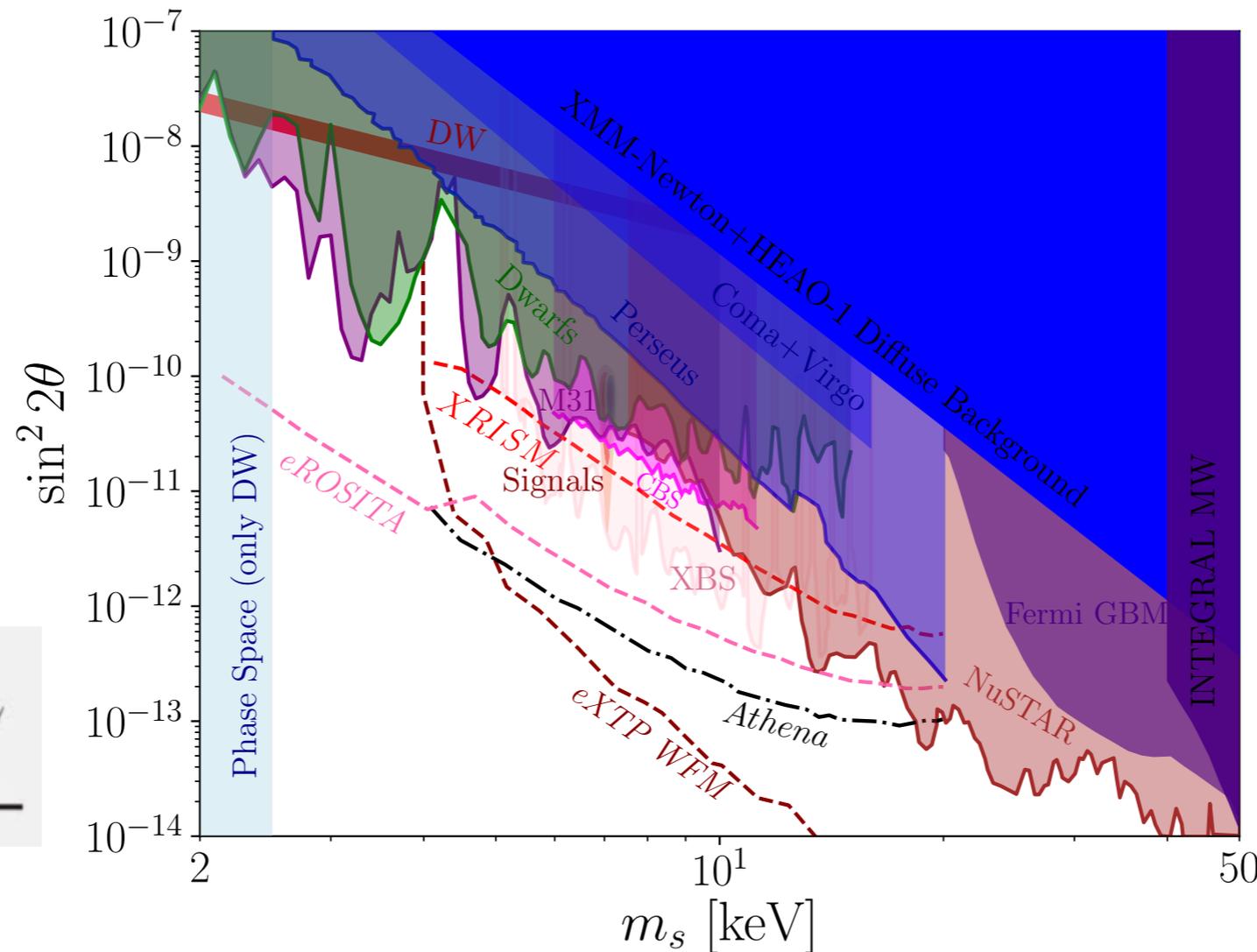
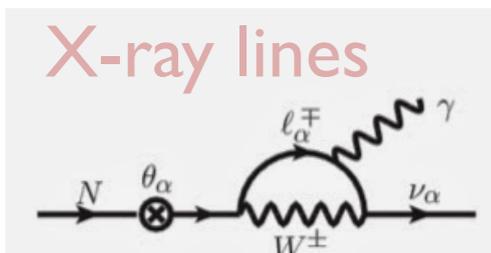


Fig.: Abazajian+, 2203.7377

- Unfortunately, this scenario is **ruled out** by observations...



Sterile neutrinos — a new life ?

TB, Depta, Hufnagel, Kersten, Ruderman & Schmidt-Hoberg, arXiv:2206.10630

- Simply add a **scalar** ϕ that only couples to the sterile neutrinos

$$\mathcal{L} \supset \frac{y}{2} \phi \bar{\nu}_s \nu_s \quad \longrightarrow \quad \frac{y}{2} \phi [\sin^2 \theta \bar{\nu}_\alpha \nu_\alpha - \sin \theta \cos \theta (\bar{\nu}_\alpha \nu_s + \bar{\nu}_s \nu_\alpha) + \cos^2 \theta \bar{\nu}_s \nu_s]$$

$$m_\phi > 2m_s$$

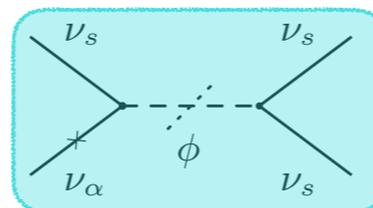
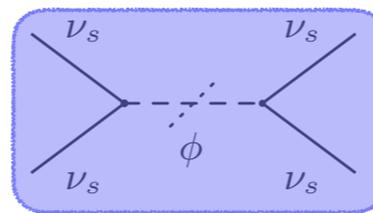
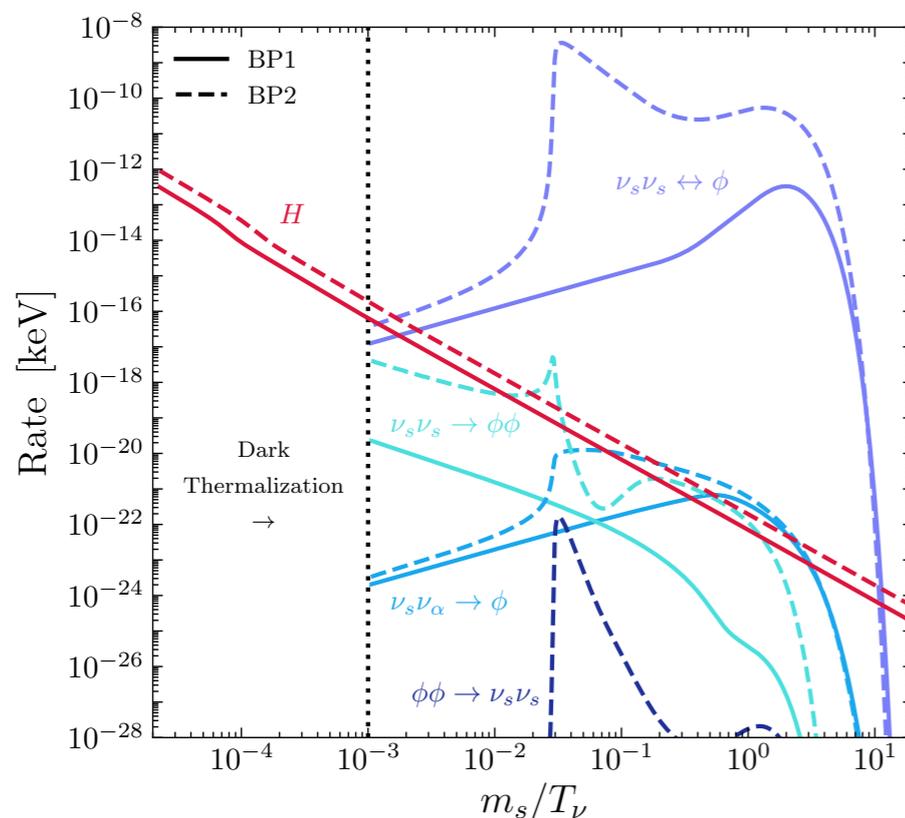
- Early times (\sim QCD PT): standard **DW** production

- Adopt resulting number & energy density as **initial condition**

from Asaka, Laine & Shaposhnikov, JHEP '15

- Soon afterwards: efficient **dark sector thermalization**

- ν_s, ϕ follow FD/BE distributions with *large* (negative) chemical potentials



Use Boltzmann equations

$$\dot{n}_s + 3Hn_s = C_{n_s}$$

$$\dot{n}_\phi + 3Hn_\phi = C_{n_\phi}$$

$$\dot{\rho} + 3H(\rho + P) = C_\rho,$$

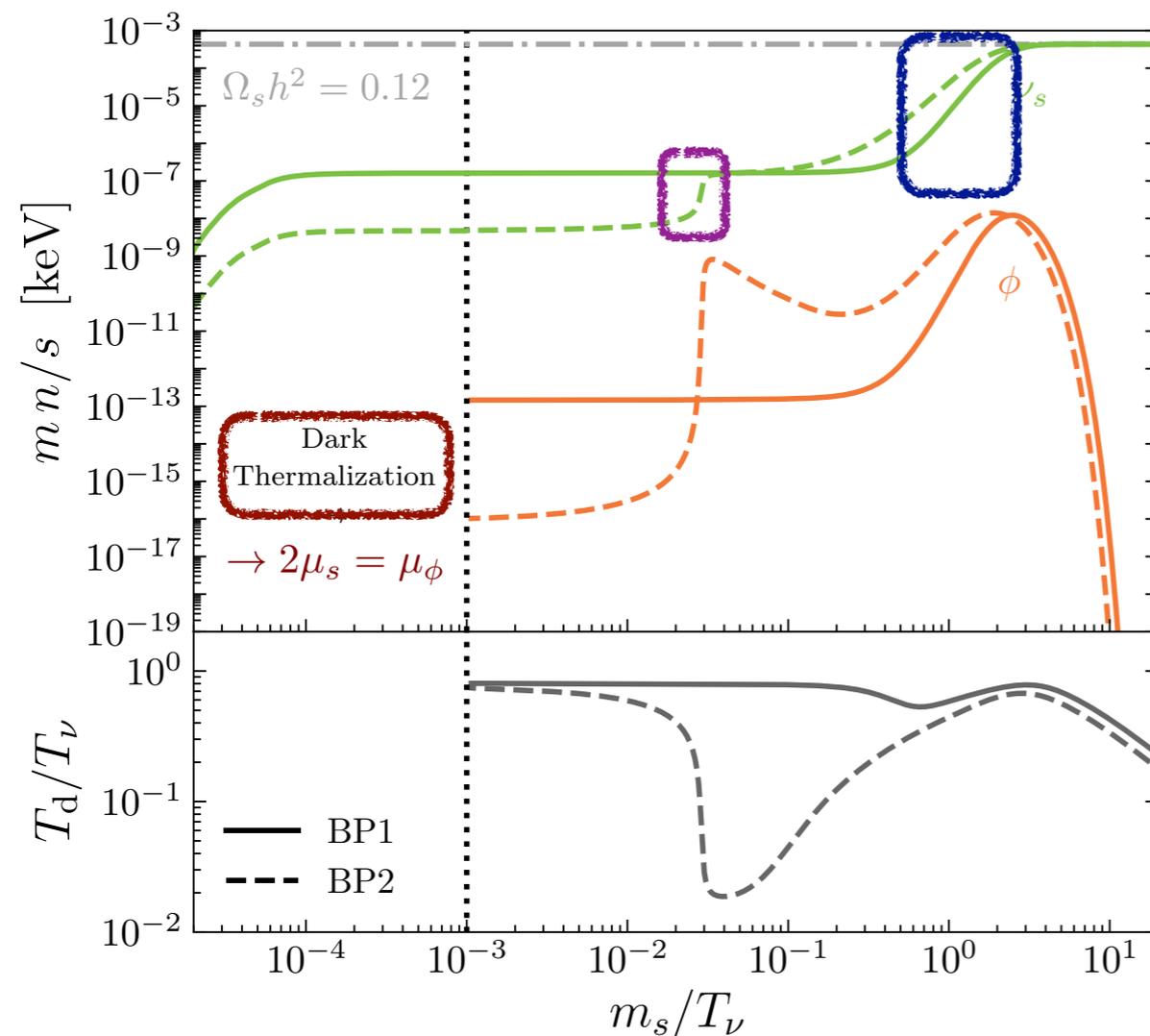
to solve for $T_d(T), \mu_s(T), \mu_\phi(T)$

Evolution and constraints

TB, Depta, Hufnagel, Kersten, Ruderman & Schmidt-Hoberg, arXiv:2206.10630

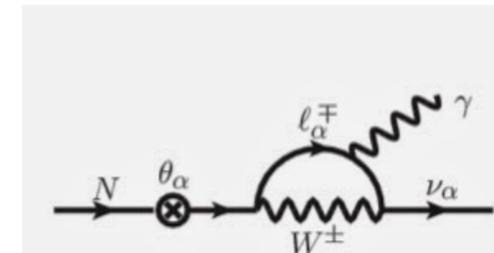
Two benchmark points

- Solid: 'large' θ , small γ
- Dashed: smaller θ , larger γ



Observational constraints

- (Standard) X-ray lines



- ν_s self-interactions

$$\sigma_T/m_s \lesssim 1 \text{ cm}^2/\text{g}$$

cf. Tulin & Yu, PR '18

maybe 0.1 possible... (?)

- Lyman- α

recast $m_{\text{WDM}} > 1.9 \text{ keV}$ to

Garzilli+, MNRAS '21

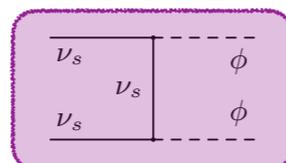
$$\lambda_{\text{FS}} < 0.24 \text{ Mpc}$$

$$r_s < 0.36 \text{ Mpc}$$

maybe $m_{\text{WDM}} > 5.3 \text{ keV}$ possible... (?)

Palanque-Delabrouille+, JCAP '20

- 3 phases (after DW): dark thermalization



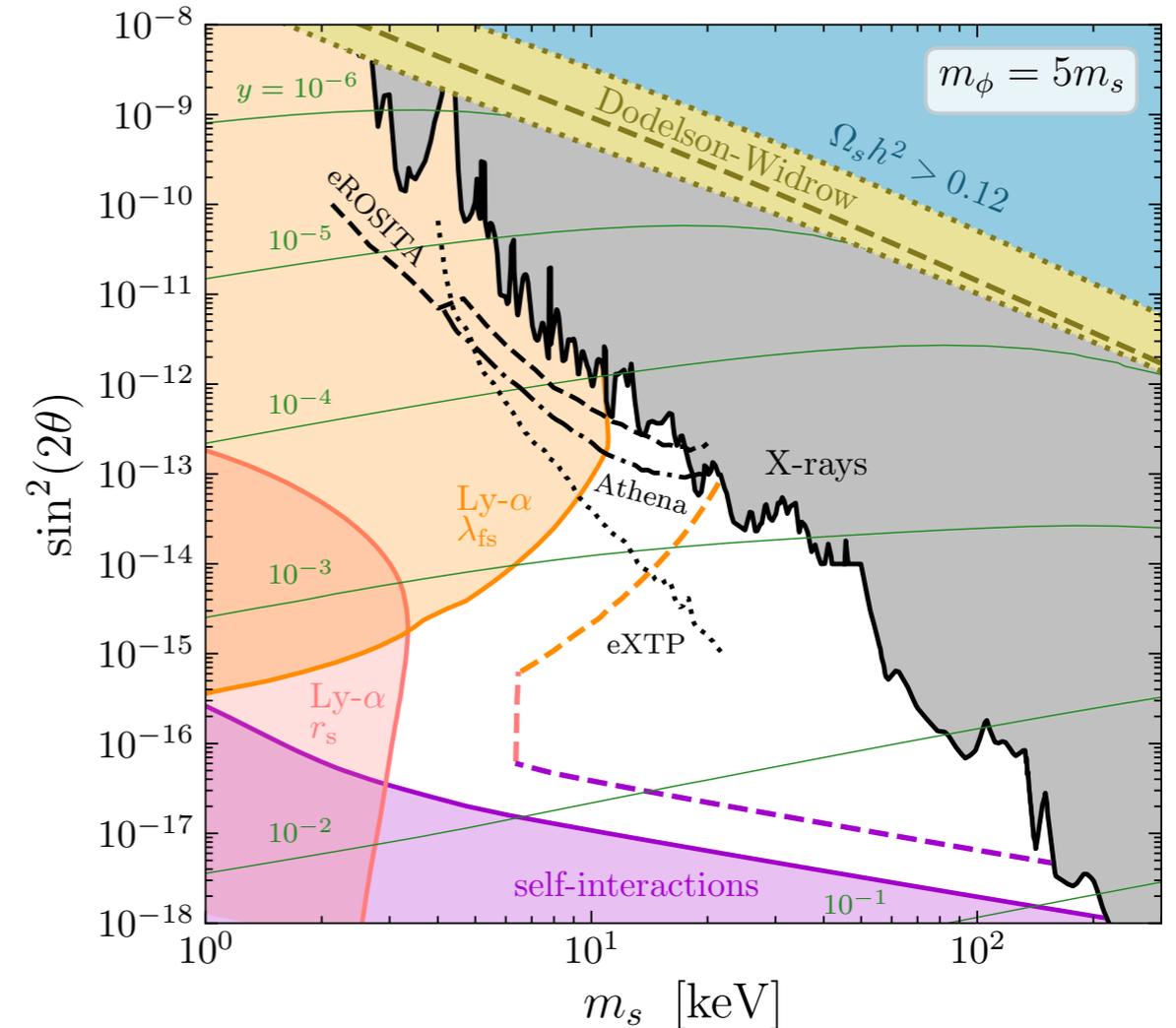
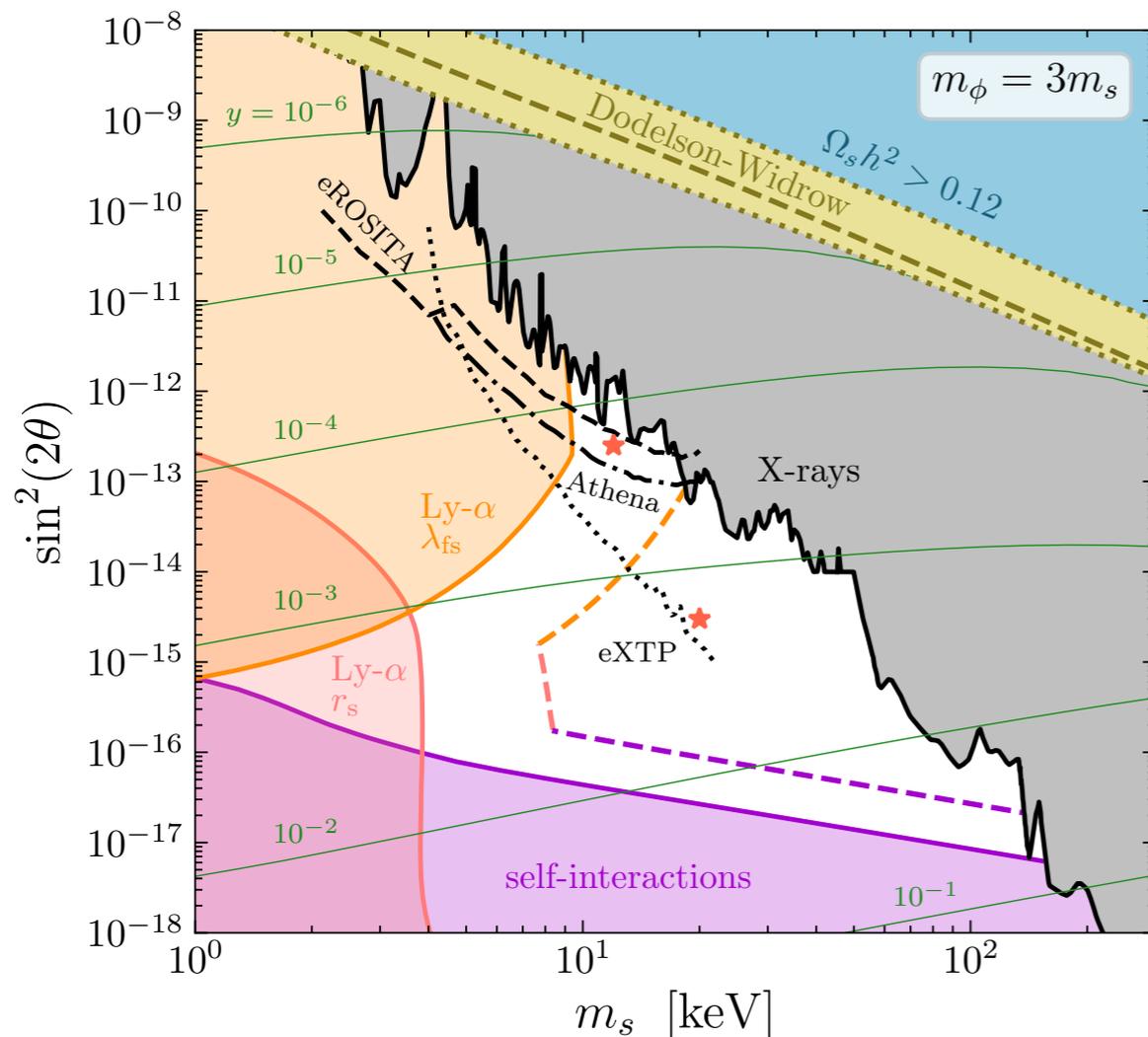
exponential growth
reproductive freeze-in

March-Russell+, PRD '20 ; Mondino+, PRD '21

Sterile neutrinos — a new life !

TB, Depta, Hufnagel, Kersten, Ruderman & Schmidt-Hoberg, arXiv:2206.10630

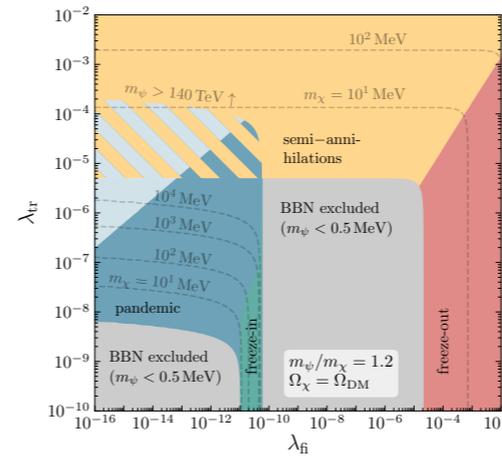
- **New parameter space** below DW line
- Bounded from above *and* below
- Significant parts are (or will be) in **observational reach**



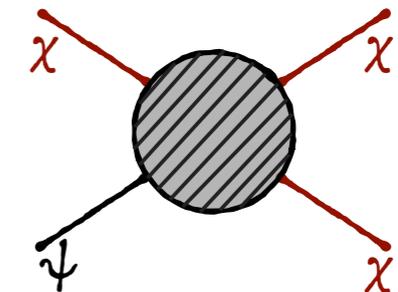
- Larger **mediator masses** widen possible window for sterile neutrino dark matter

Conclusions

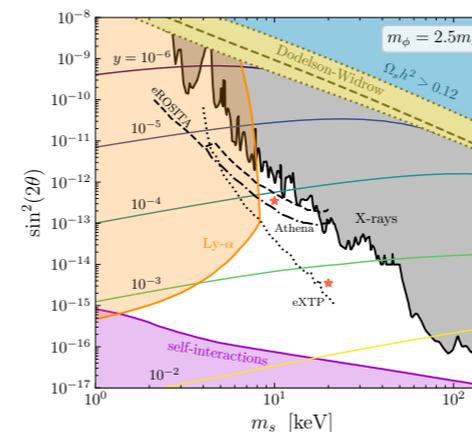
- There are a handful of **generic ways of DM genesis** from the primordial heat bath



- Pandemic dark matter is a novel such mechanism. **Mathematical analogy** to spread of diseases works almost scarily well



- Maybe this even brings new life to **sterile neutrino DM** ?!



Thanks for your attention!

DarkSUSY



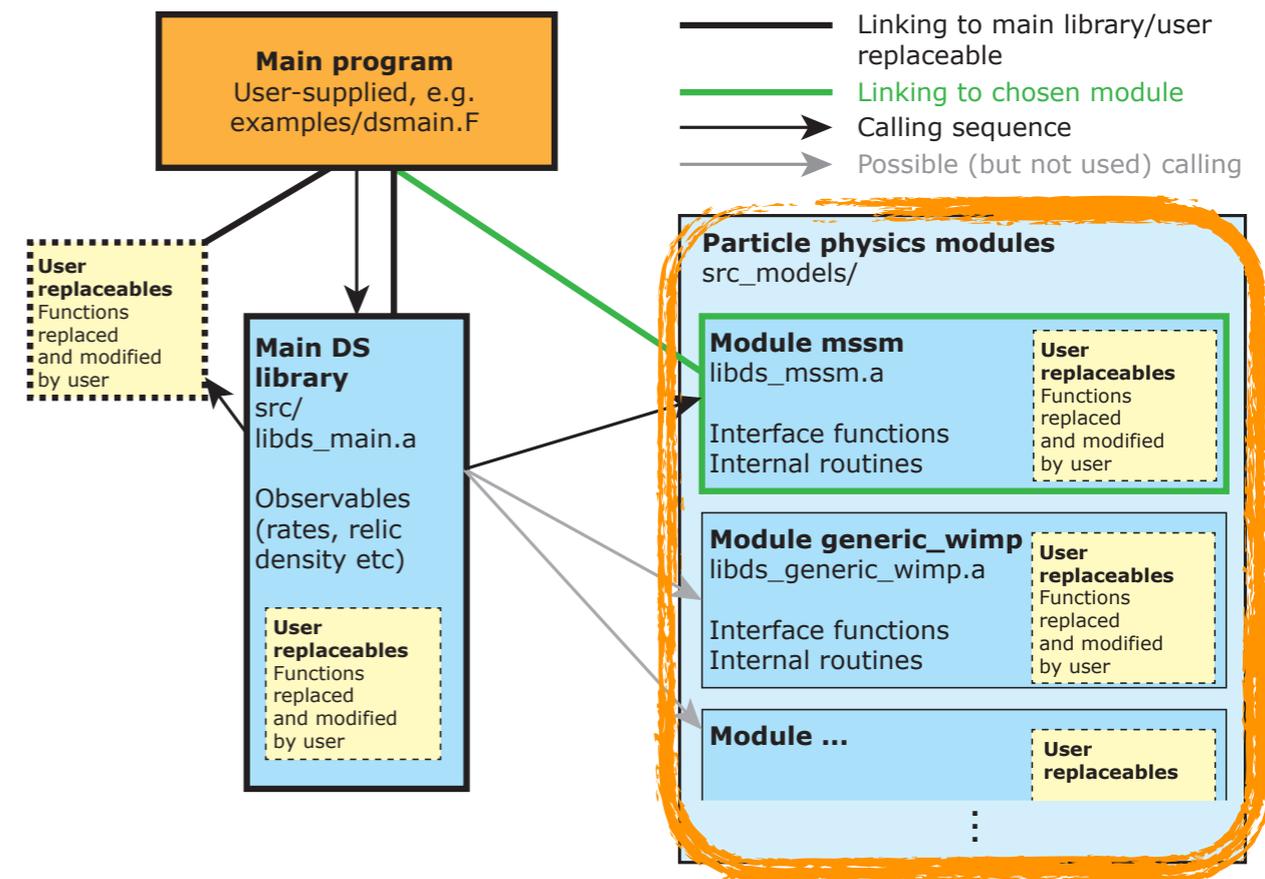
TB, Edsjö, Gondolo,
Ullio & Bergström,
JCAP '18

[http://
darksusy.hepforge.org](http://darksusy.hepforge.org)

**Since version 6:
no longer restricted to
supersymmetric DM !**

Numerical package to calculate 'all' DM related quantities:

- relic density + kinetic decoupling
(also for $T_{\text{dark}} \neq T_{\text{photon}}$)
- generic SUSY models + laboratory constraints implemented
- cosmic ray propagation
- particle yields for generic DM annihilation or decay
- indirect detection rates: gammas, positrons, antiprotons, neutrinos
- direct detection rates
- ...



- since 6.1: DM self-interactions**
- since 6.2: 'reverse' direct detection**
(incl. full Q^2 -dependent scattering!)
- since 6.3: freeze-in**