

# Maximal temperature of strongly coupled dark sectors

Helena Kolešová (AEC, ITP, University of Bern)



**Joint work with Simona Procacci and Mikko Laine**  
**ArXiv: [2303.17973](https://arxiv.org/abs/2303.17973)**

Supported by the SNSF under grant 200020B-188712

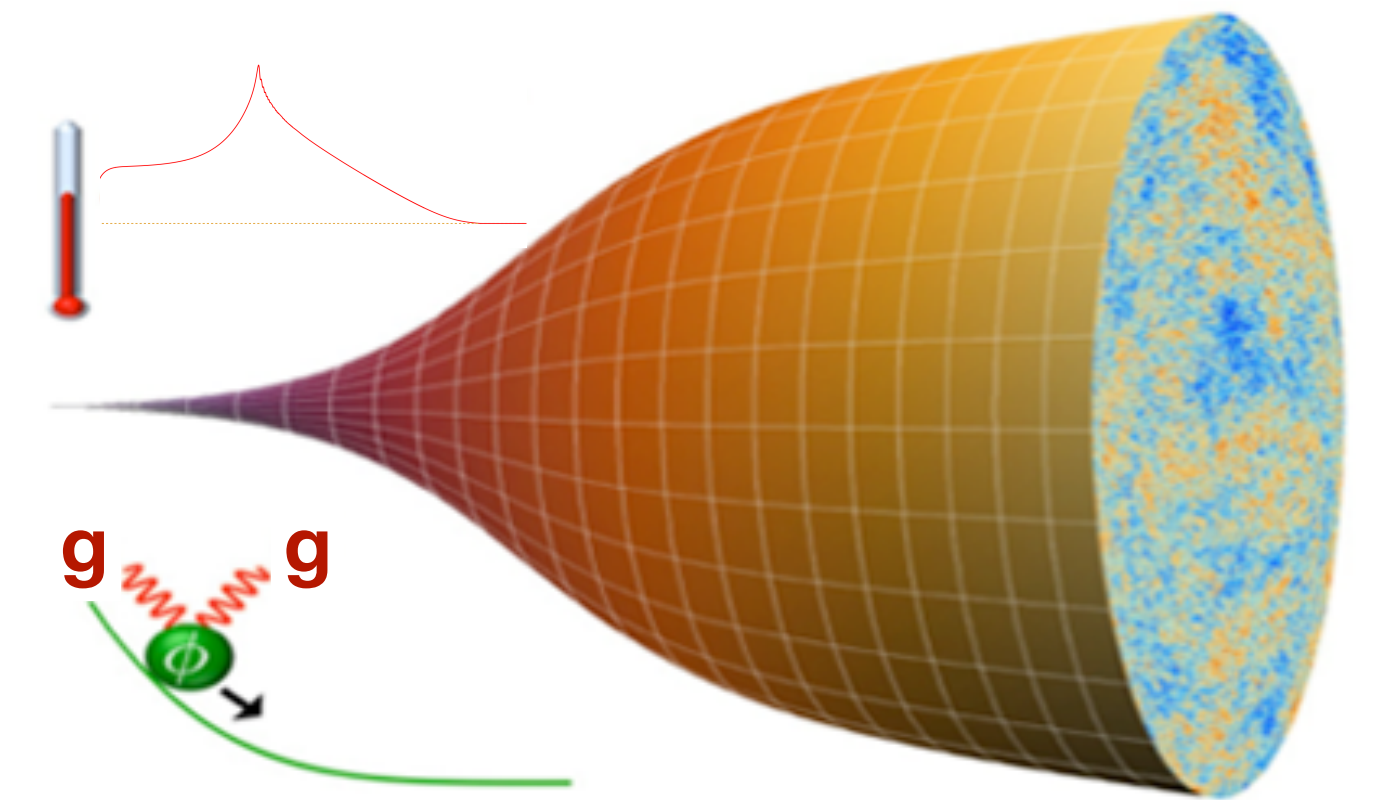
***u*<sup>b</sup>**

---

<sup>b</sup>  
**UNIVERSITÄT  
BERN**

# Summary

- Inflaton coupled to non-abelian dark sector (parametrised by confinement scale  $\Lambda_{\text{IR}}$ )
- Thermal bath present throughout the inflation!
- Working example: Axion inflation coupled to SU(3) gauge sector  
[Laine, Procacci: 2102.09913]  
[Klose, Laine, Procacci: 2201.02317]  
[Klose, Laine, Procacci: 2210.11710]
- [HK, Laine, Procacci: 2303.17973] - evolution of dark sector temperature studied



Credit: João G. Rosa/University of Aveiro; ESA and the Planck collaboration

# Why is high temperature interesting?

- Possible gravitational wave signal:

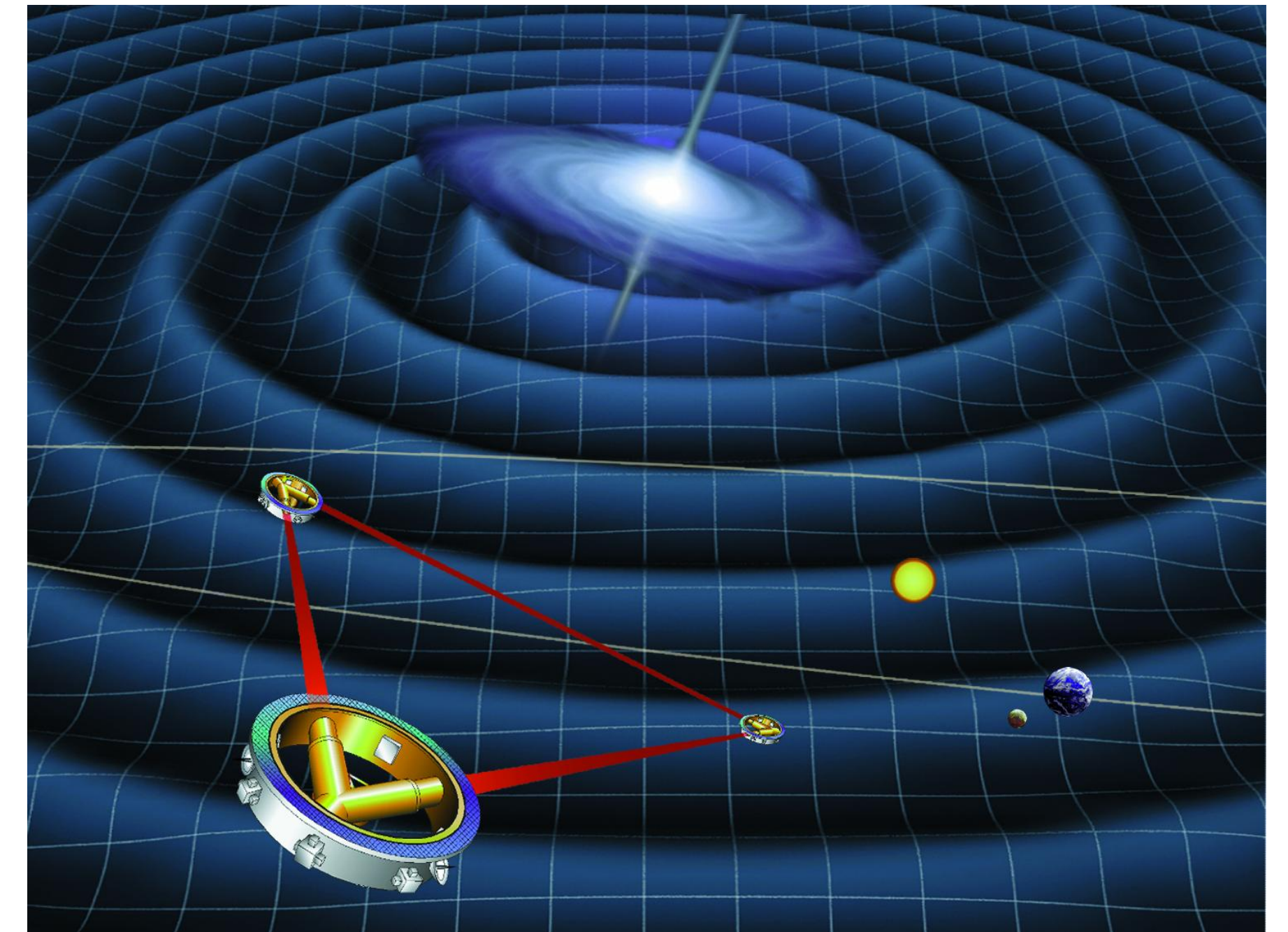
## Thermal fluctuations

[Ghiglieri, Laine: 1504.02569]

[Klose, Laine, Procacci: 2201.02317, 2210.11710]

## Dark sector confinement phase transition

[Caprini et al.: 1910.13125] [Schwaller:1504.07263]



Credit: <http://lisa.jpl.nasa.gov/gallery/lisa-waves.html>

# Why is high temperature interesting?

- Possible gravitational wave signal:

## Thermal fluctuations

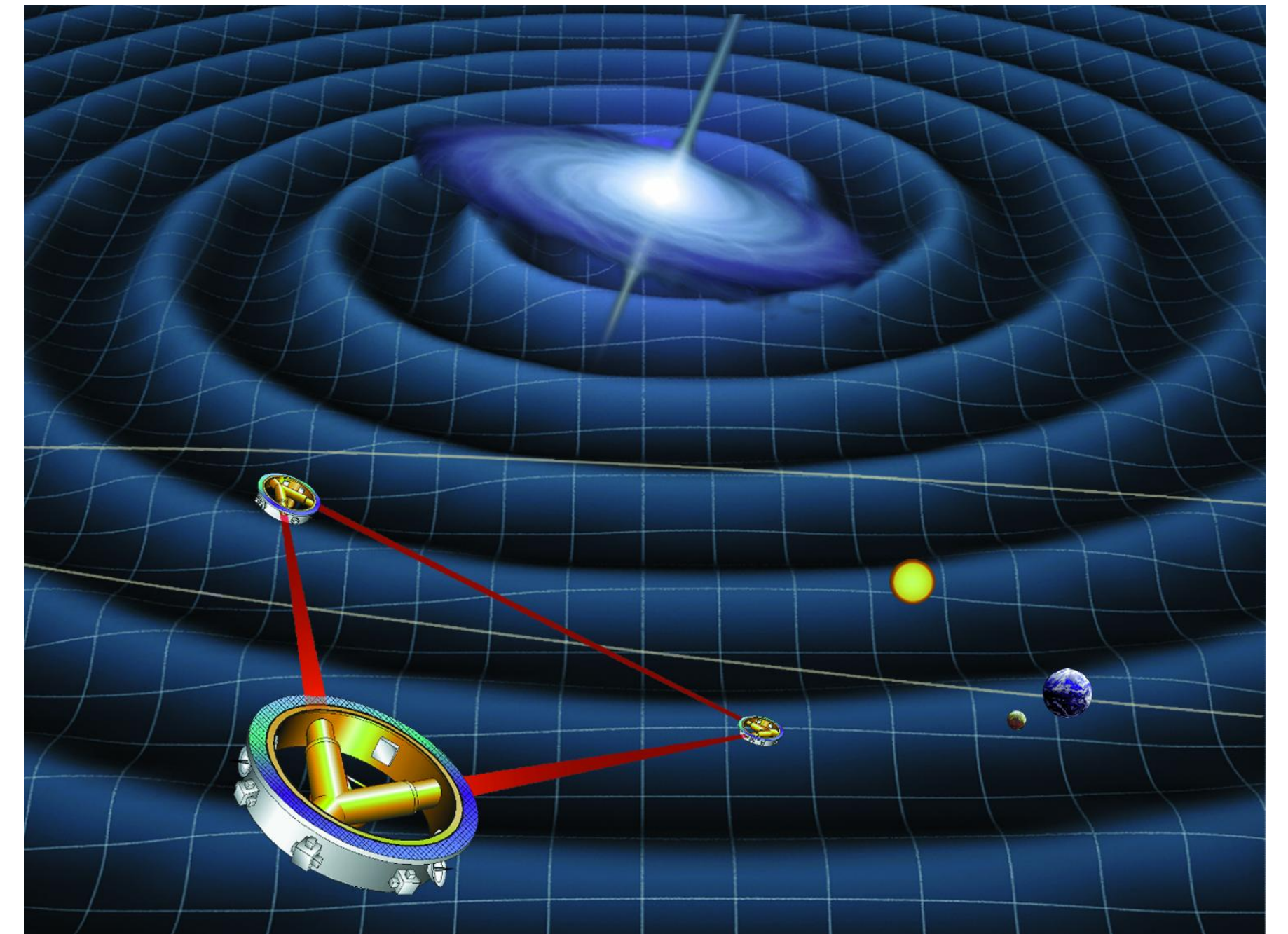
[Ghiglieri, Laine: 1504.02569]

[Klose, Laine, Procacci: 2201.02317, 2210.11710]

## Dark sector confinement phase transition

[Caprini et al.: 1910.13125] [Schwaller:1504.07263]

- But has the dark sector heated up above the confinement temperature?



Credit: <http://lisa.jpl.nasa.gov/gallery/lisa-waves.html>

# Why is high temperature interesting?

- Possible gravitational wave signal:

## Thermal fluctuations

[Ghiglieri, Laine: 1504.02569]

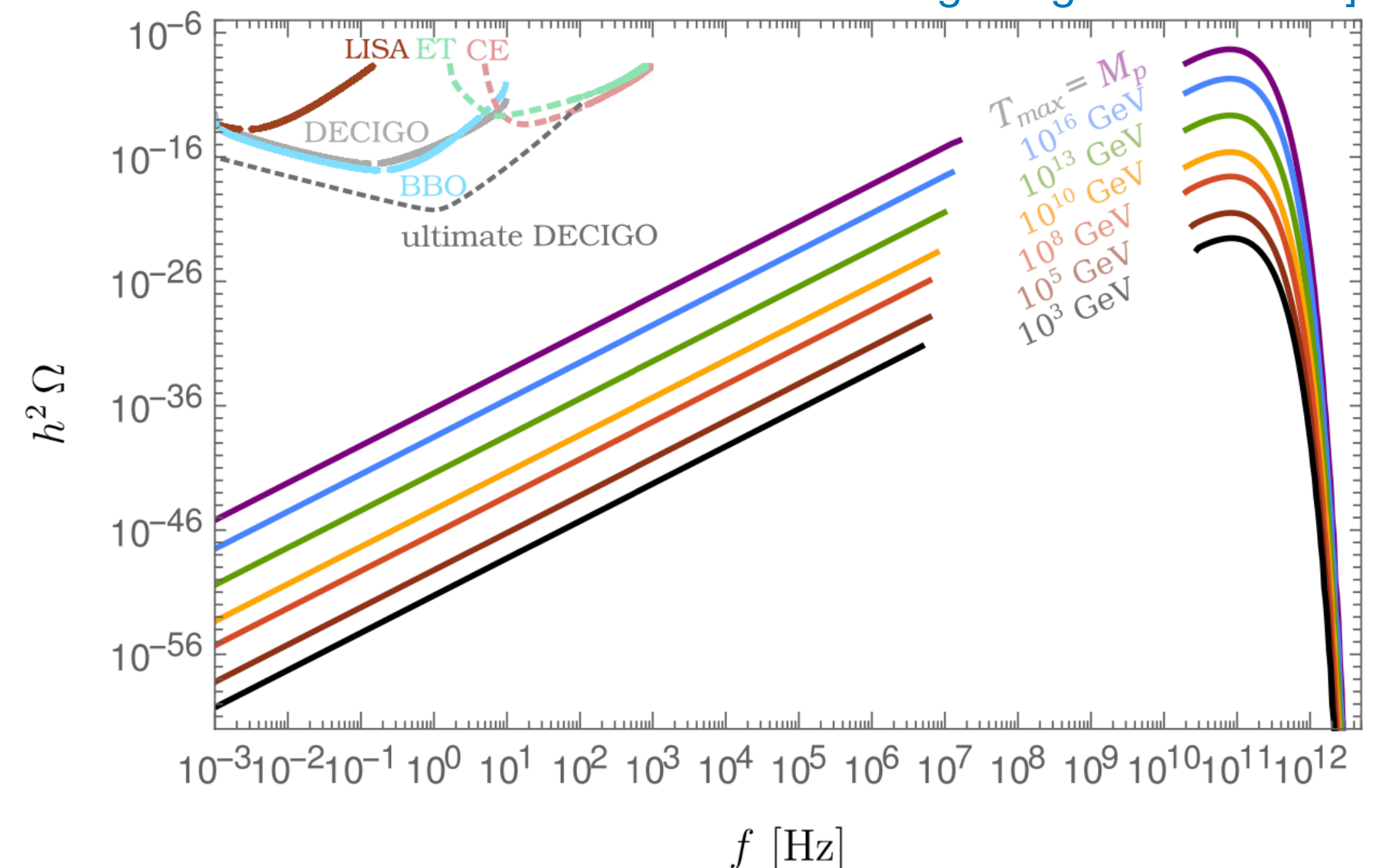
[Klose, Laine, Procacci: 2201.02317, 2210.11710]

## Dark sector confinement phase transition

[Caprini et al.: 1910.13125] [Schwaller:1504.07263]

- But has the dark sector heated up above the confinement temperature?
- Amplitude of gravitational waves from thermal fluctuations grows significantly with maximum reached temperature!

[Ringwald, Schütte-Engel, Tamarit: 2011.04731  
Gravitational waves as a big bang thermometer]



# NB: abelian vs non-abelian dark sector

- Pseudoscalar inflaton coupled to gauge fields:

See talks by Juan, Matteo and Martino!

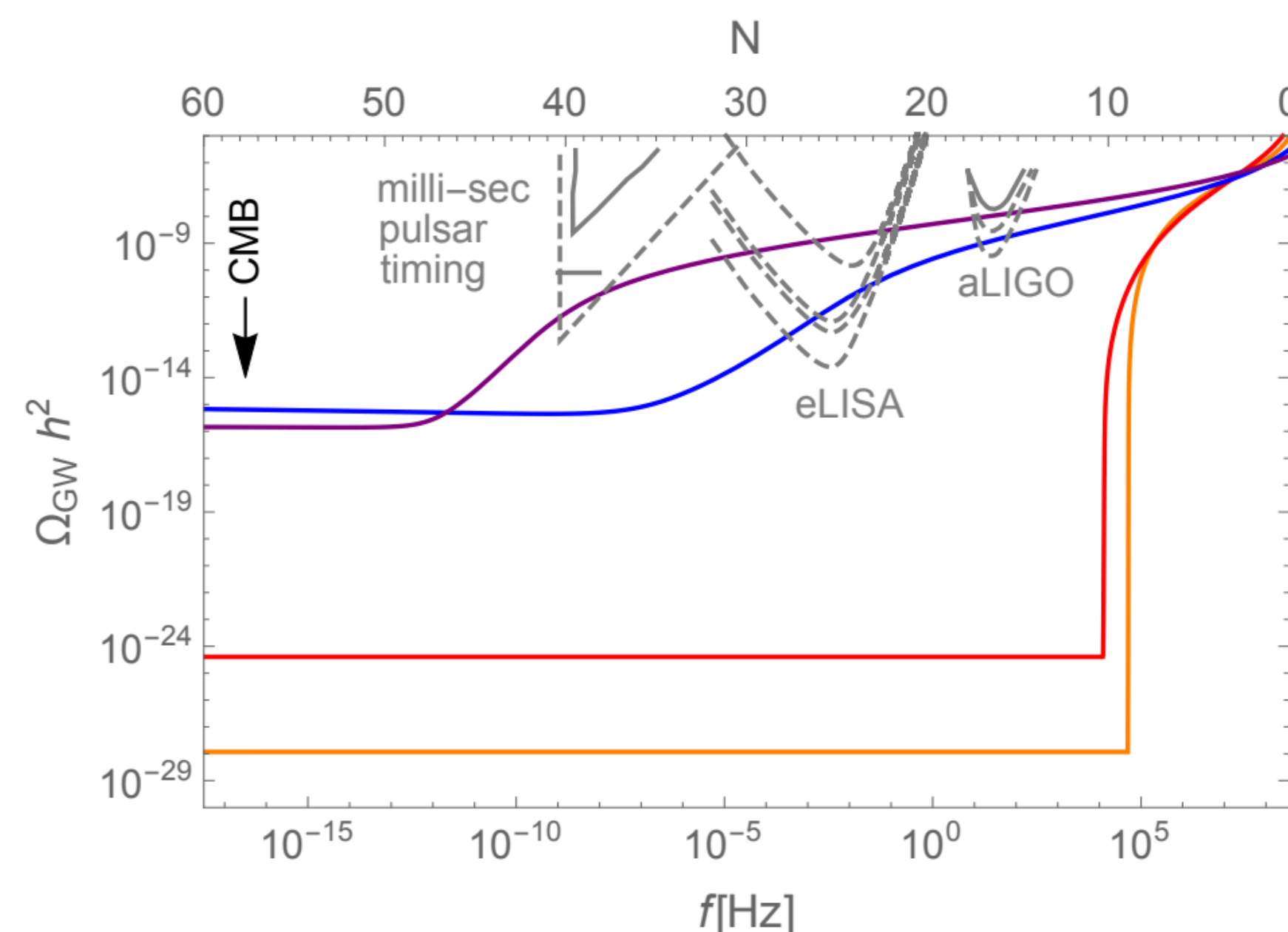
$$\mathcal{L} \supset - \frac{\alpha \varphi \epsilon^{\mu\nu\rho\sigma} F_{\mu\nu} F_{\rho\sigma}}{16\pi f_a}$$

- Abelian case: exponential growth of one helicity mode of the vector field  $\Rightarrow$  GW, PBH, CMB non-gaussianities...

[Sorbo: 1101.1525;  
Cook, Sorbo: 1101.1525;  
Barnaby, Pajer, Peloso: 1110.3327;  
Domcke, Pieroni, Binétruy: 1603.01287...]

- Discussion about back-reaction [... Figuera et al.: 2303.17436]

- Non-Abelian case: thermalisation assumption simplifies the back-reaction modeling!



[Domcke: 1605.06364;  
Domcke, Pieroni, Binétruy: 1603.01287]

# NB: abelian vs non-abelian dark sector

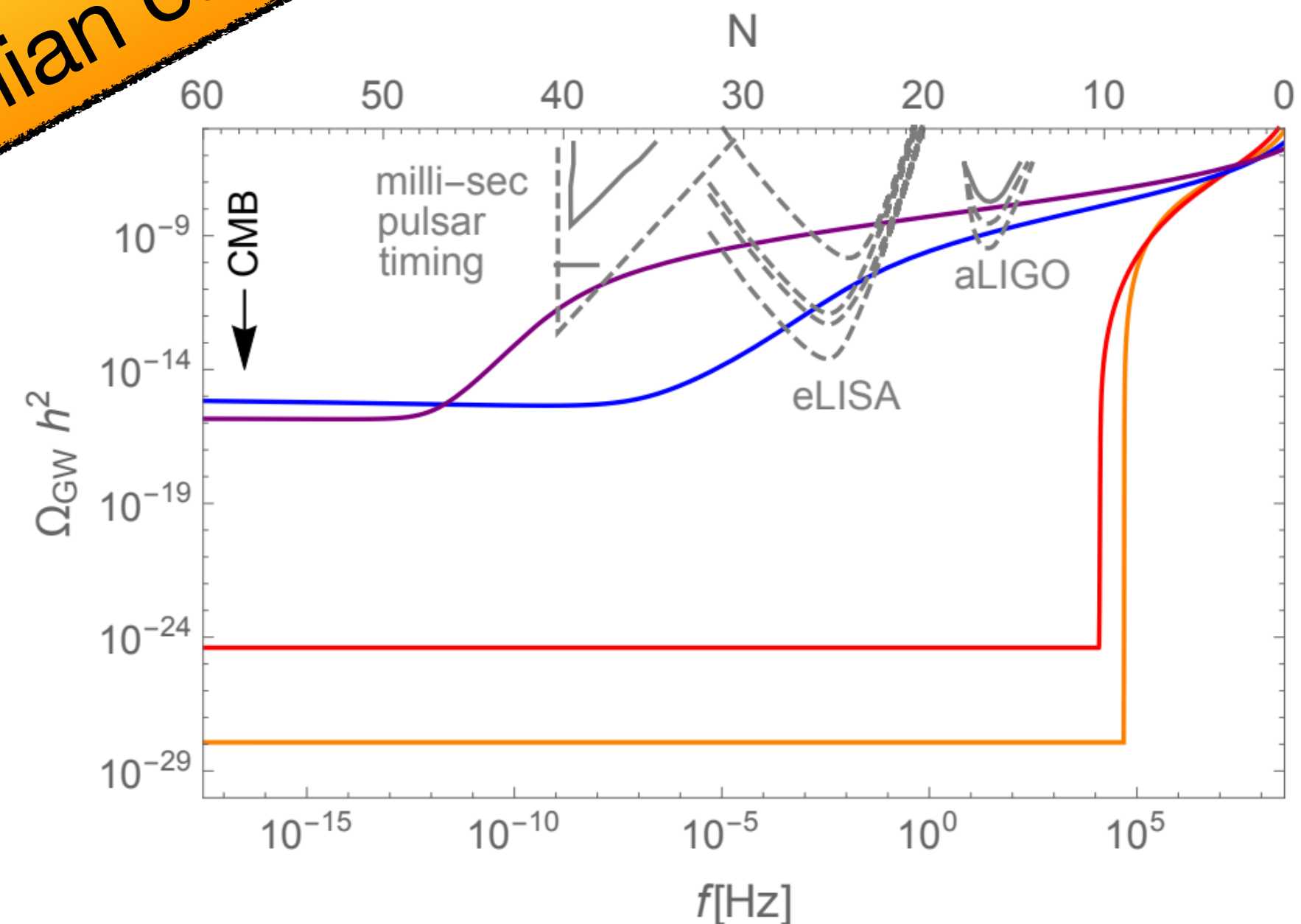
- Pseudoscalar inflaton coupled to gauge fields: [Sorbo, Pajer, and Martino](#) and [Domcke and Martino!](#)

$$\mathcal{L} \supset - \frac{\alpha \varphi \epsilon^{\mu\nu\rho\sigma} F_{\mu\nu} F_{\rho\sigma}}{16\pi f_a}$$

Absent in thermal non-abelian case!

- Abelian case: exponential growth of one helicity mode of the vector field  $\Rightarrow$  GW, PBH, CMB non-gaussianities...

[Sorbo: 1101.1525;  
Cook, Sorbo: 1101.1525;  
Barnaby, Pajer, Peloso: 1110.3327;  
Domcke, Pieroni, Binétruy: 1603.01287...]

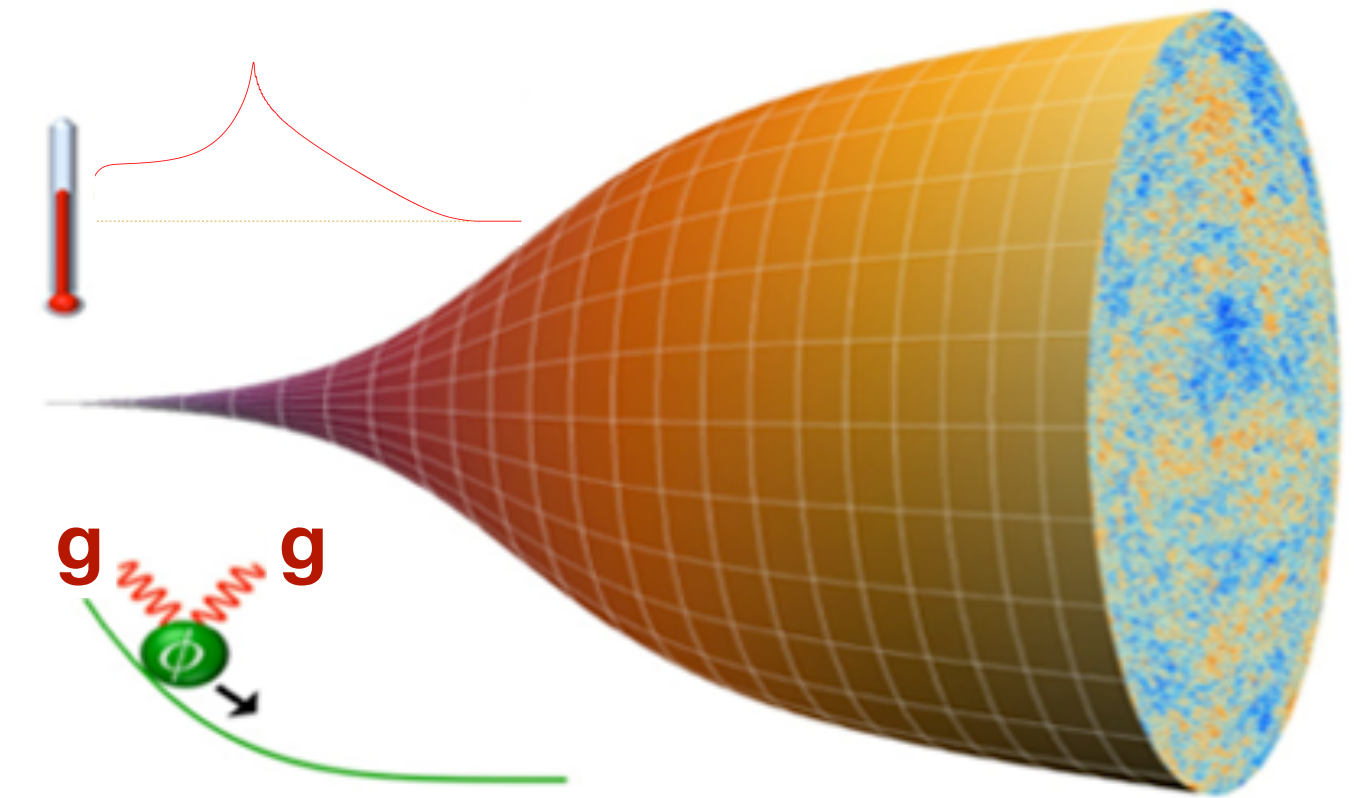


[Domcke: 1605.06364;  
Domcke, Pieroni, Binétruy: 1603.01287]

- Discussion about back-reaction [... Figuera et al.: 2303.17436]
- Non-Abelian case: thermalisation assumption simplifies the back-reaction modeling!

# Model setup

- Inflaton coupled to non-abelian dark sector (parametrised by confinement scale  $\Lambda_{\text{IR}}$ )



Credit: João G. Rosa/University of Aveiro; ESA and the Planck collaboration

$$\mathcal{L} \supset \frac{1}{2} \partial^\mu \varphi \partial_\mu \varphi - V_0(\varphi) - \frac{\alpha \varphi \epsilon^{\mu\nu\rho\sigma} F_{\mu\nu}^c F_{\rho\sigma}^c}{16\pi f_a}$$

Inflaton field

Inflaton potential:

$$V_0 \simeq m^2 f_a^2 \left[ 1 - \cos\left(\frac{\bar{\varphi}}{f_a}\right) \right]$$

“Natural/axion inflation”  
[Freese, Frieman, Olinto:  
Phys.Rev.Lett. 65 (1990)]

Gauge coupling

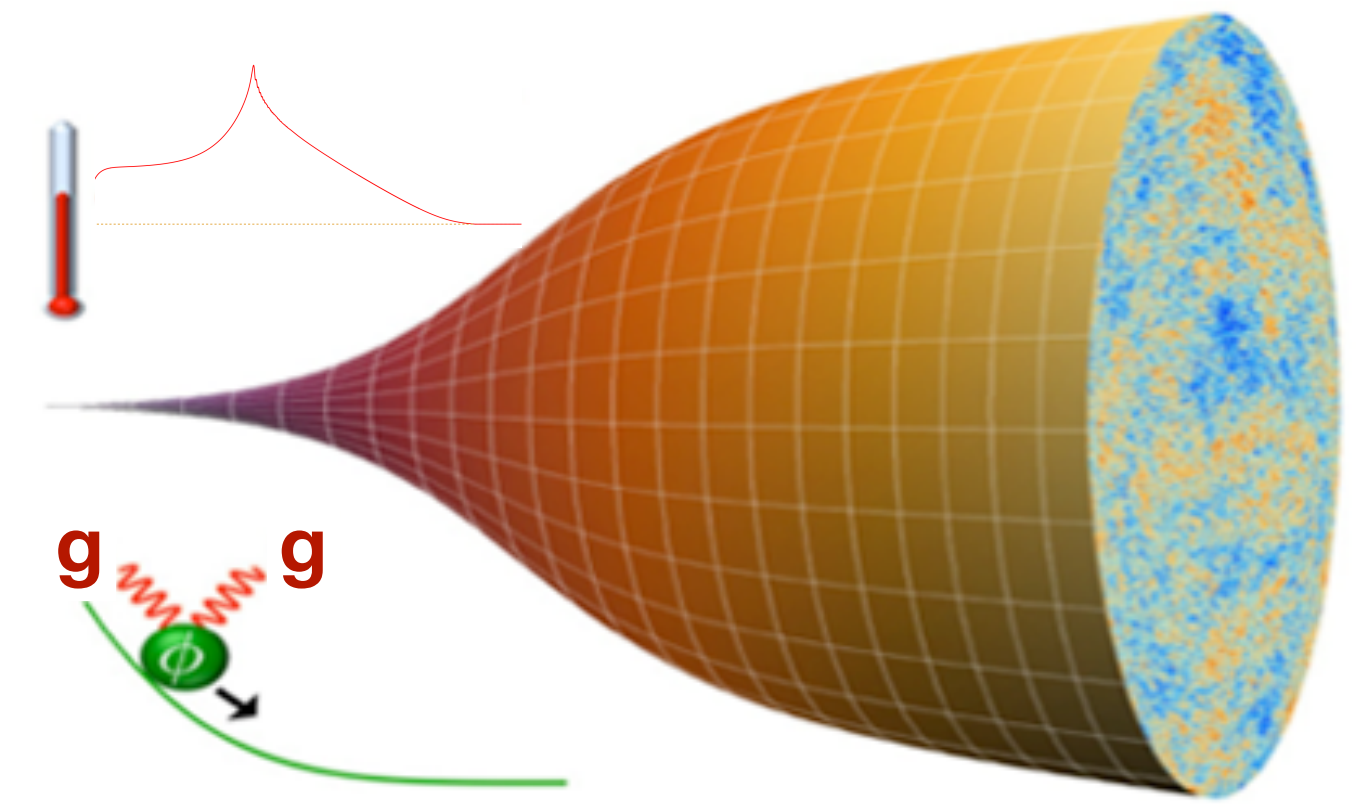
Yang-Mills  
field strength

Axion decay  
constant



# Model setup

- Inflaton coupled to non-abelian dark sector (parametrised by confinement scale  $\Lambda_{\text{IR}}$ )



Credit: João G. Rosa/University of Aveiro; ESA and the Planck collaboration

$$\mathcal{L} \supset \frac{1}{2} \partial^\mu \varphi \partial_\mu \varphi - V_0(\varphi) - \frac{\alpha \varphi \epsilon^{\mu\nu\rho\sigma} F_{\mu\nu}^c F_{\rho\sigma}^c}{16\pi f_a}$$

Inflaton field

Inflaton potential:

$$V_0 \simeq m^2 f_a^2 \left[ 1 - \cos\left(\frac{\bar{\varphi}}{f_a}\right) \right]$$

“Natural/axion inflation”  
[Freese, Frieman, Olinto:  
Phys.Rev.Lett. 65 (1990)]

Gauge coupling

Yang-Mills  
field strength

Axion decay  
constant

Hubble rate

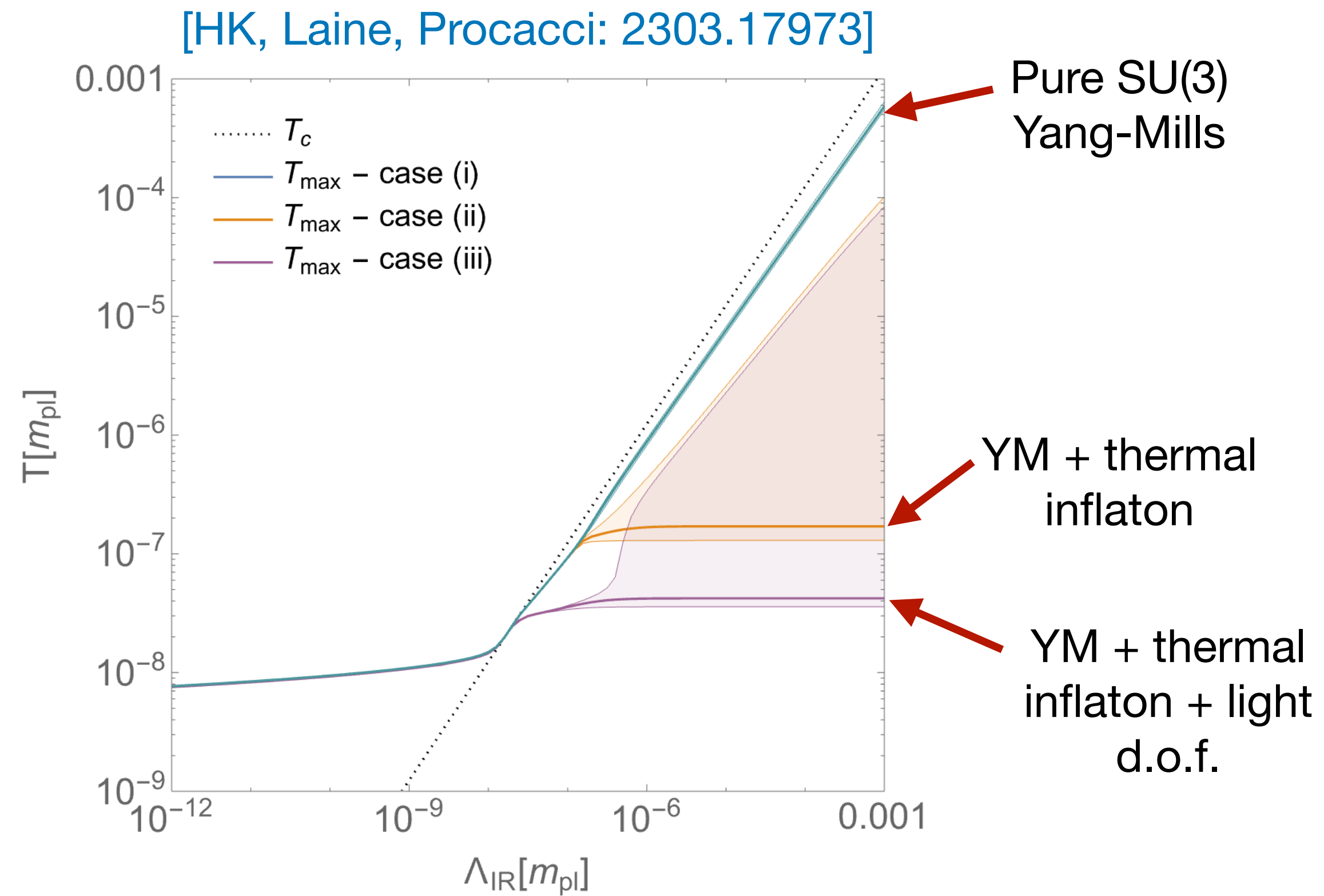
Friction due to inflaton  
coupling to dark sector!  
[Laine, Niemi, Procacci,  
Rummukainen: 2209.13804]

$$\ddot{\bar{\varphi}} + (3H + \Upsilon)\dot{\bar{\varphi}} + V_\varphi \simeq 0$$

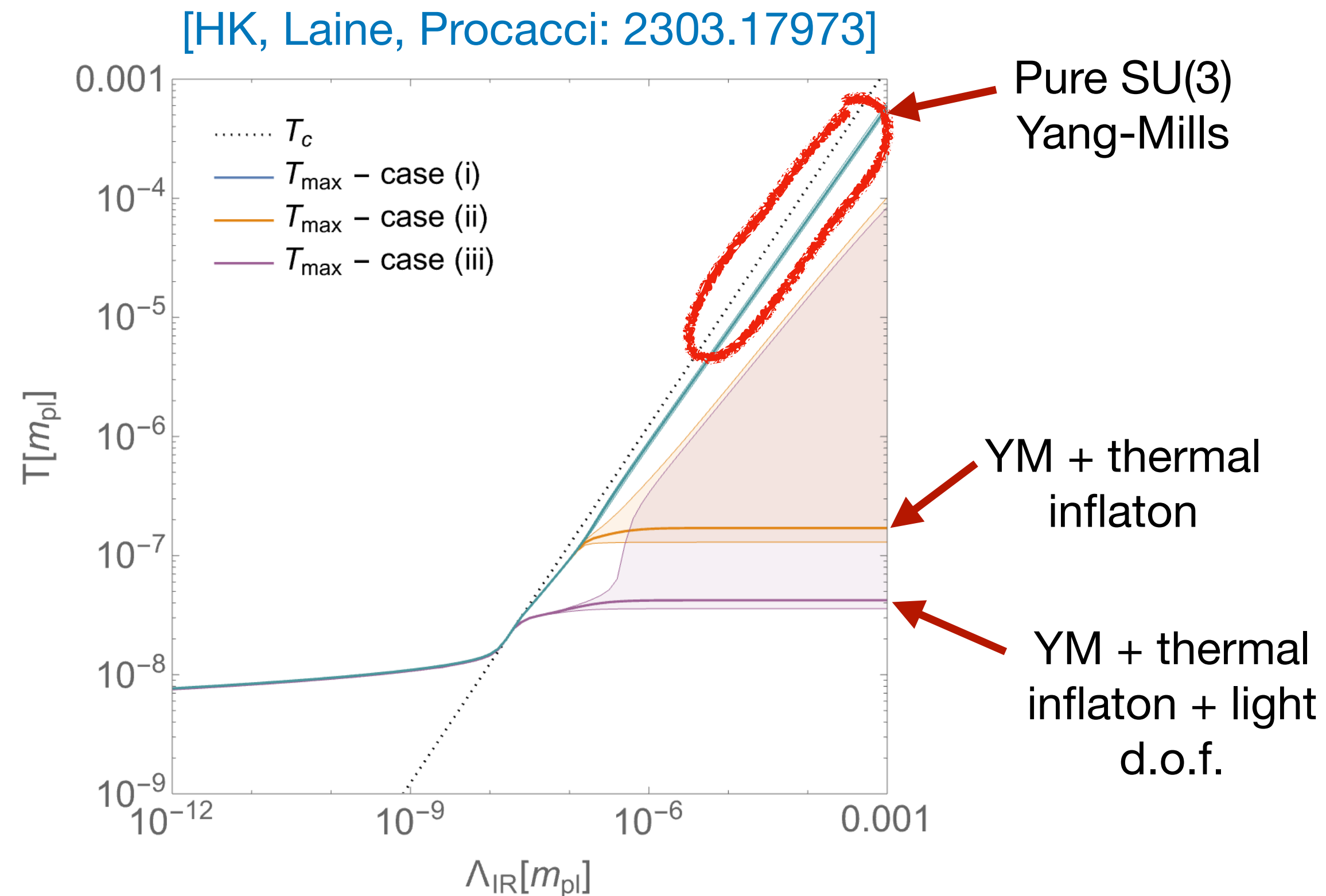
$$\dot{e}_r + 3H(e_r + p_r) \simeq \Upsilon \dot{\bar{\varphi}}^2$$

Dark radiation energy  
and pressure densities

# Results: Maximal temperature of the dark sector



# Results: Maximal temperature of the dark sector

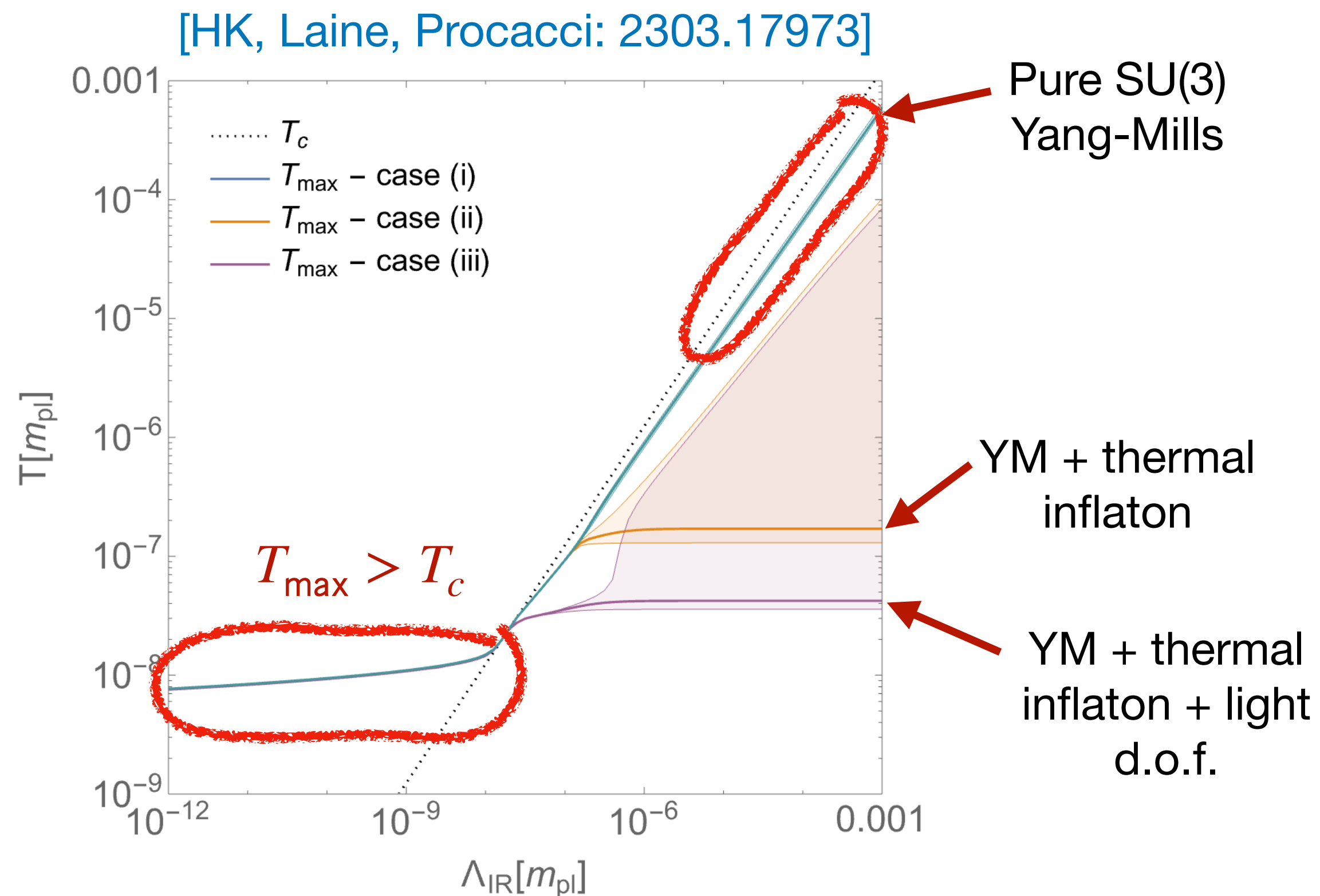


**Message 1:** Temperatures up to  $10^{-3} m_{\text{pl}}$  can be reached

- For  $\Lambda_{\text{IR}}$  up to  $10^{-3} m_{\text{pl}}$
- If Yang-Mills plasma not coupled to extra light d.o.f.

⇒ Hope for potentially detectable GW from thermal plasma in ET, LISA frequency range!

# Results: Maximal temperature of the dark sector



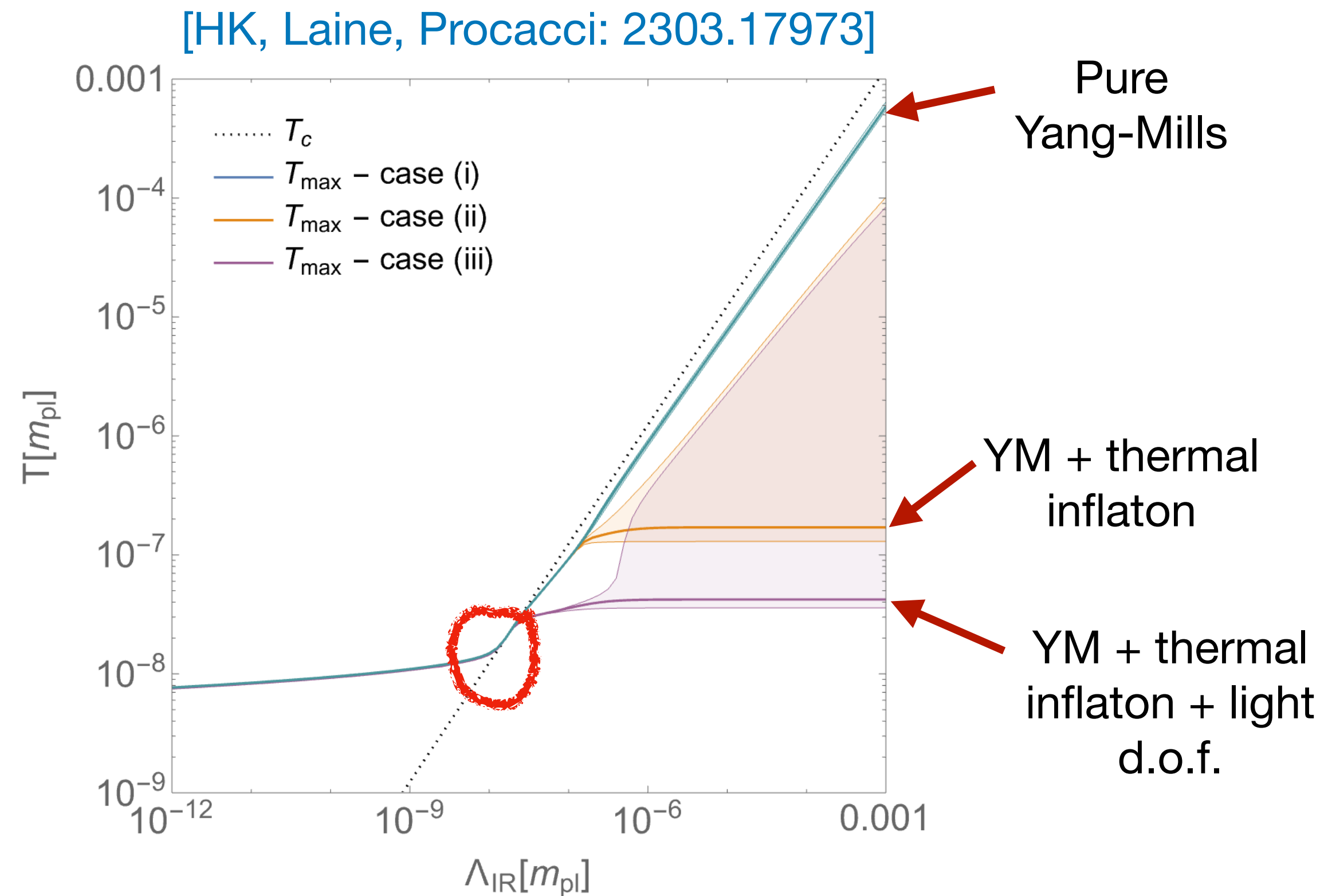
**Message 1:** Temperatures up to  $10^{-3} m_{\text{pl}}$  can be reached

- For  $\Lambda_{\text{IR}}$  up to  $10^{-3} m_{\text{pl}}$
- If Yang-Mills plasma not coupled to extra light d.o.f.

⇒ Hope for potentially detectable GW from thermal plasma in ET, LISA frequency range!

**Message 2:** For lower  $\Lambda_{\text{IR}}$  the dark sector heats up above  $T_c$   
 ⇒ undergoes a phase transition! Possible further GW signal!

# Results: Maximal temperature of the dark sector



**Message 3: There can even be two phase transitions!**

# Example: $\Lambda_{\text{IR}} = 10^{-8} m_{\text{pl}}$

Benchmark parameter choice  
(axion inflation consistent with CMB data)

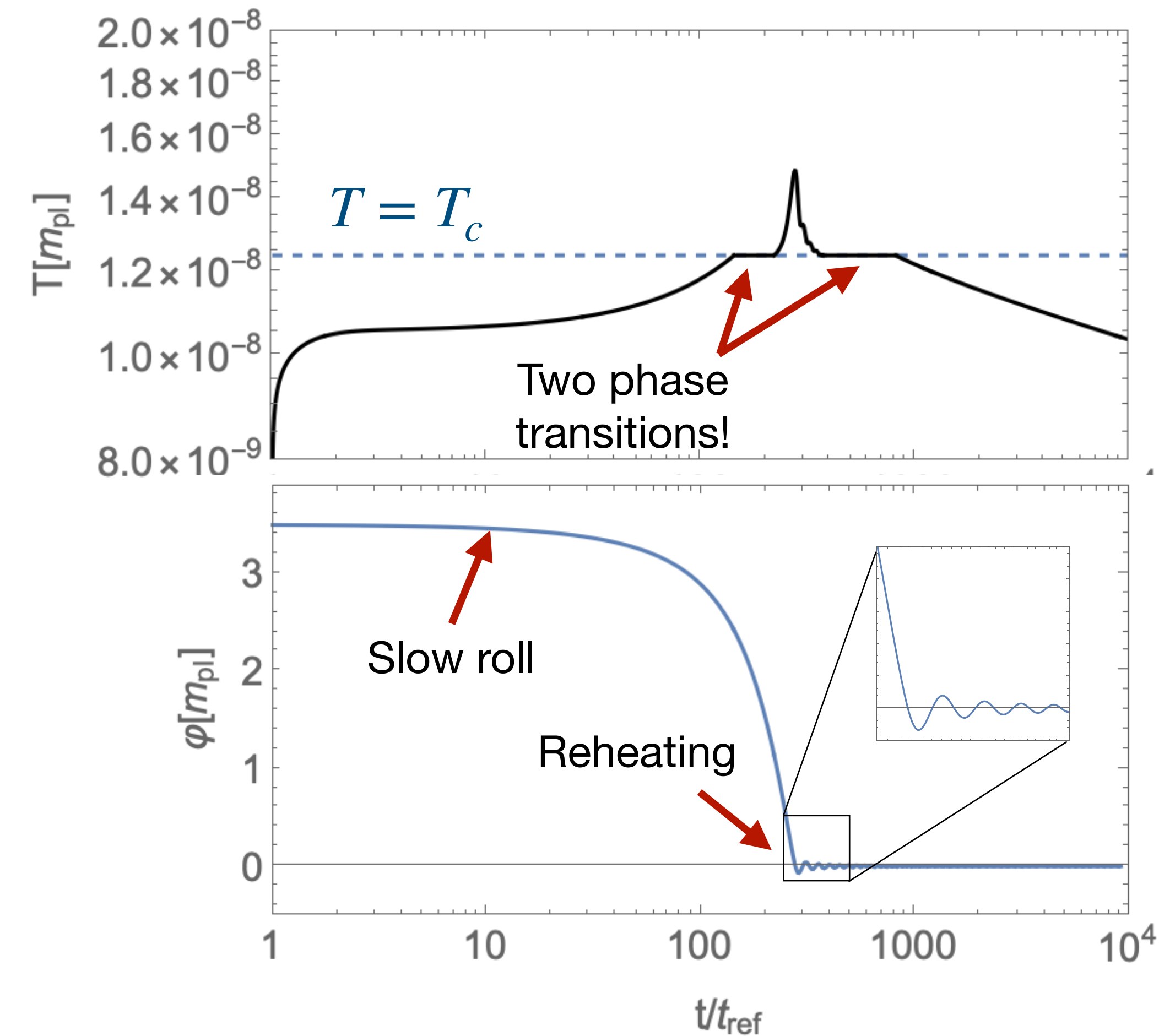
[Klose, Laine, Procacci: 2201.02317]:

axion mass:  $m = 1.09 \times 10^{-6} m_{\text{pl}}$ ,

axion decay constant:  $f_a = 1.25 m_{\text{pl}}$ ,

initial time:  $t_{\text{ref}} \sim H_{\text{initial}}^{-1}$

$$\Rightarrow \Upsilon \sim 10^{-23} m_{\text{pl}} \ll H_{\text{initial}}$$



# Example: $\Lambda_{\text{IR}} = 10^{-8} m_{\text{pl}}$

Benchmark parameter choice  
(axion inflation consistent with CMB data)

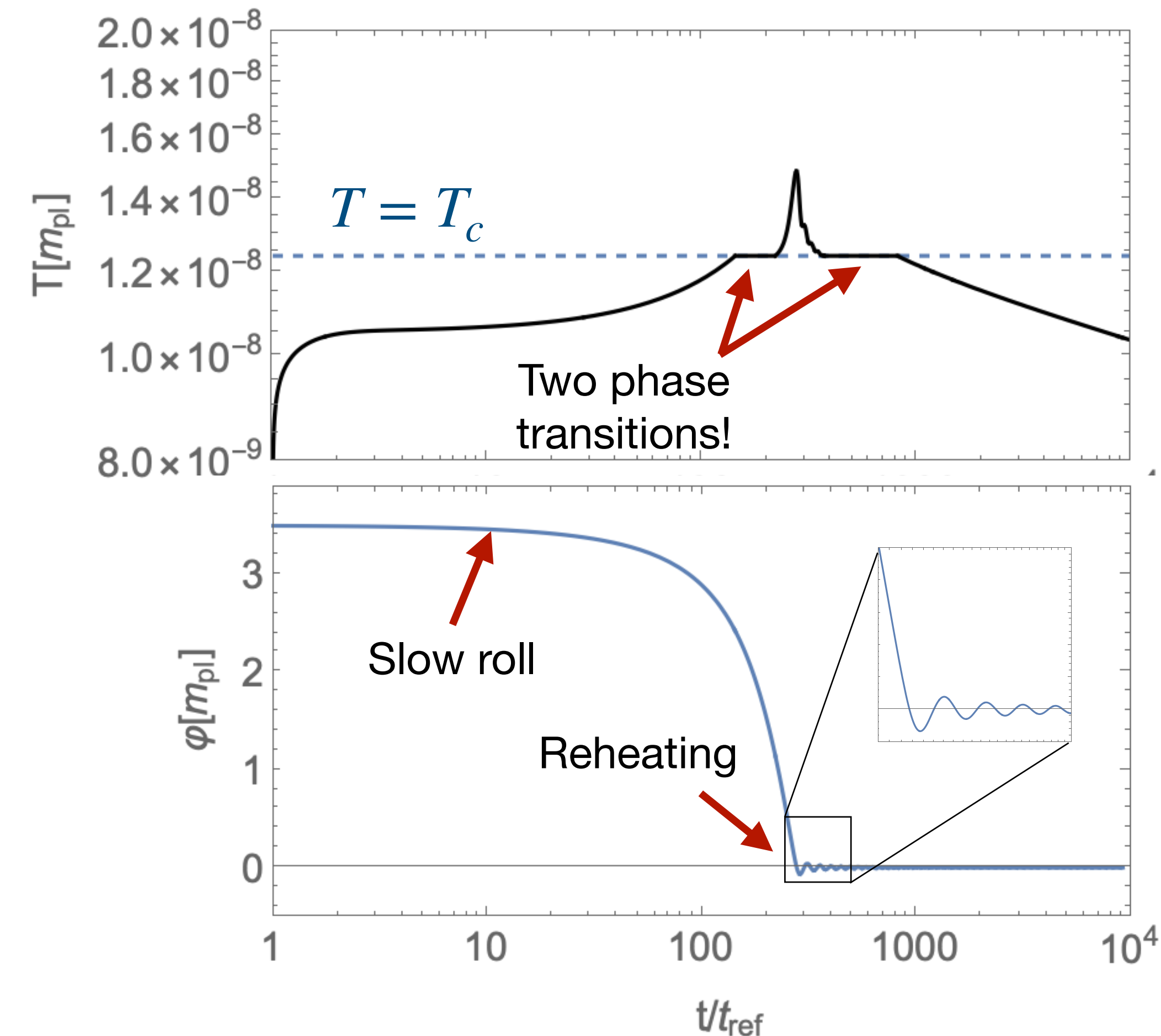
[Klose, Laine, Procacci: 2201.02317]:

axion mass:  $m = 1.09 \times 10^{-6} m_{\text{pl}}$ ,

axion decay constant:  $f_a = 1.25 m_{\text{pl}}$ ,

initial time:  $t_{\text{ref}} \sim H_{\text{initial}}^{-1}$

$$\Rightarrow \Upsilon \sim 10^{-23} m_{\text{pl}} \ll H_{\text{initial}}$$



- Heating and cooling phase transitions may bring interesting GW signatures!

[Buen-Abad, Chang, Hook: 2305.09712]

# Example: $\Lambda_{\text{IR}} = 10^{-8} m_{\text{pl}}$

Benchmark parameter choice  
(axion inflation consistent with CMB data)

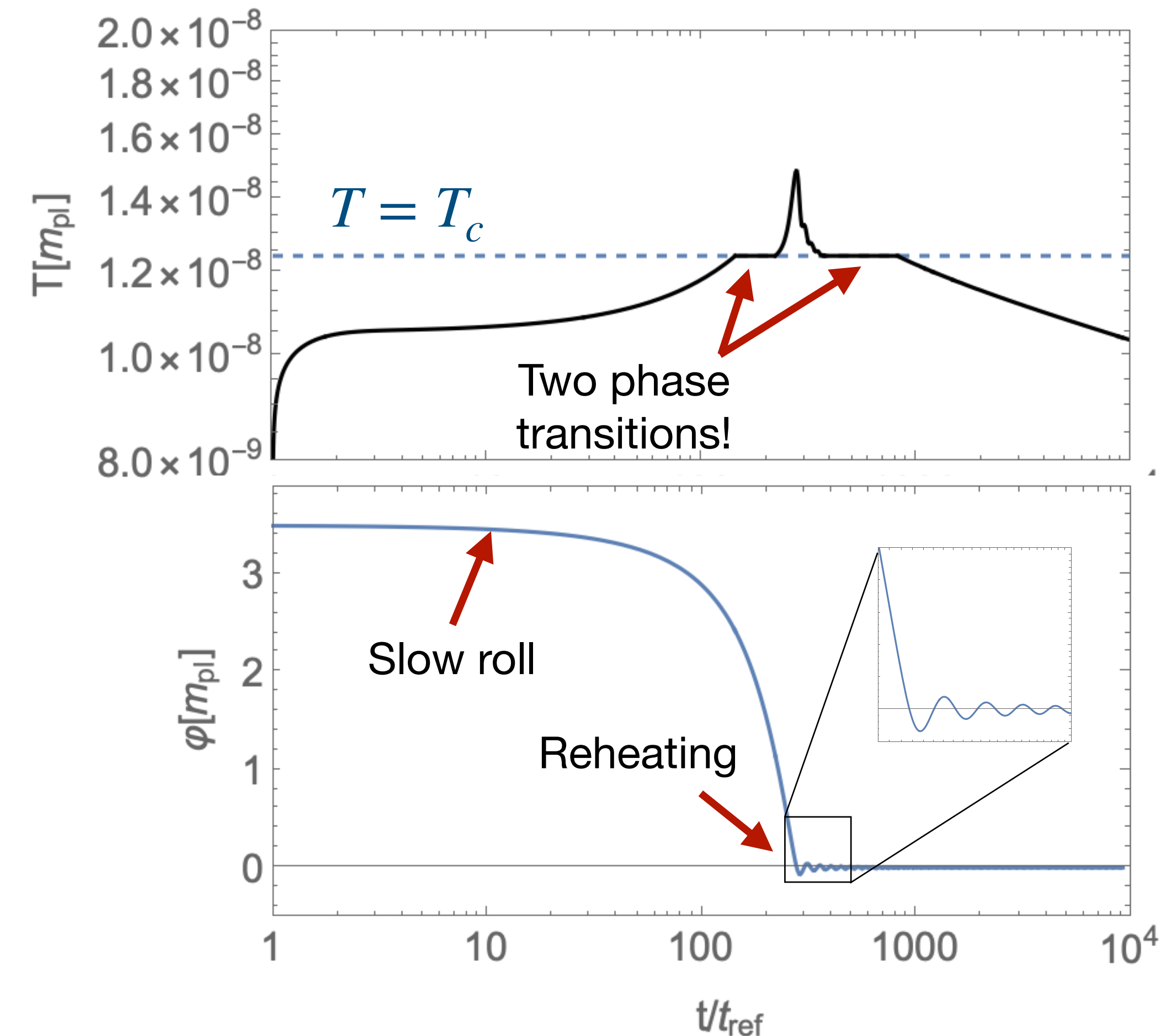
[Klose, Laine, Procacci: 2201.02317]:

axion mass:  $m = 1.09 \times 10^{-6} m_{\text{pl}}$ ,

axion decay constant:  $f_a = 1.25 m_{\text{pl}}$ ,

initial time:  $t_{\text{ref}} \sim H_{\text{initial}}^{-1}$

$$\Rightarrow \Upsilon \sim 10^{-23} m_{\text{pl}} \ll H_{\text{initial}}$$



- Heating and cooling phase transitions may bring interesting GW signatures!

[Buen-Abad, Chang, Hook: 2305.09712]

**BUT!**

- $\Upsilon \ll H \Rightarrow$  Long inflaton domination  $\Rightarrow$  more significant GW dilution
- How and when SM is reheated?



# Example: $\Lambda_{\text{IR}} = 10^{-8} m_{\text{pl}}$

Benchmark parameter choice  
(axion inflation consistent with CMB data)

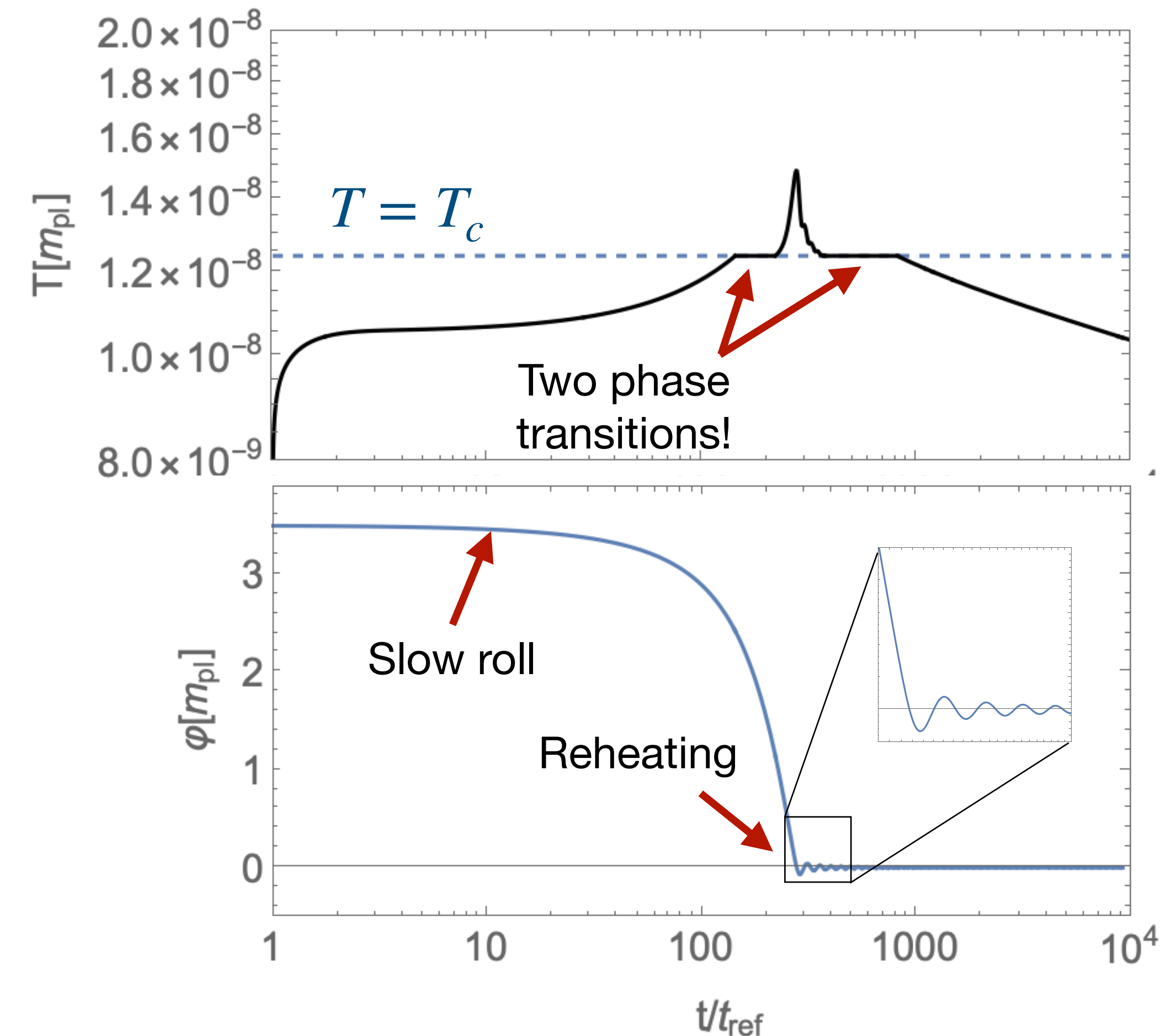
[Klose, Laine, Procacci: 2201.02317]:

axion mass:  $m = 1.09 \times 10^{-6} m_{\text{pl}}$ ,

axion decay constant:  $f_a = 1.25 m_{\text{pl}}$ ,

initial time:  $t_{\text{ref}} \sim H_{\text{initial}}^{-1}$

$$\Rightarrow \Upsilon \sim 10^{-23} m_{\text{pl}} \ll H_{\text{initial}}$$



- Heating and cooling phase transitions may bring interesting GW signatures!

[Buen-Abad, Chang, Hook: 2305.09712]


**BUT!**

- $\Upsilon \ll H \Rightarrow$  Long inflaton domination more significant GW dilution

- How and when SM is re...

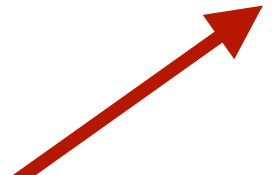
**To be worked out!**

# Conclusions

- Evolution of a **dark Yang-Mills sector** coupled to axion inflation studied
- Different qualitative behaviour depending on the **dark confinement scale**  $\Lambda_{\text{IR}}$ :
  -   $\Lambda_{\text{IR}} \lesssim 10^{-8} m_{\text{pl}} : T_{\text{max}} \sim 10^{-8} m_{\text{pl}} > T_c \Rightarrow$  Phase transition
  - $\Lambda_{\text{IR}} \gtrsim 10^{-8} m_{\text{pl}} : T_{\text{max}}$  slightly below  $T_c \Rightarrow$  Large temperatures achieved
- Possible **gravitational wave signal** in both cases

# Conclusions

- Evolution of a **dark Yang-Mills sector** coupled to axion inflation studied
- Different qualitative behaviour depending on the **dark confinement scale**  $\Lambda_{\text{IR}}$ :

  $\Lambda_{\text{IR}} \lesssim 10^{-8} m_{\text{pl}} : T_{\text{max}} \sim 10^{-8} m_{\text{pl}} > T_c \Rightarrow$  Phase transition

  $\Lambda_{\text{IR}} \gtrsim 10^{-8} m_{\text{pl}} : T_{\text{max}}$  slightly below  $T_c \Rightarrow$  Large temperatures achieved

- Possible **gravitational wave signal** in both cases

**Stay tuned :) Thanks for your attention!**

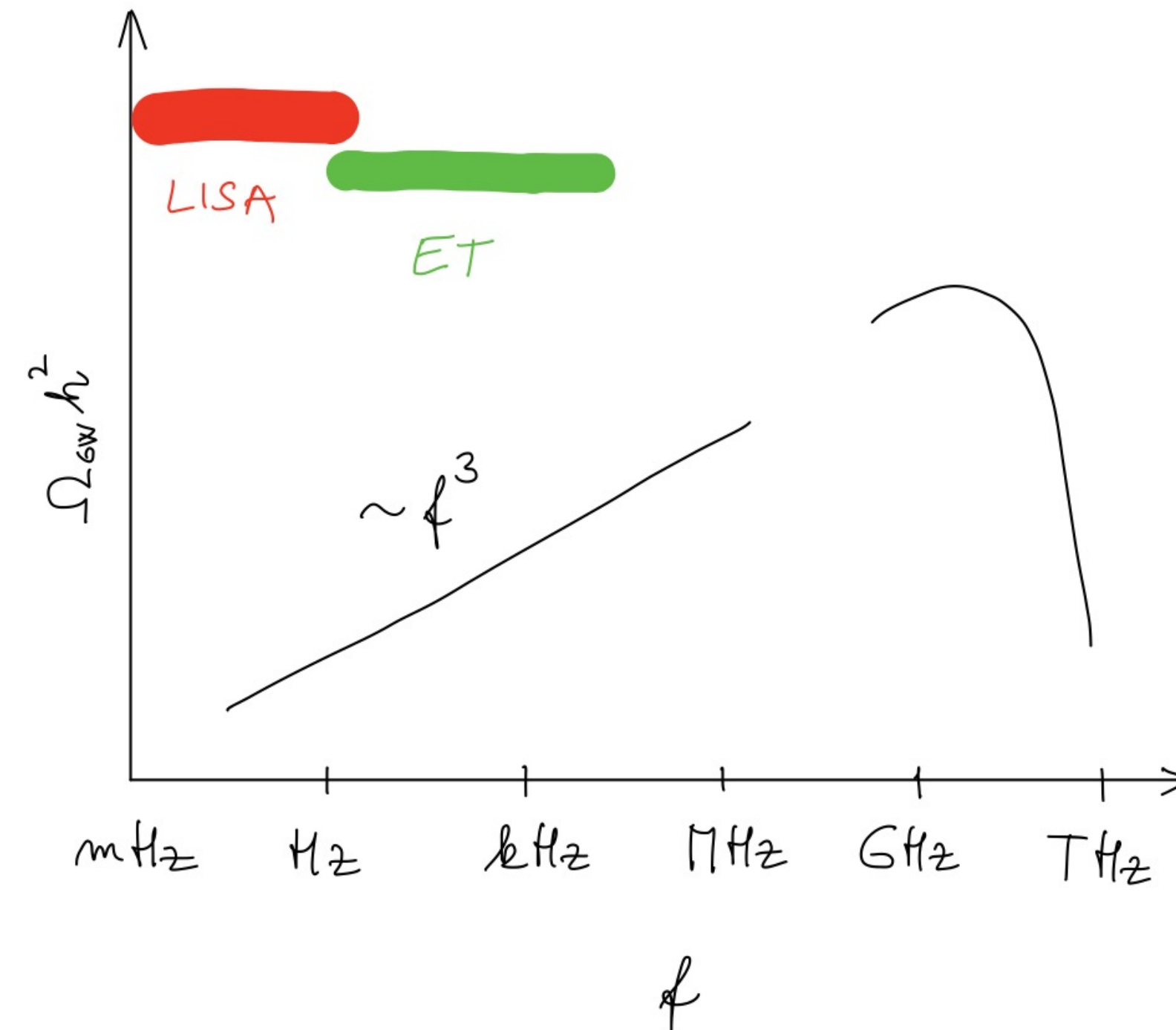
**Back up**

# GW from thermal plasma

[Ghiglieri, Laine: 1504.02569]

[Klose, Laine, Procacci: 2201.02317]

[Klose, Laine, Procacci: 2210.11710]

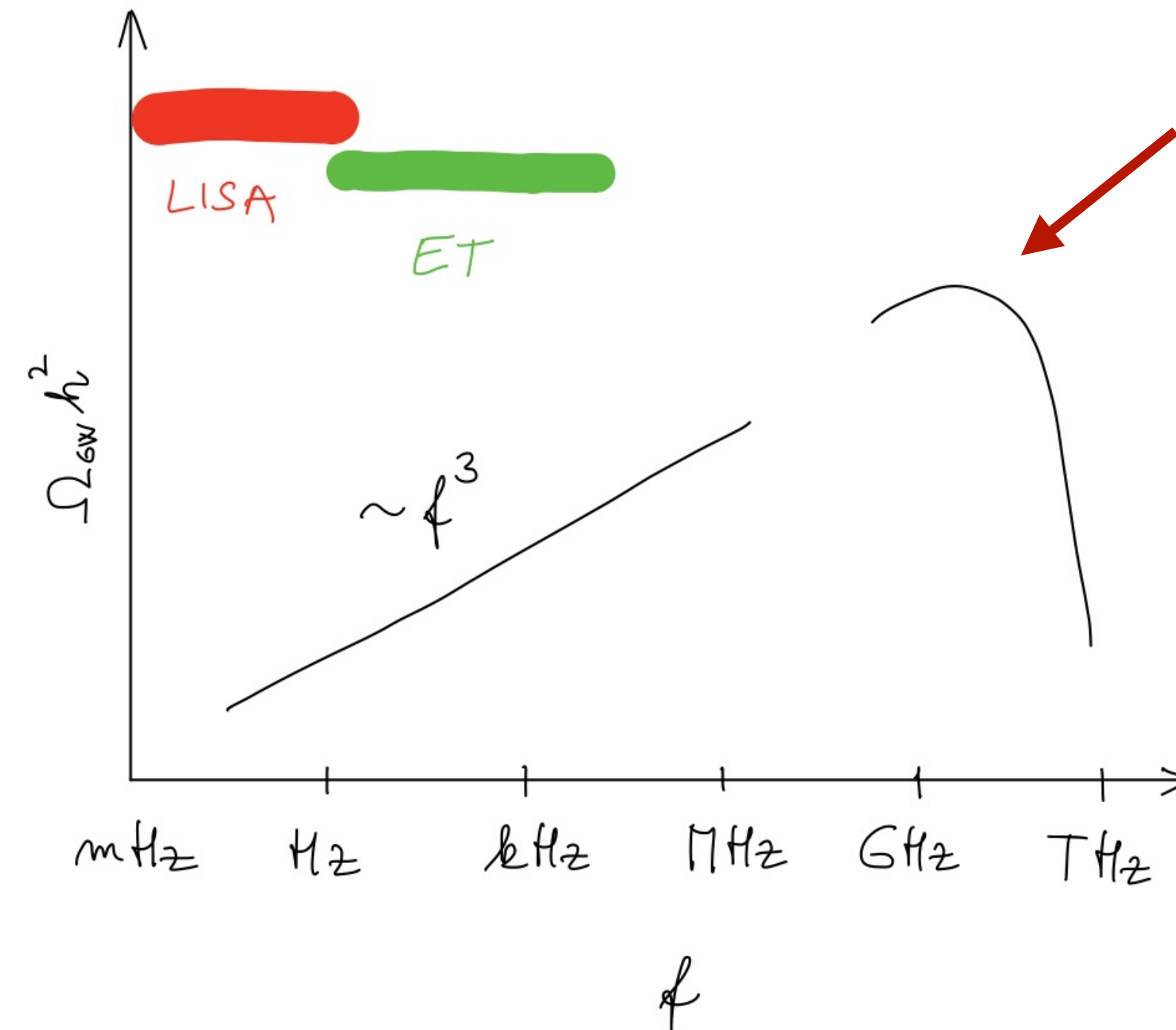


# GW from thermal plasma

[Ghiglieri, Laine: 1504.02569]

[Klose, Laine, Procacci: 2201.02317]

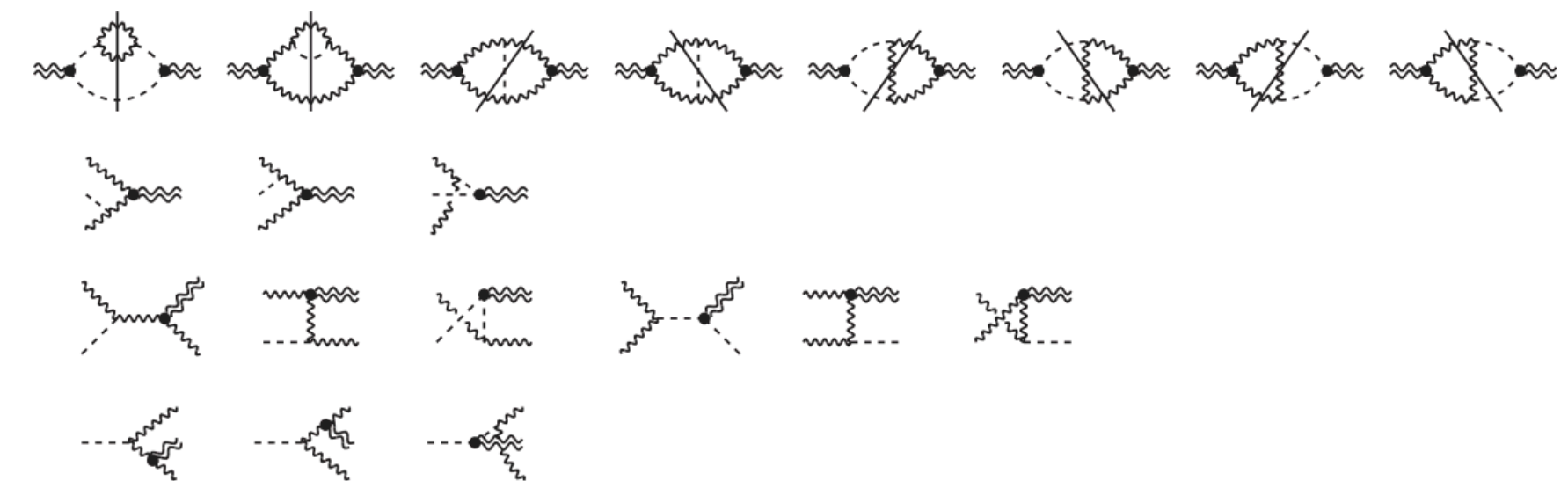
[Klose, Laine, Procacci: 2210.11710]



“Boltzmann domain” - determines total energy density in GW  $\Leftrightarrow$  constraints from  $\Delta N_{\text{eff}}$  at BBN

$$(\Delta N_{\text{eff}} \lesssim 10^{-3} \Rightarrow T_{\text{max}} \lesssim 10^{17} \text{ GeV})$$

[Klose, Laine, Procacci: 2201.02317]

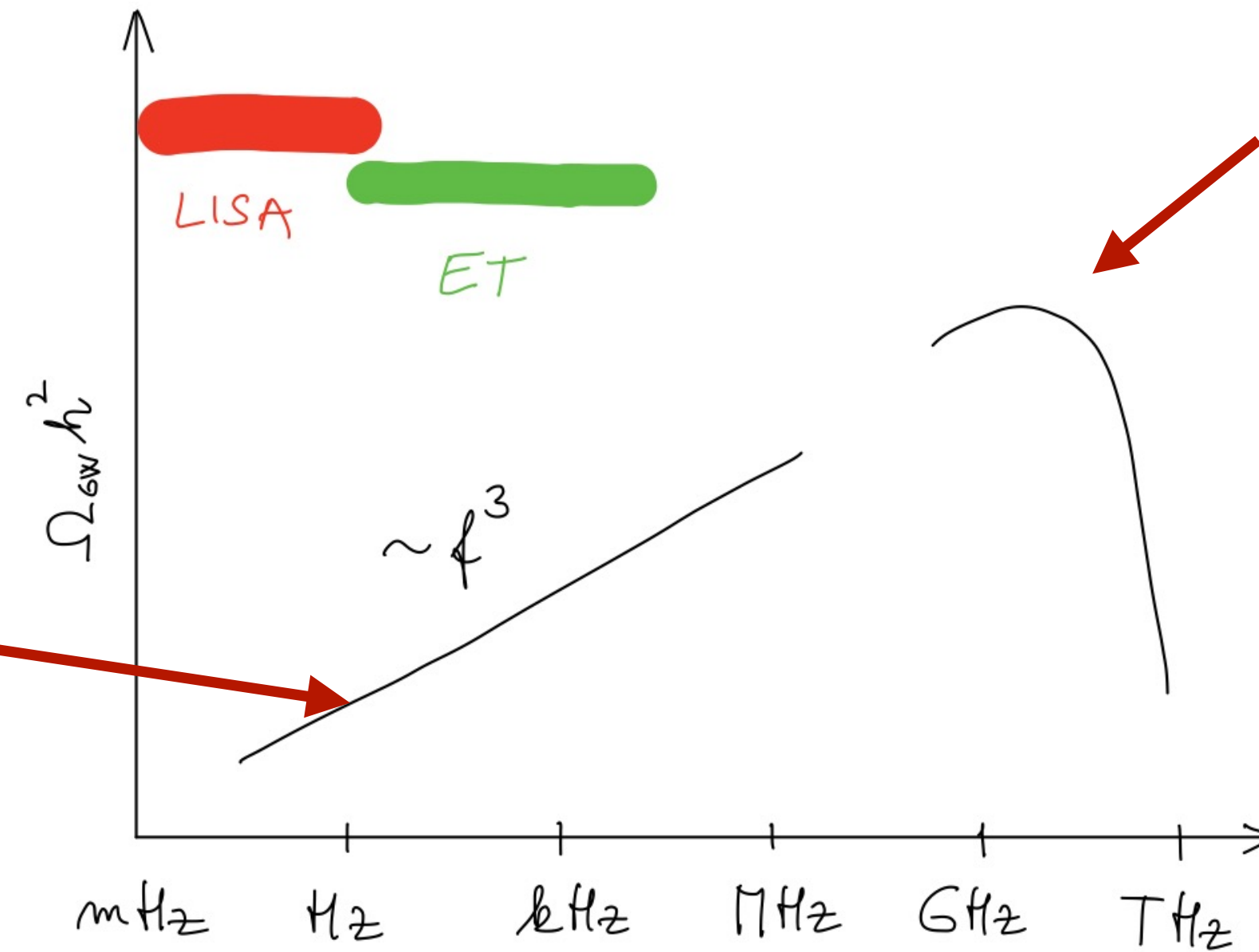


# GW from thermal plasma

[Ghiglieri, Laine: 1504.02569]

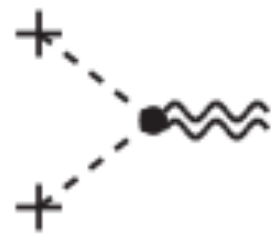
[Klose, Laine, Procacci: 2201.02317]

[Klose, Laine, Procacci: 2210.11710]

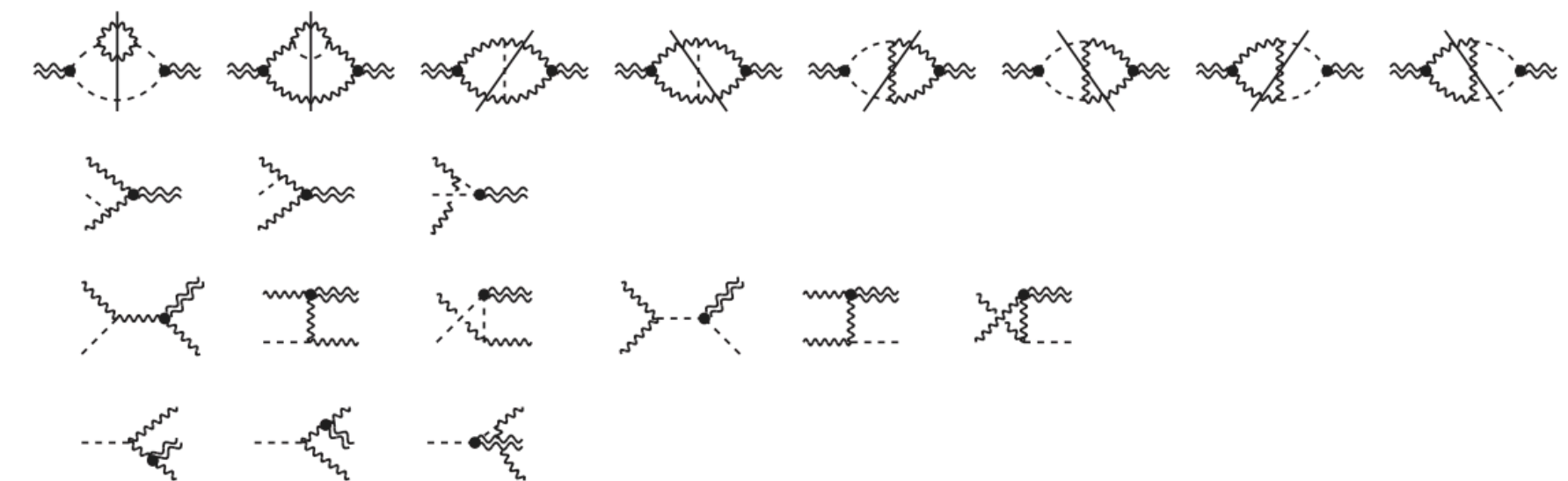


“Boltzmann domain” - determines total energy density in GW  $\Leftrightarrow$  constraints from  $\Delta N_{\text{eff}}$  at BBN  
 $(\Delta N_{\text{eff}} \lesssim 10^{-3} \Rightarrow T_{\text{max}} \lesssim 10^{17} \text{ GeV})$

“Hydrodynamic domain”: GW signal driven by shear viscosity  $\eta$



[Klose, Laine, Procacci: 2201.02317]



$$\Omega_{\text{GW}} h^2 \supset A \left( \frac{f}{\text{Hz}} \right)^3 \frac{(T\eta)_{\text{max}}}{m_{\text{pl}}^4}, \quad (T\eta)_{\text{max}} \sim \frac{T_{\text{max}}^4}{\alpha_{\text{min}}^2}$$

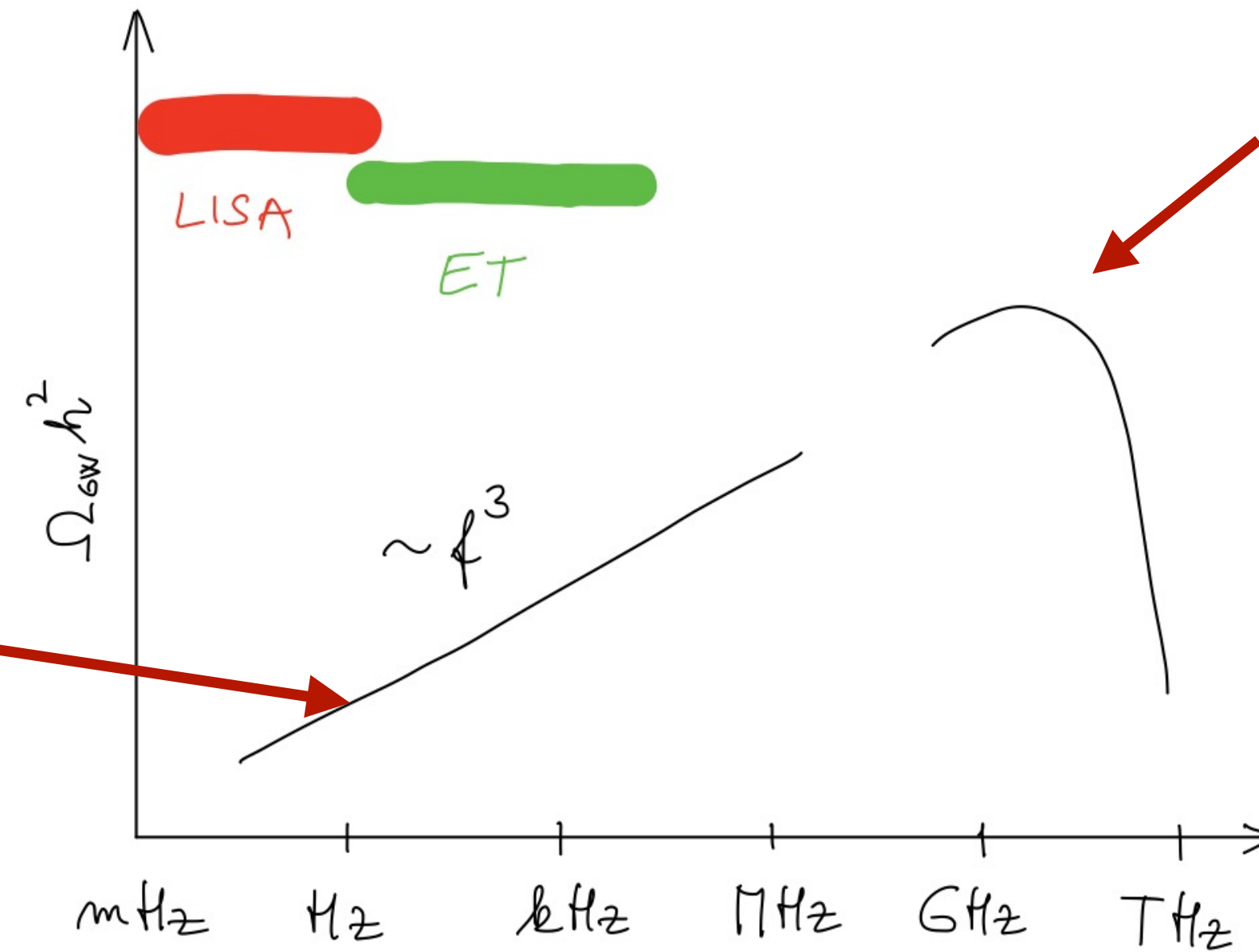
$A \sim 10^{-9}$  - based on calculation performed for  $\Lambda \sim \text{GeV}$

# GW from thermal plasma

[Ghiglieri, Laine: 1504.02569]

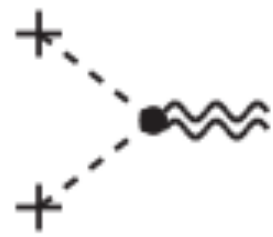
[Klose, Laine, Procacci: 2201.02317]

[Klose, Laine, Procacci: 2210.11710]

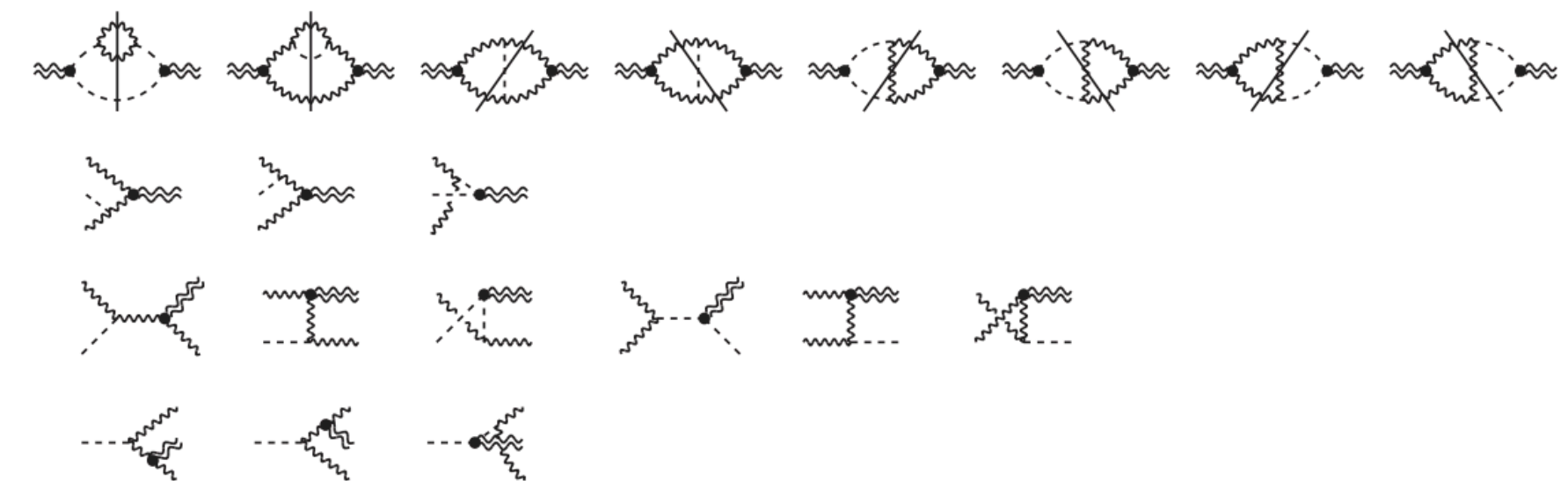


“Boltzmann domain” - determines total energy density in GW  $\Leftrightarrow$  constraints from  $\Delta N_{\text{eff}}$  at BBN  
 $(\Delta N_{\text{eff}} \lesssim 10^{-3} \Rightarrow T_{\text{max}} \lesssim 10^{17} \text{ GeV})$

“Hydrodynamic domain”: GW signal driven by shear viscosity  $\eta$



[Klose, Laine, Procacci: 2201.02317]



$$\Omega_{\text{GW}} h^2 \supset A \left( \frac{f}{\text{Hz}} \right)^3 \frac{(T\eta)_{\text{max}}}{m_{\text{pl}}^4}, \quad (T\eta)_{\text{max}} \sim \frac{T_{\text{max}}^4}{\alpha_{\text{min}}^2}$$

$A \sim 10^{-9}$  - based on calculation performed for  $\Lambda \sim \text{GeV}$

Might be measurable if  $T_{\text{max}} \sim 10^{-3} m_{\text{pl}}$  but more detailed calculation needed!




# Why large temperatures?

$$\ddot{\phi} + (3H + \Upsilon)\dot{\phi} + V_\phi \simeq 0$$

$$\dot{e}_r + 3H(e_r + p_r) \simeq \Upsilon \dot{\phi}^2$$

Parametrize  
 $e_r, p_r$  by  $T$



$$c_r(T) \dot{T} + 3H[e_r(T) + p_r(T)] \simeq \Upsilon \dot{\phi}^2$$

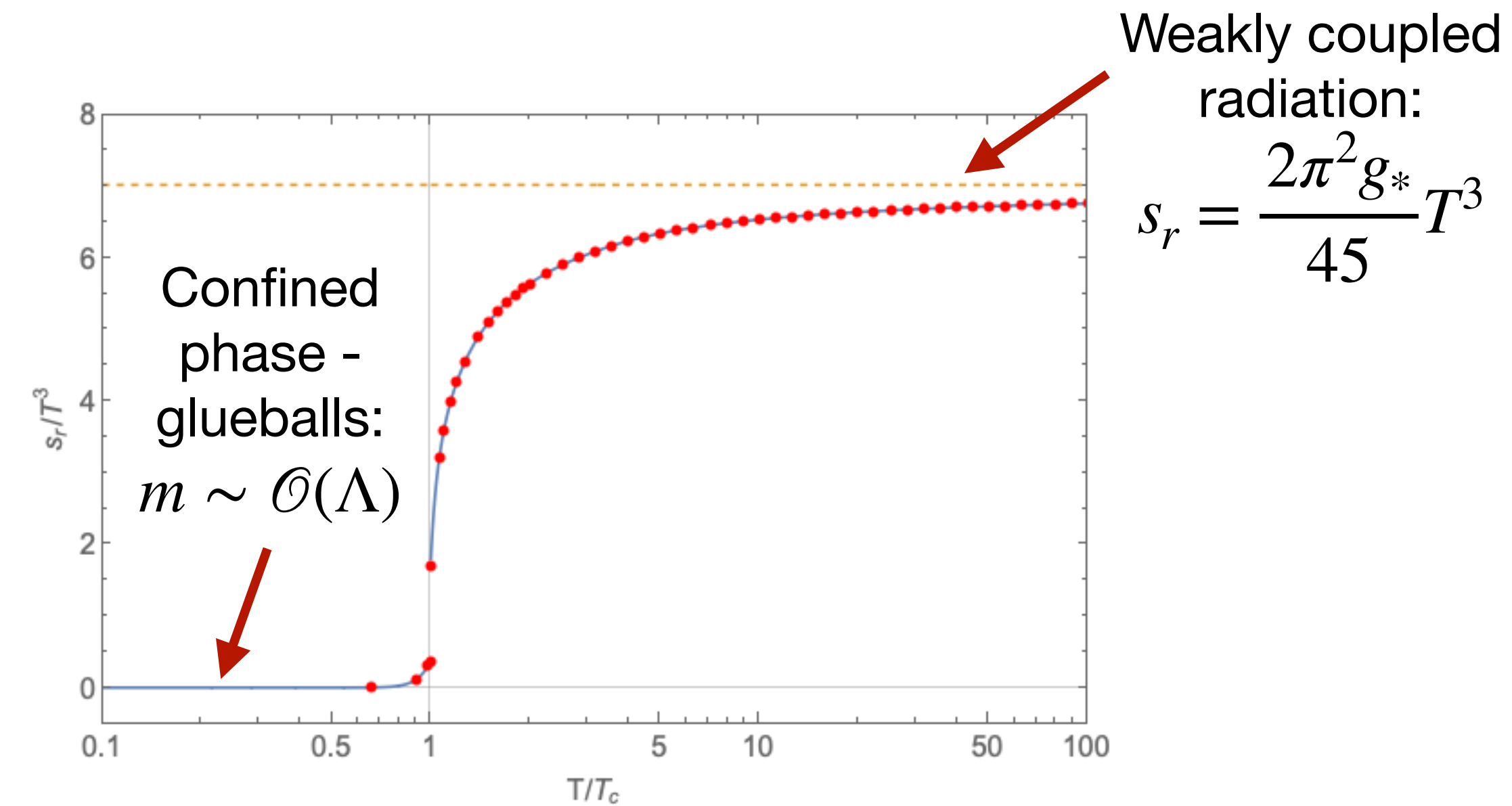
# Why large temperatures?

$$\ddot{\phi} + (3H + \Upsilon)\dot{\phi} + V_\phi \simeq 0$$

$$\dot{e}_r + 3H(e_r + p_r) \simeq \Upsilon\dot{\phi}^2$$

$$c_r(T)\dot{T} + 3H[e_r(T) + p_r(T)] \simeq \Upsilon\dot{\phi}^2$$

Parametrize  
 $e_r, p_r$  by  $T$



Entropy density of pure  $SU(3)$  measured on lattice  
[Giusti, Pepe: 1612.00265]

# Why large temperatures?

$$\ddot{\phi} + (3H + \Upsilon)\dot{\phi} + V_\phi \simeq 0$$

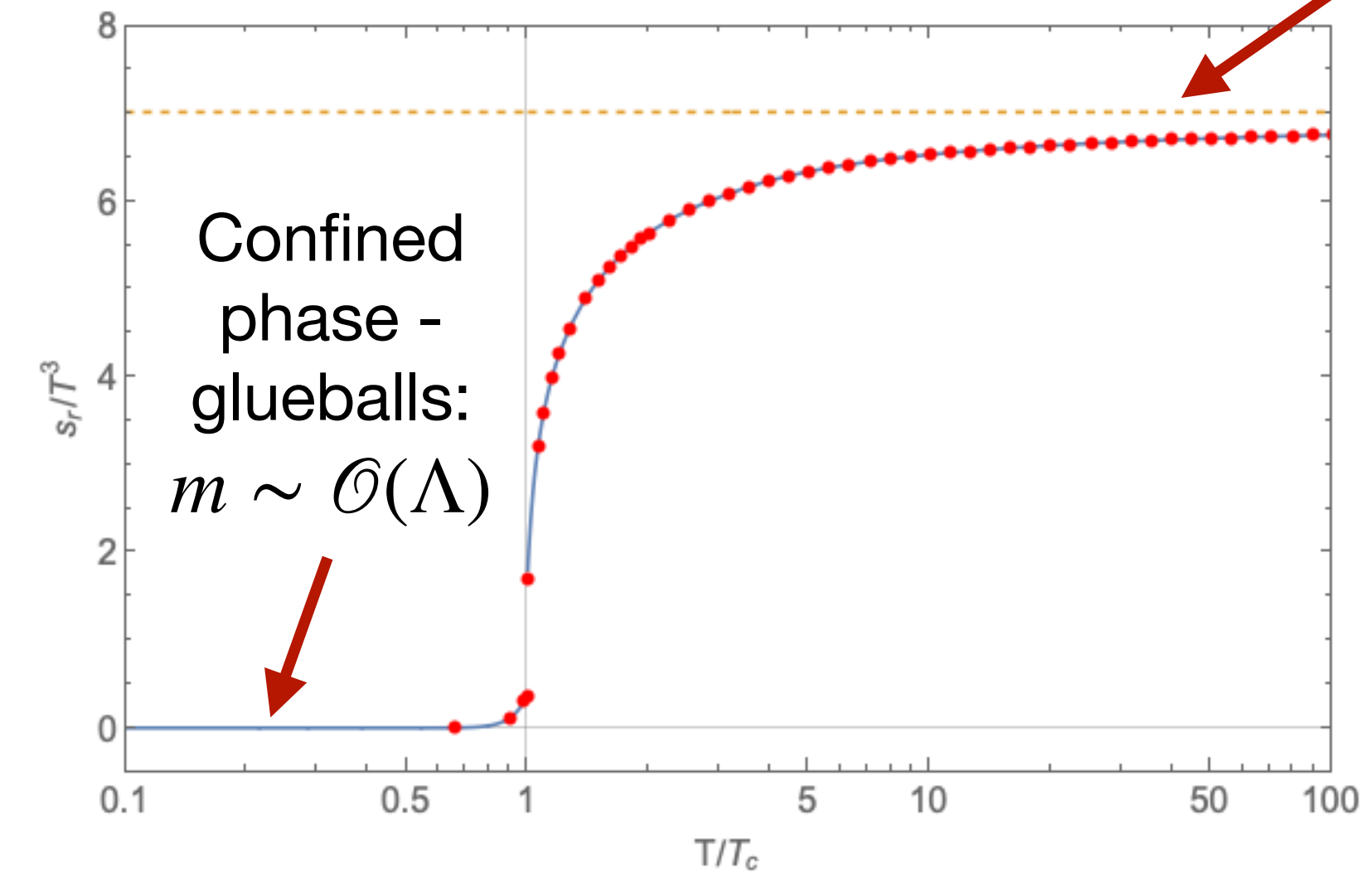
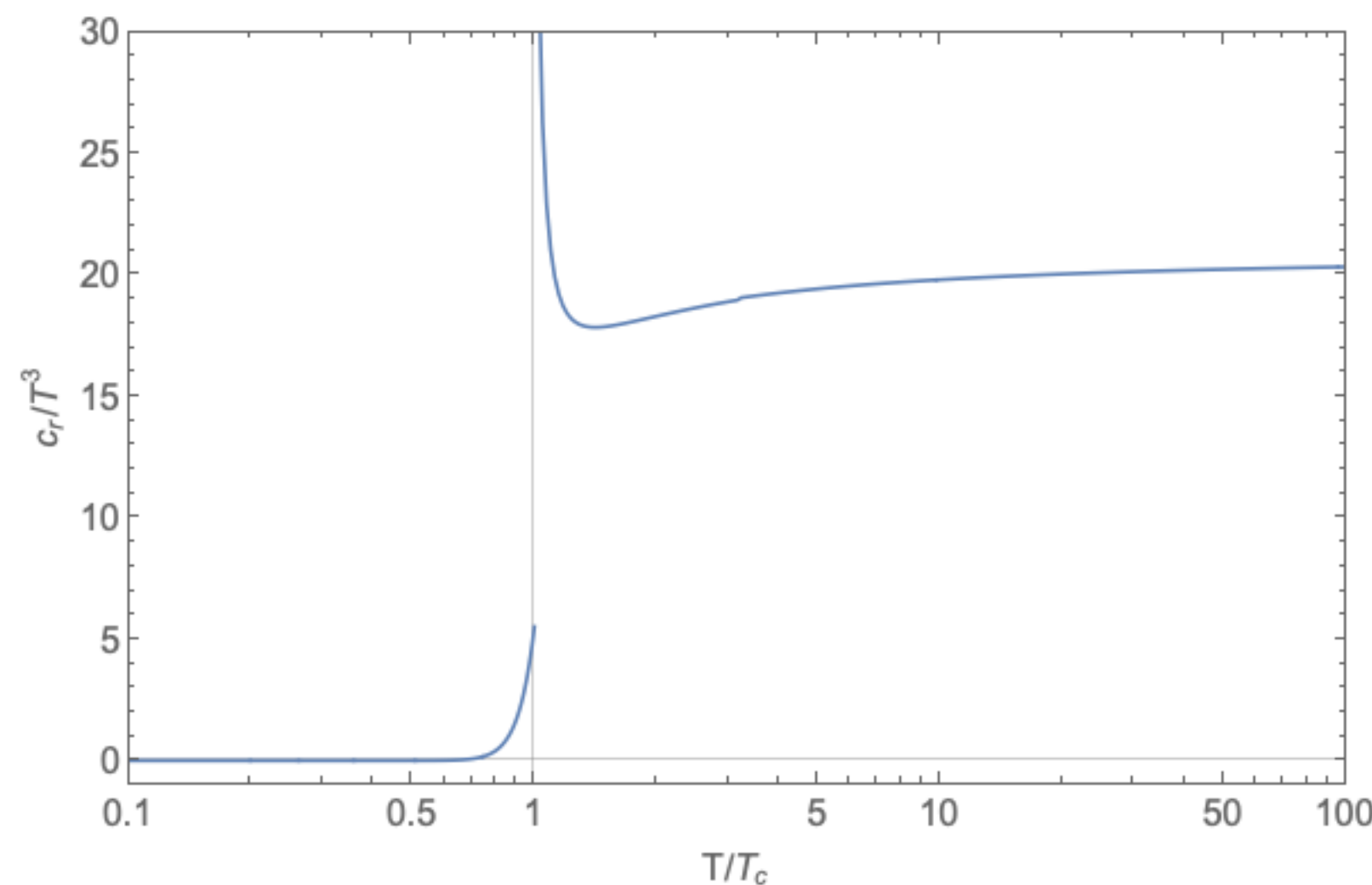
$$\dot{e}_r + 3H(e_r + p_r) \simeq \Upsilon\dot{\phi}^2$$

Parametrize  
 $e_r, p_r$  by  $T$

$$c_r(T)\dot{T} + 3H[e_r(T) + p_r(T)] \simeq \Upsilon\dot{\phi}^2$$

Heat capacity:

$$c_r = \partial_T e_r$$



Weakly coupled  
radiation:

$$s_r = \frac{2\pi^2 g_*}{45} T^3$$

Entropy density of pure  $SU(3)$  measured on lattice  
[Giusti, Pepe: 1612.00265]

Fit of lattice data for  $s_r$ ,

$$s_r = \partial_T p_r$$

$$e_r = Ts_r - p_r$$

# Why large temperatures?

$$\ddot{\phi} + (3H + \Upsilon)\dot{\phi} + V_\phi \simeq 0$$

$$\dot{e}_r + 3H(e_r + p_r) \simeq \Upsilon \dot{\phi}^2$$

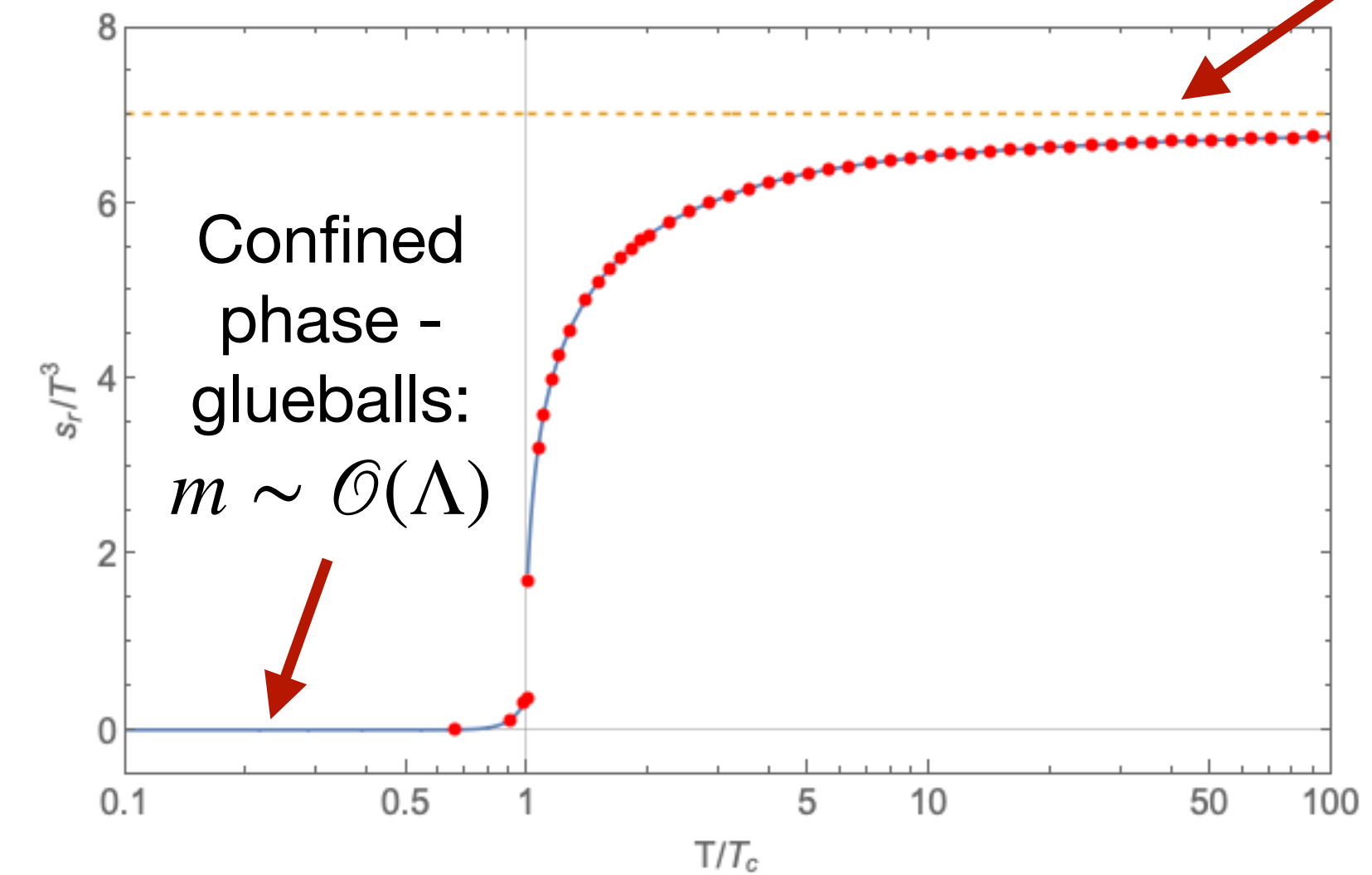
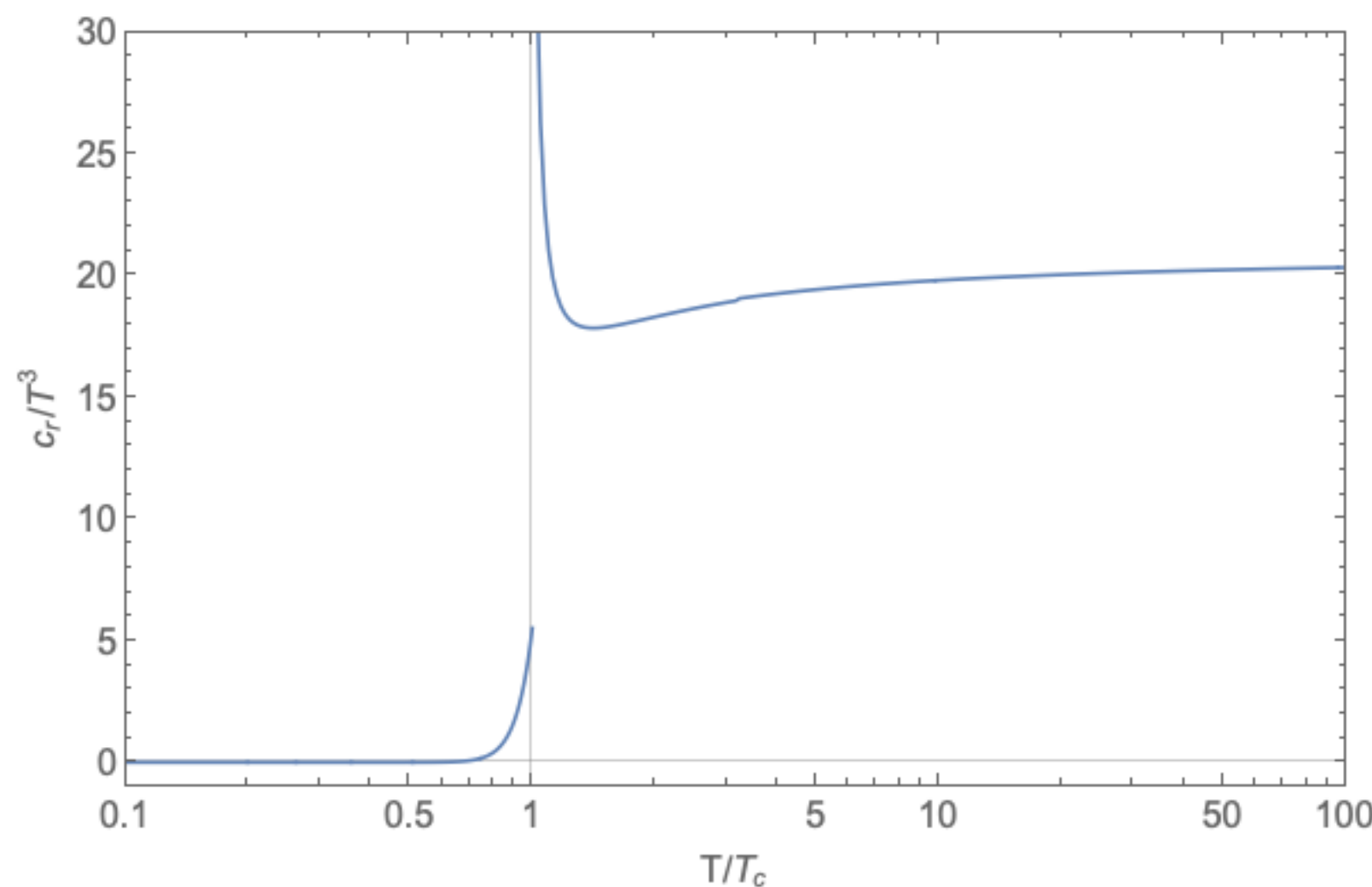
Parametrize  
 $e_r, p_r$  by  $T$

$$c_r(T) \dot{T} + 3H[e_r(T) + p_r(T)] \simeq \Upsilon \dot{\phi}^2$$

Heat capacity:

$$c_r = \partial_T e_r$$

$c_r$  exponentially small well below  $T_c \Rightarrow$  rapid temperature growth!



Weakly coupled radiation:

$$s_r = \frac{2\pi^2 g_*}{45} T^3$$

Entropy density of pure  $SU(3)$  measured on lattice  
[Giusti, Pepe: 1612.00265]

Fit of lattice data for  $s_r$ ,

$$s_r = \partial_T p_r$$

$$e_r = Ts_r - p_r$$