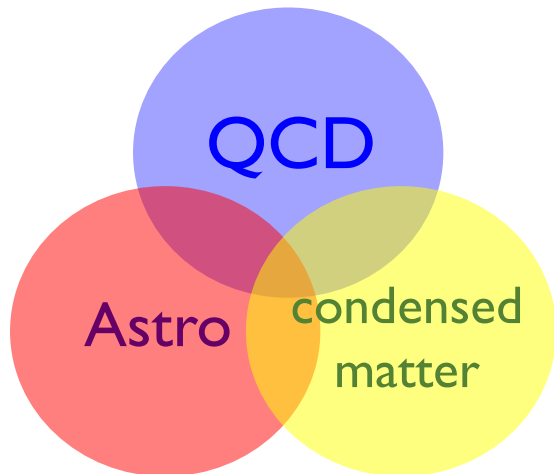


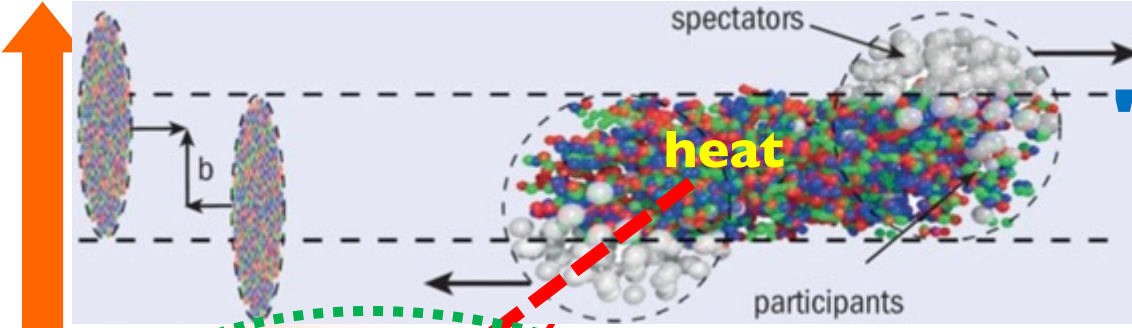
Neutron stars & multi-messenger physics

Toru Kojo

(**Tohoku Univ.**)



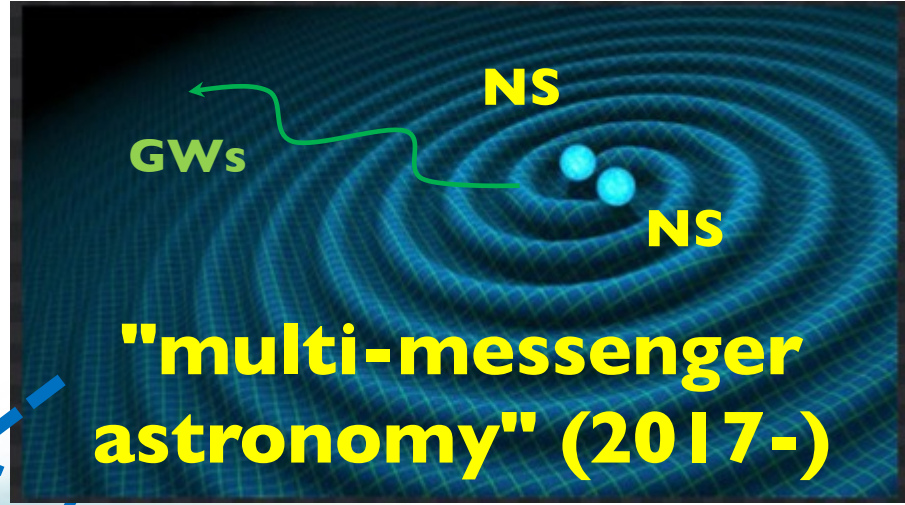
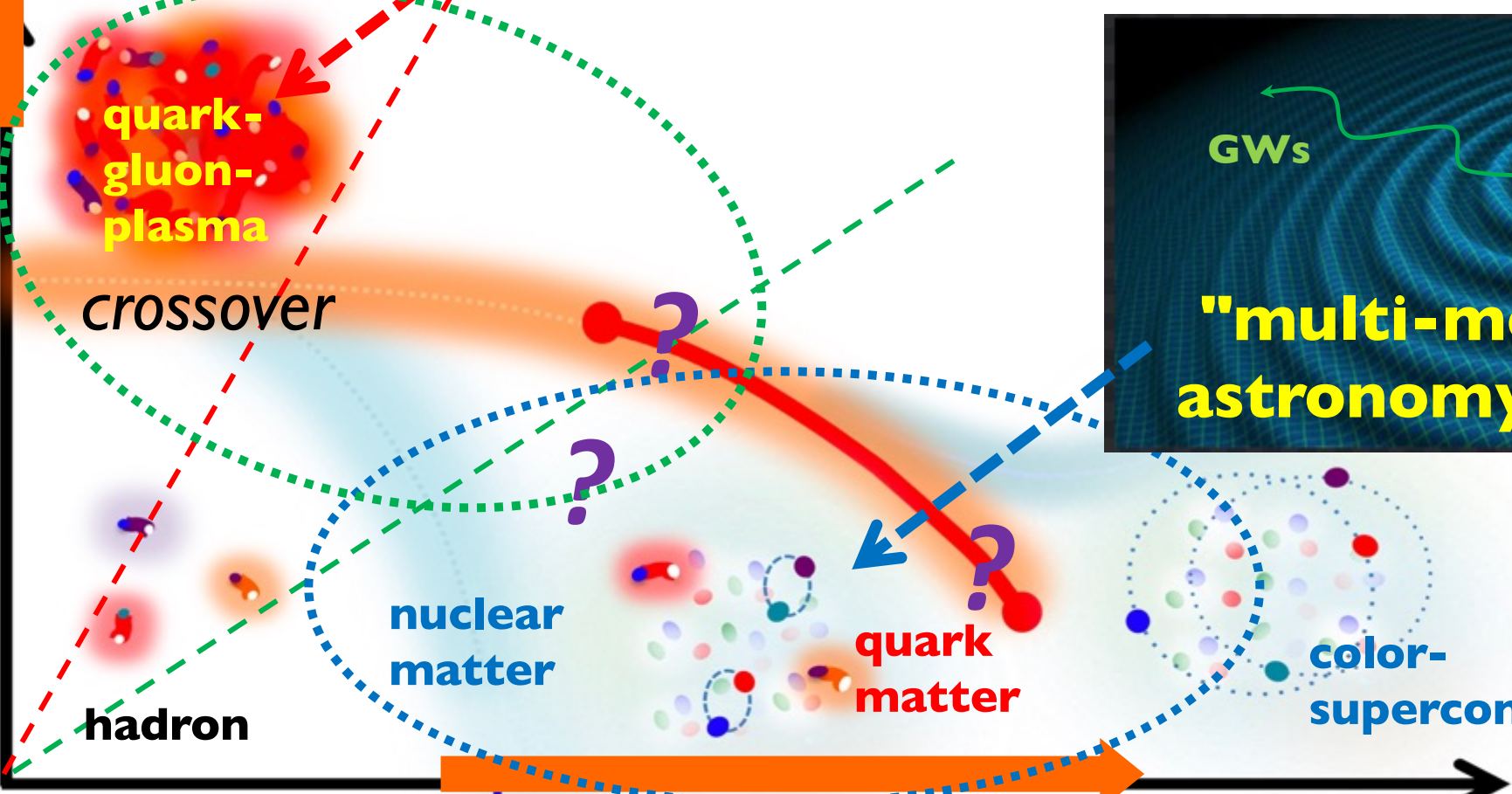
lattice QCD



"heavy-ion collisions"
('80s-)

T

 $\sim 155 \text{ MeV}$



"multi-messenger astronomy"
(2017-)

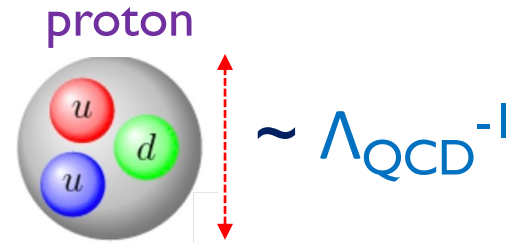
M_N **Neutron stars ('30s-)**

μ_B

Scales & units to be used

dynamical scale

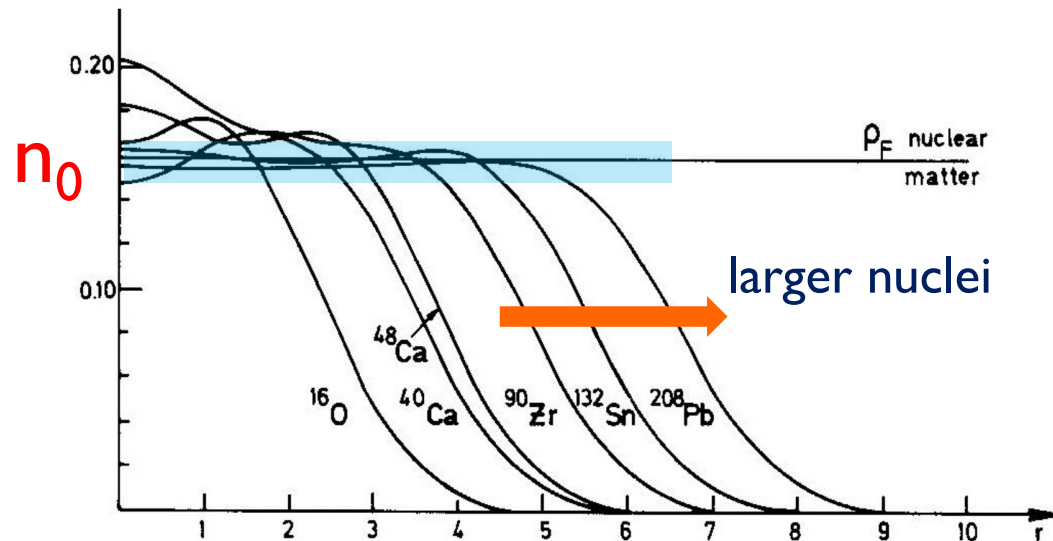
- $\Lambda_{\text{QCD}} \sim 200 \text{ MeV} \sim 1 \text{ fm}^{-1}$



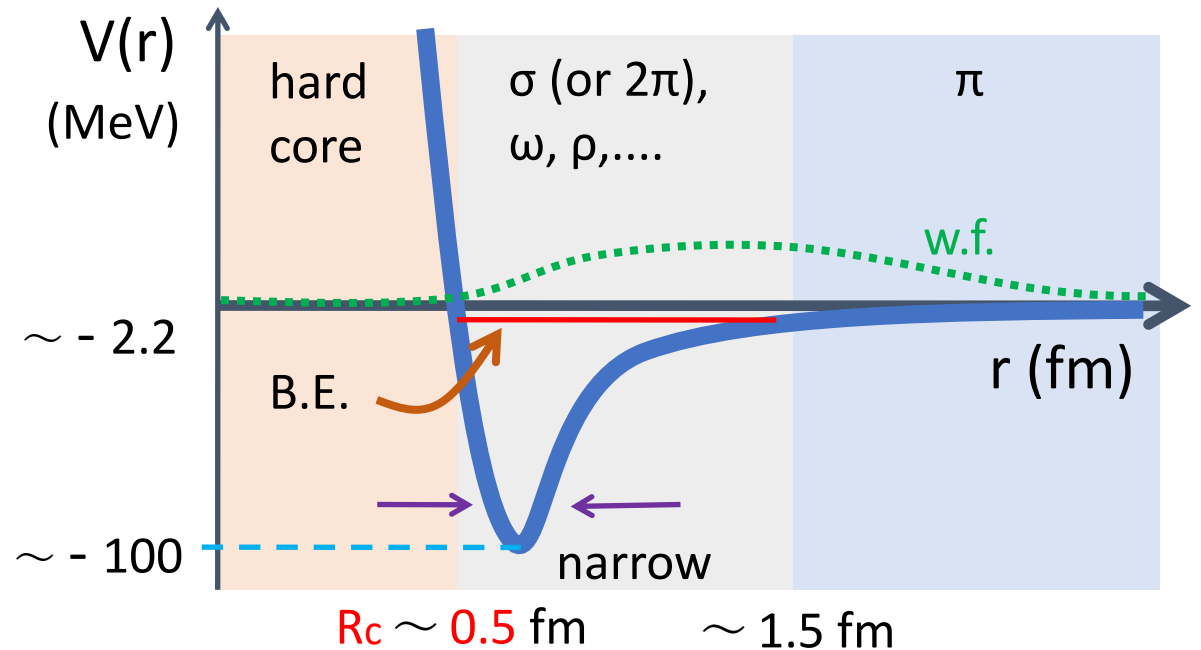
mass

$$M_N \sim N_c \Lambda_{\text{QCD}} \sim 1 \text{ GeV}$$

- saturation density: $n_0 = 0.16 \text{ fm}^{-3}$
(\sim central density of nuclei)



- nuclear binding E : $O(1-10) \text{ MeV}$
(unnaturally small !!)

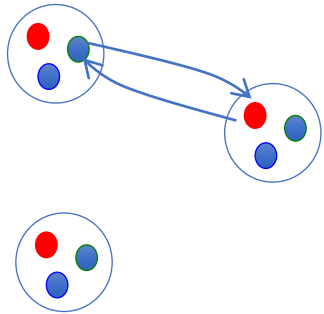


State of matter: **overview**

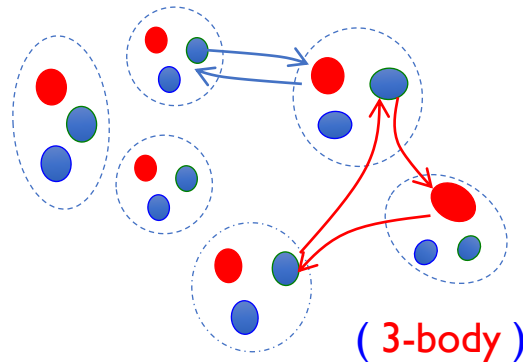
$(n_0 = 0.16 \text{ fm}^{-3})$

[Masuda+ '12; TK+ '14]

- few meson exchange
- nucleons **only**

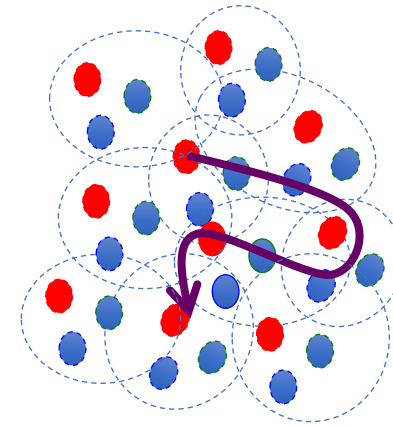


- many-quark exchange
- structural change,...
- hyperons, Δ , ...



most difficult
(d.o.f ??)

- Baryons overlap
- Quark Fermi sea



strongly correlated
(d.o.f : quasi-particles??)

not explored well

➔
(pQCD)

[Freedman-McLerran, Kurkela+, Fujimoto+...]

ab-initio nuclear cal.
laboratory experiments

steady progress

$\sim 1.4 M_{\odot}$

$\sim 2 M_{\odot}$

n_B

$\sim 2n_0$

Hints from NS

$\sim 5n_0$

$\sim 40n_0$



Lect 1) Overview: NS phenomenology & QCD

- glancing at NS properties
- M-R relation and EOS
- $R_{1.4}$ & low density EOS

Lect 2) NS-NS mergers

- gravitational waves
- pre-mergers [inspiral & tidal deformation]
- post-mergers [EM-counterparts]

Lect 3) From hadrons to quarks in NS

- quark matter
- 3-window modeling
- stiffening of matter in quark-hadron continuity

Plan

Glancing at neutron star properties

Neutron stars (NSs)

rot. period : $P \sim 1 \text{ ms} - 1 \text{ s}$ ($\sim 10^{-9} - 10^{-6} P_{\odot}$)

Mass : $M \sim 1-2 M_{\odot}$

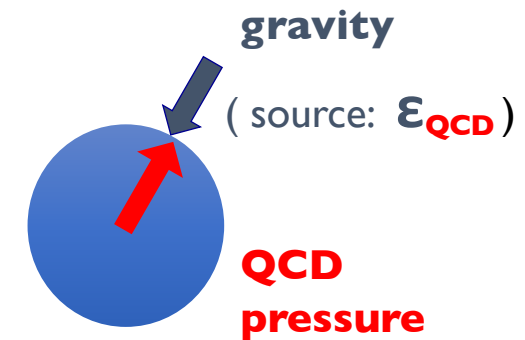
Radius : $R \sim 11-13 \text{ km}$ ($\sim 10^{-6} R_{\odot}$)

n_p/n_B : $Y_p \sim 0.05$ (neutron rich)

Temp. : $T \sim \text{KeV}$ ($\ll p_F$ of nucleons)

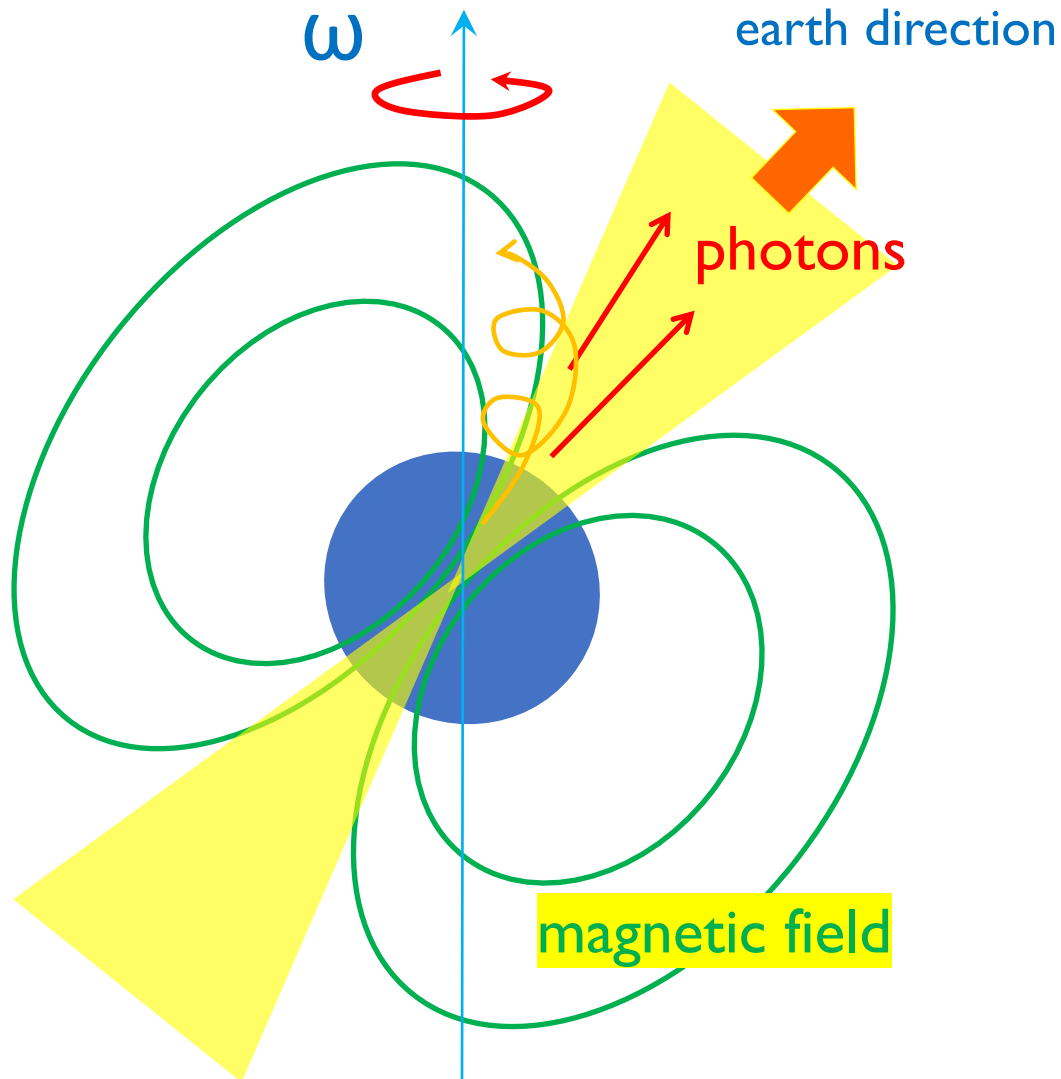
mag. field : $B \sim 10^6 - 10^{15} \text{ Gauss}$ ($\sim 10^5 - 10^{11} B_{\odot}$)

how can we infer
these extreme properties ?



gravitationally bound

fast rotating pulsars & NS



the **existence** of pulsars with
pulse period (for earth) $< 1\text{s}$
 (very quickly rotating, large ω)

with constraints

$$v_{\text{surf}} = \omega R < c \quad (\text{light velocity})$$

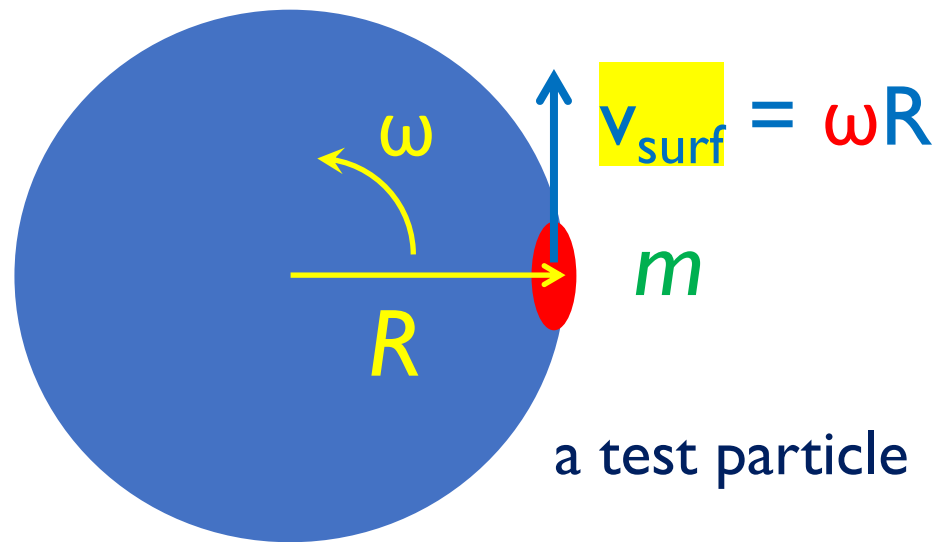
➔ $R \sim 10\text{-}100 \text{ km (!)}$

compact objects !!

fast rotating pulsars & NS

pulse periods are known to be extremely **stable** & **precise**

*how can such **compact** objects be **stable** ?*



$$GMm/R^2 = mR\omega^2 \quad \text{mechanical balance}$$

$$\rightarrow \omega = [GM/R^3]^{1/2} \quad \text{Kepler frequency}$$

$$\text{large } \omega \rightarrow M_{\text{pulsar}} \sim M_{\text{sun}}$$

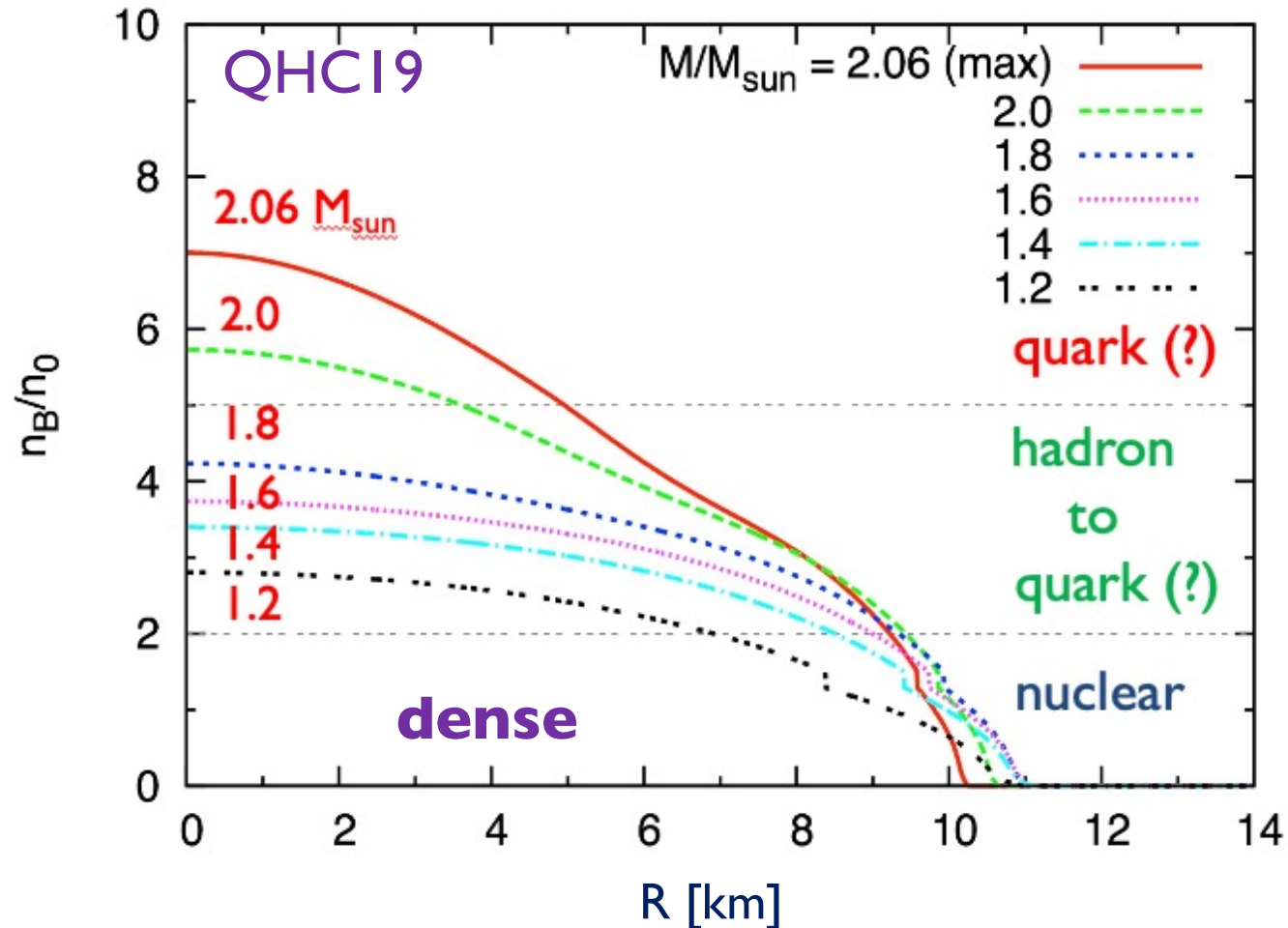
(unless R is extremely small $< \sim 10$ km)

massive & compact \rightarrow **hardly affected by other objects**

How dense?

very dense

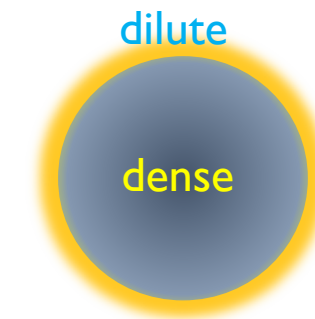
For $M \sim M_{\text{sun}}$ & $R \sim 10 \text{ km}$ $\rightarrow n_B \sim M/(4\pi R^3 \times m_N) \sim O(0.1-1) \text{ fm}^{-3}$



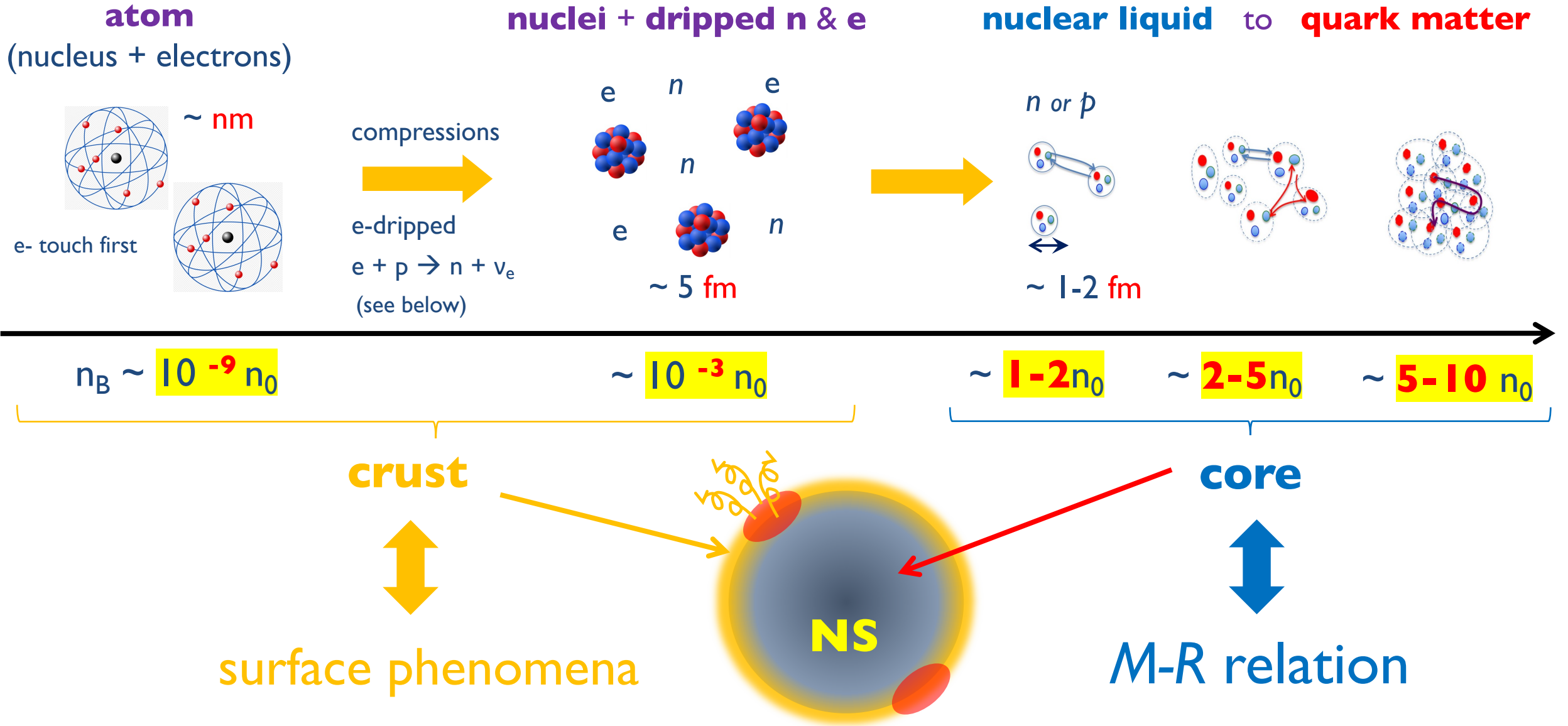
density distribution
from the surface to the center

for $M \sim M_{\text{sun}}$

the majority of matter
is very dense, $n_B > n_0$



gravitational compression & matter content



why “neutron” stars ?

assuming ideal gases of rela. electrons + non-rela. nucleons

charge neutrality $\rightarrow n_p (= Y_p n_B) = n_e$ or $p_F^p = p_F^e$

β -equilibrium $\rightarrow m_N + (p_F^n)^2 / 2M_N = m_N + \underbrace{(p_F^p)^2 / 2M_N}_{\text{small}} + p_F^e + \dots$
 $(\mu_n = \mu_p + \mu_e)$ $(= p_F^p)$

$m_N \gg p_F^p$ large scale

$\rightarrow \underline{(p_F^n)^2 = 2M_N p_F^p + \dots}$ $(p_F^n \gg p_F^p)$ **neutron rich !**

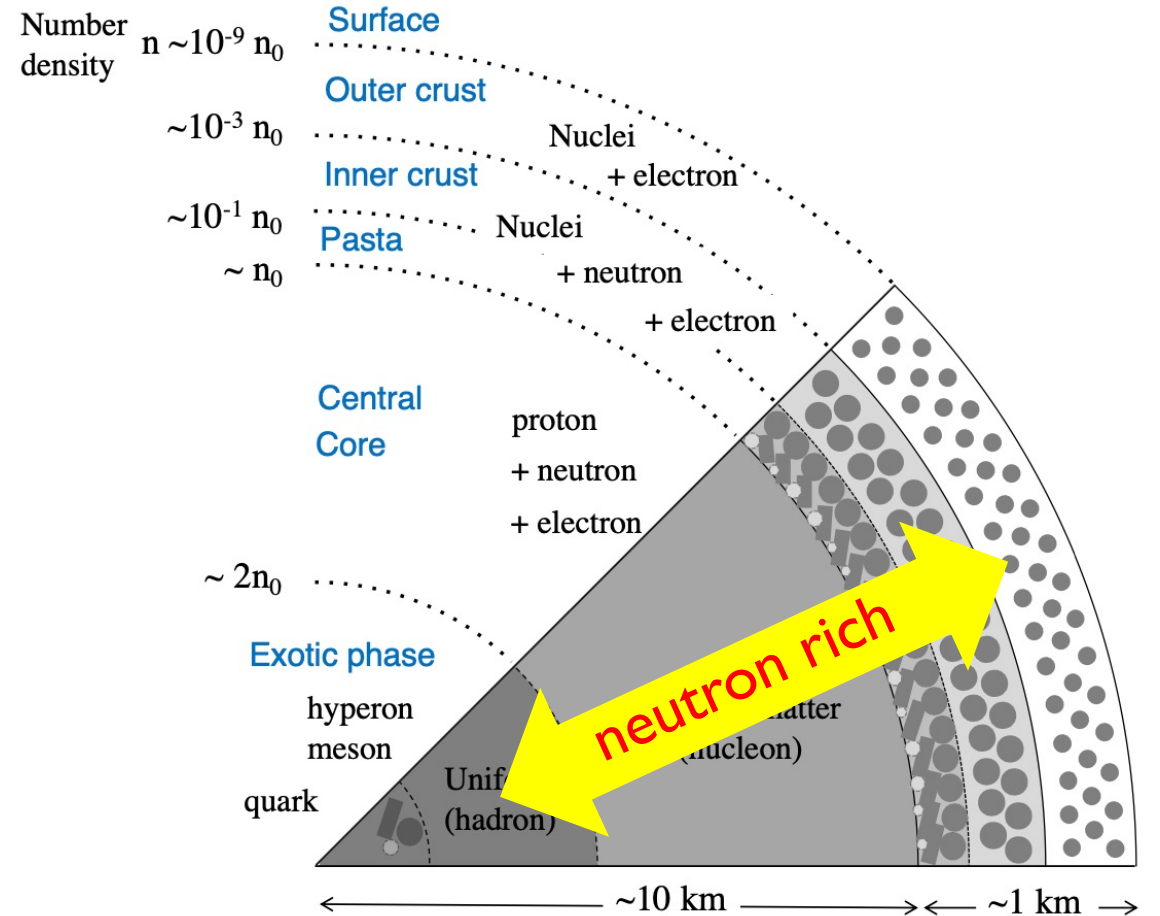
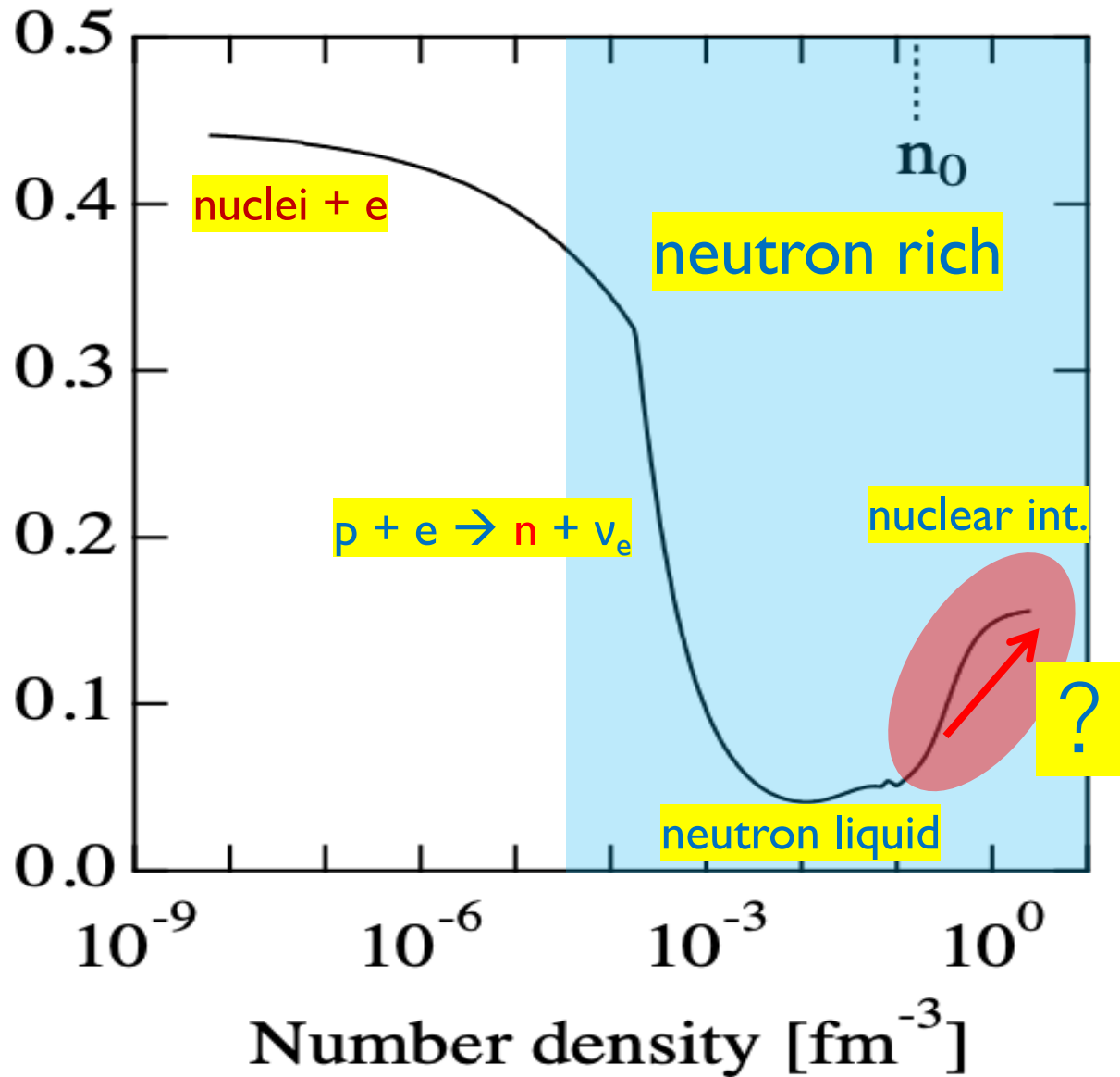
in intuitive terms,

electron kin. energy cost is sensitive to Y_p , while nucleon kin E cost is NOT (tempered by $1/M_N$)

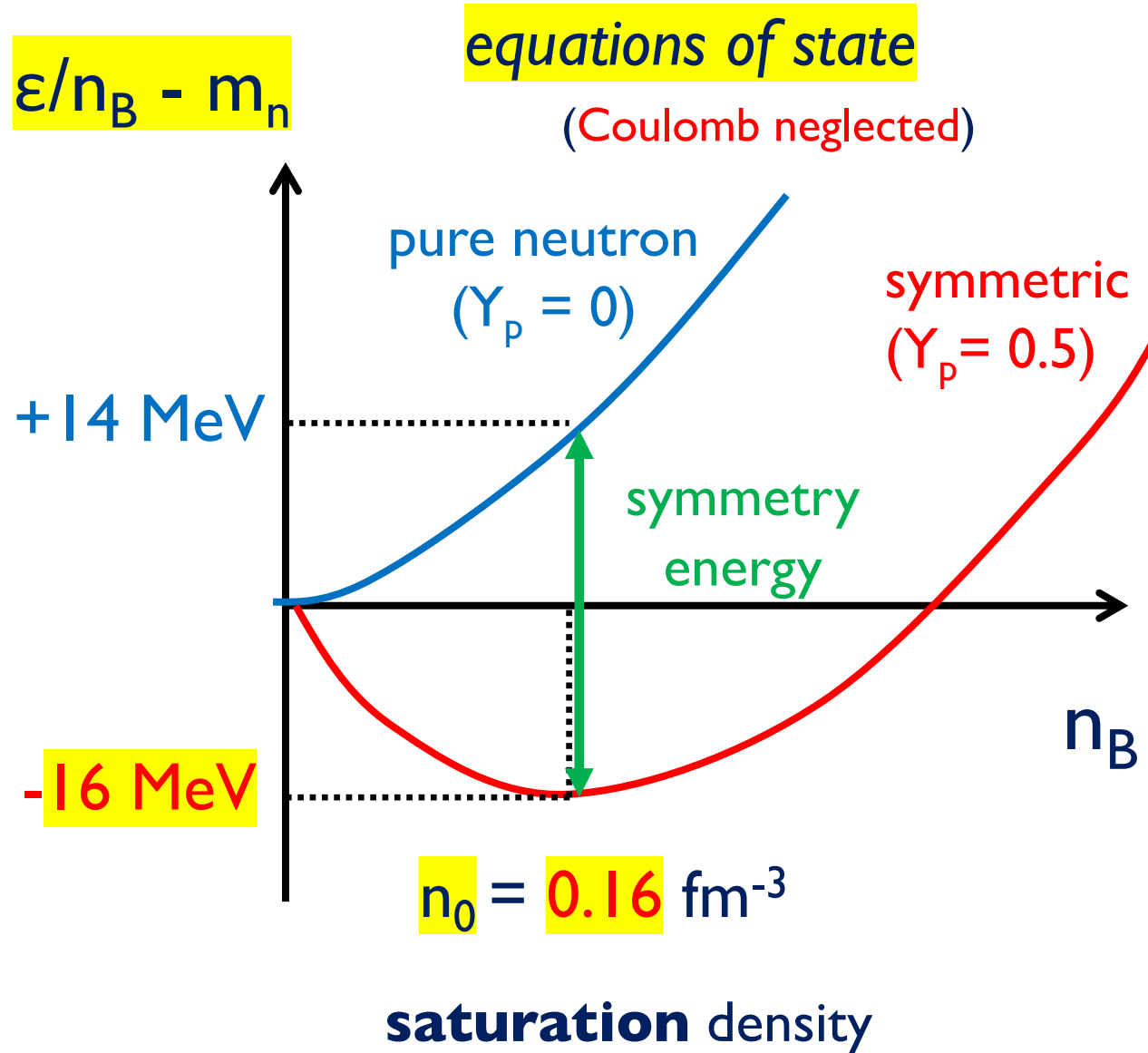
\rightarrow the priority is to reduce electron density as much as possible \rightarrow small Y_p

the trend continues until nucleon interactions become important

how wide neutron rich domain?



Y_p in nuclear sector becomes important at $\sim n_0$



- *nuclear liquid* at $n_B > \sim 0.5n_0$
(clusters are all dissolved)

- **SNM** is **self-bound**

$$P = \mu_B n_B - \epsilon = n_B^2 \partial(\epsilon/n_B)/\partial n_B$$

$$\rightarrow P(n_0) = 0$$

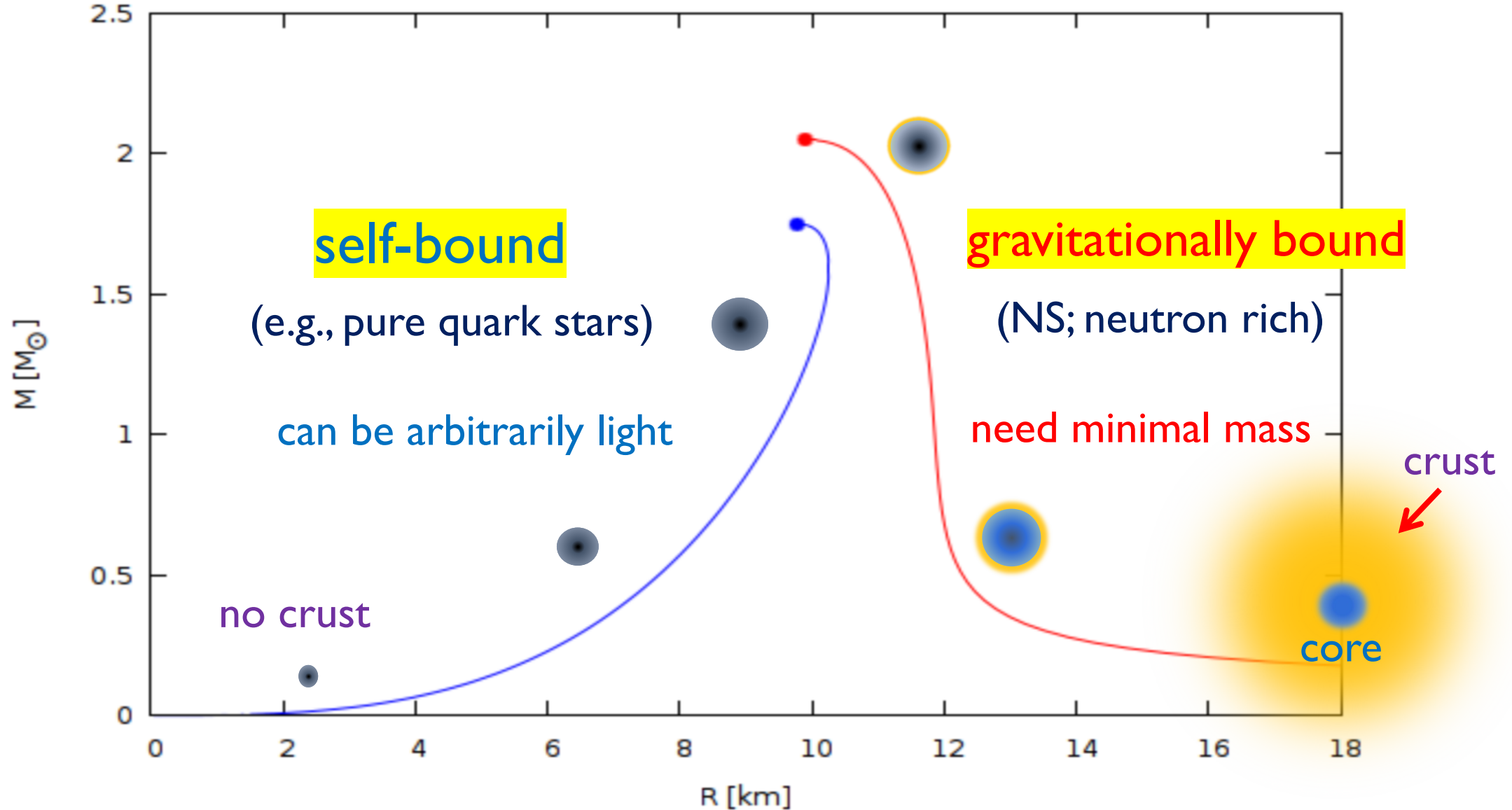
(finite density **without** pressure!)

- **PNM** is **NOT** self-bound

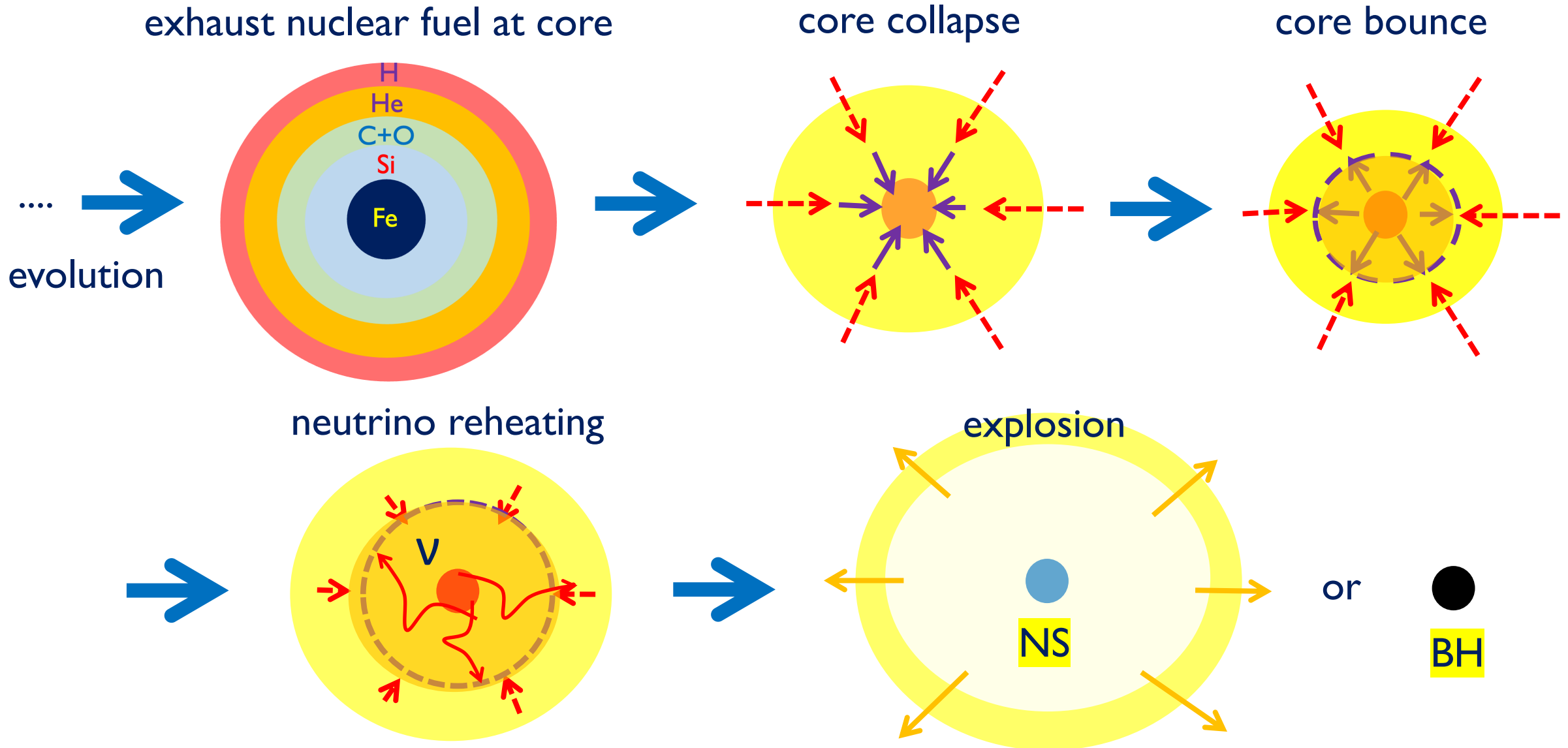
(finite density needs **external pressure**!)

the external force \rightarrow **gravity**

self-bound vs gravitationally bound stars

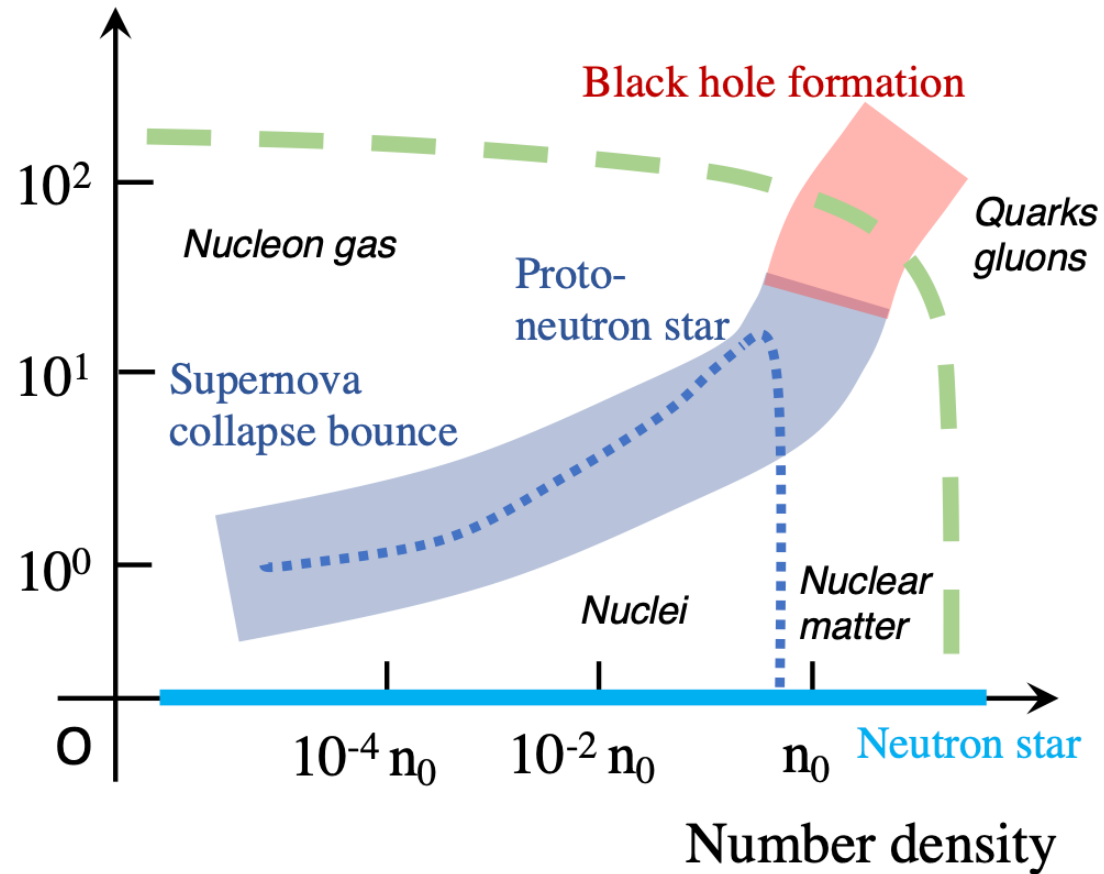


Core collapse supernovae: birth of NS

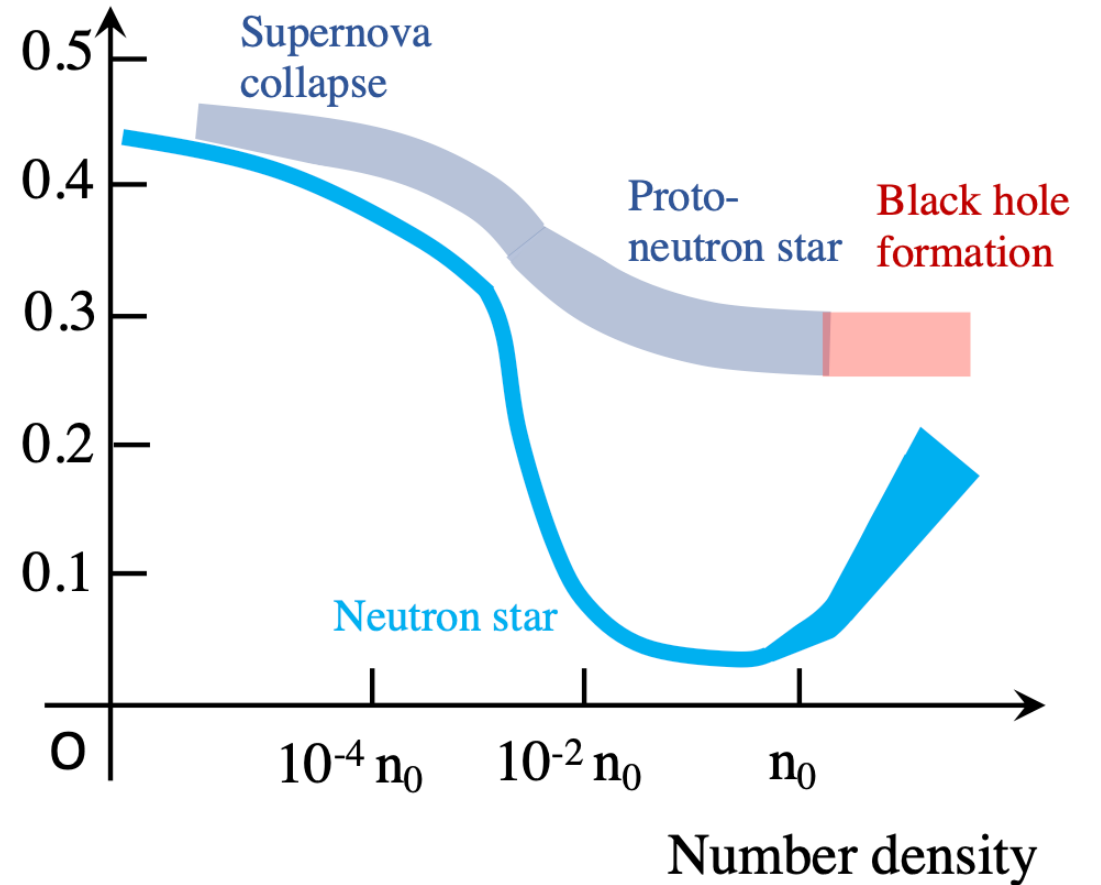


Supernova matter & proto-NS

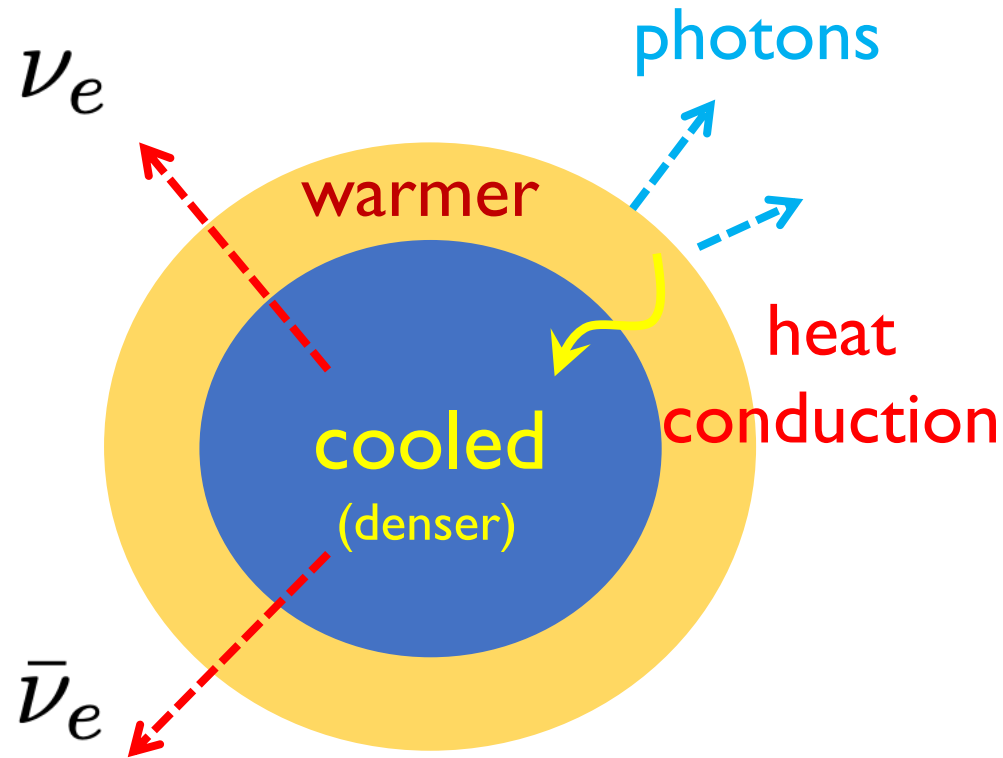
Temperature [MeV]



Proton fraction



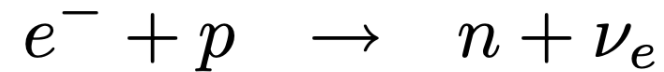
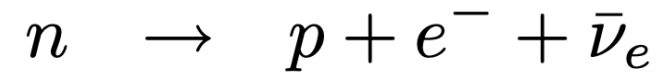
Temperature in NSs



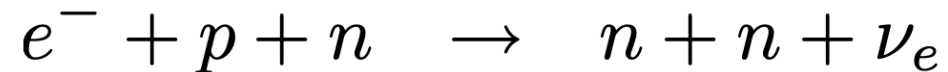
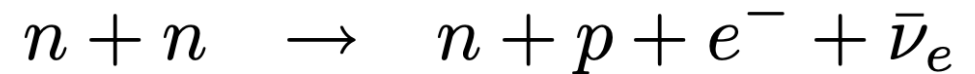
in the **core**:

- **thermal** nucleons lose energy via

direct URCA (Fast):

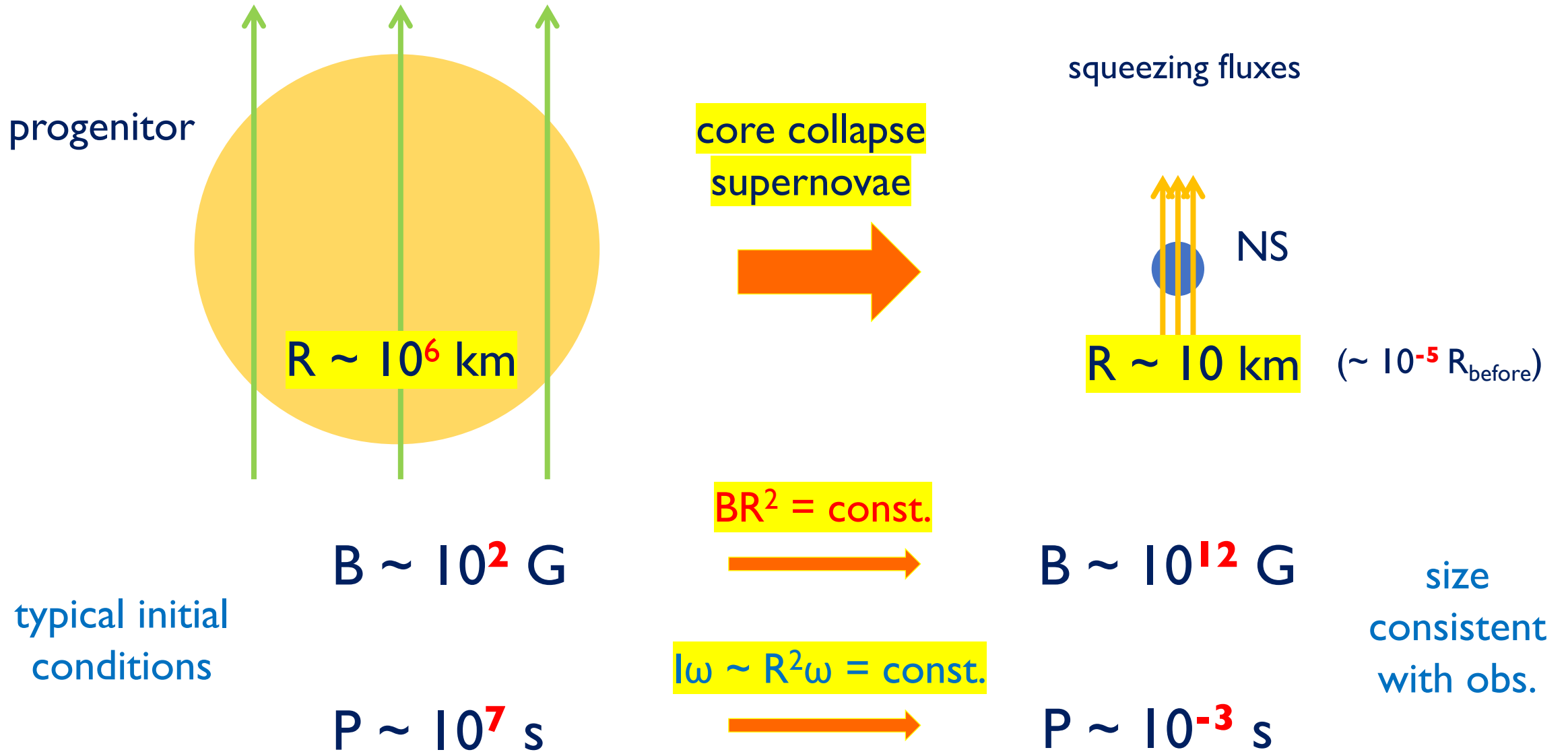


modified URCA (Slow):



~10-100s after the birth, the T_{surf} drops to **~ keV**

magnetic fields & angular momentum



Neutron stars (NSs)

rot. period : $P \sim 1 \text{ ms} - 1 \text{ s}$ ($\sim 10^{-9} - 10^{-6} P_{\odot}$)

Mass : $M \sim 1-2 M_{\odot}$

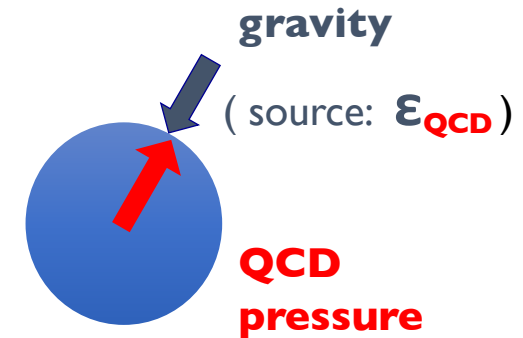
Radius : $R \sim 11-13 \text{ km}$ ($\sim 10^{-6} R_{\odot}$)

n_p/n_B : $Y_p \sim 0.05$ (neutron rich)

Temp. : $T \sim \text{KeV}$ ($\ll p_F$ of nucleons)

mag. field : $B \sim 10^6 - 10^{15} \text{ Gauss}$ ($\sim 10^5 - 10^{11} B_{\odot}$)

now we have (very) rough ideas on NS



gravitationally bound

How many NSs?

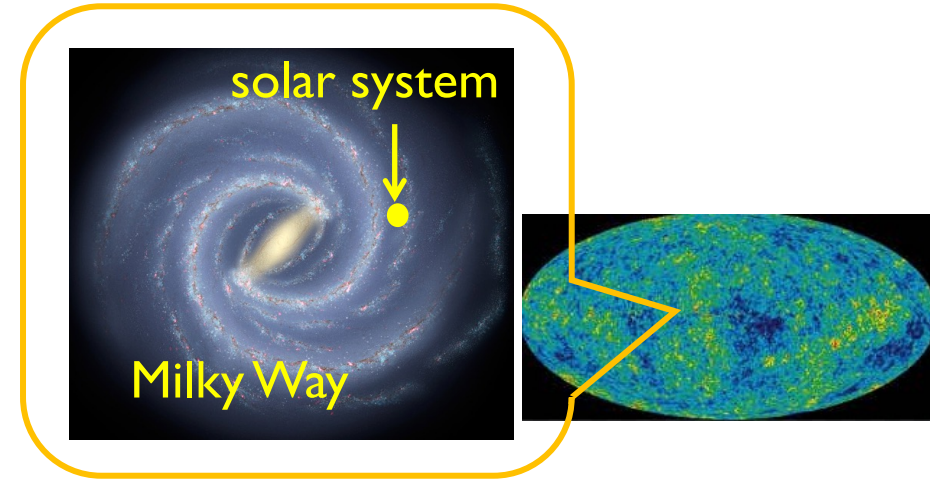
1 parsec (**pc**) = 3.26.. light year

earth to sun: $\sim 10^{-8}$ pc $\sim 10^8$ km

diameter of **Milky Way galaxy**: $\sim 10^4$ pc

range of GW detection (aLIGO): $\sim 10^8$ pc

diameter of **visible universe (CMB)**: $\sim 10^{10}$ pc



In our Milky Way galaxy:

- NS formation (SNe) \rightarrow a few events per century
- age of the galaxy $\rightarrow \sim 10^9$ years

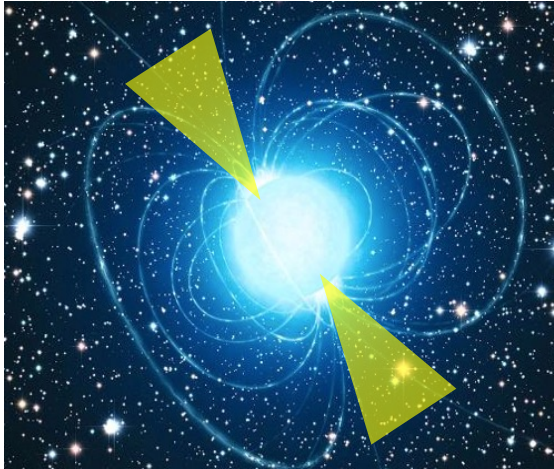


the num. of NSs $\sim 10^8 - 10^9$; only $\sim 2 \times 10^3$ were observed

Types of NSs?

observed $\sim 2 \times 10^3$ pulsars

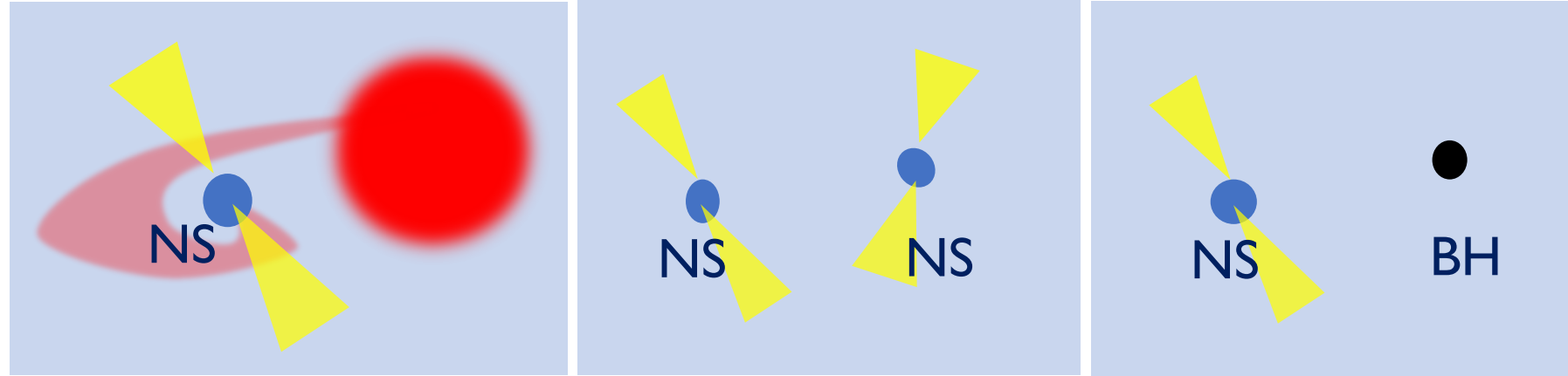
isolated radio pulsars



> 1500 NS

(including ~ 400 millisecc.
pulsars)

pulsars in binaries



~ 100 NSs (pulsars) : member of binaries

... + many invisible NSs (**old**, age $> \sim 10^8$ yrs)

which have exhausted the rot. energy to power radio pulses

EOS & NS structure

NS: brief history --- where are we now?

1934) Zwicky & Baade (prediction) [1931: neutron discovered (Chadwick)]

"*supernova: ordinary star → NS*"

1967) Bell & Hewish: discovery of a pulsar (NS)

"*existence of NS*"

1974) Taylor & Hulse: a binary pulsar (double NSs)

"*indirect confirmation of gravitation waves (GWs)*"

2010) Demorest+: discovery of $2M_{\odot}$ -NS

"*new guide on high density equations of state*"

2017) LIGO-Virgo: GWs from NS mergers (+ electromagnetic counterparts)

"*dawn of multi-messenger (GW, EM, neutrino) astronomy*"

The **first** (established) $\sim 2M_{\odot}$ NS : **PSR J1614-2230**

NS – WD (white dwarf) binary [v1) Demorest+ (2010); v2) Fonseca+ (2016)]

spin period = 2.15 ms; orbital period = 8.7 days

inclination angle = **89.17 ± 0.02** ; edge on

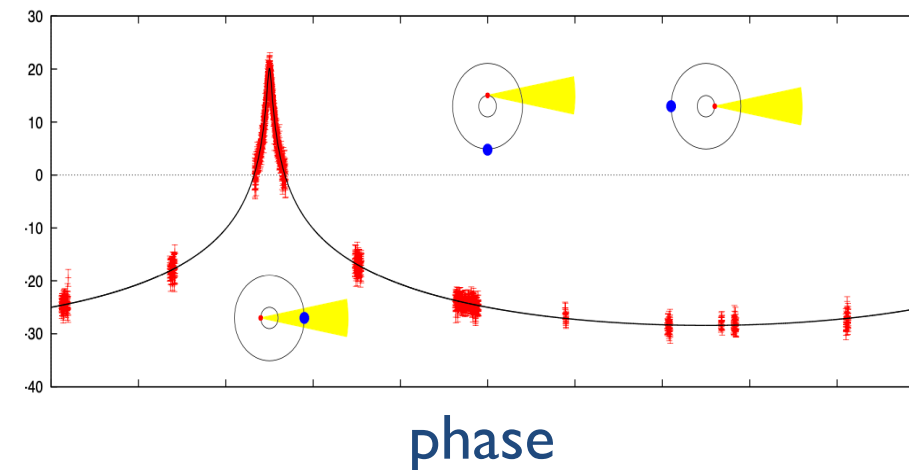
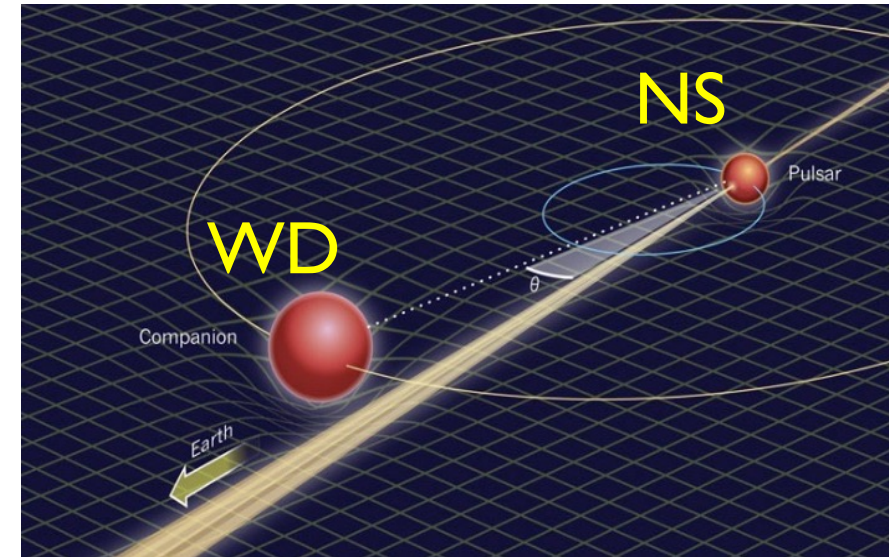
$$M_{\text{NS}} = \mathbf{1.928 \pm 0.017 M_{\odot}}$$

$$M_{\text{WD}} = 0.500 \pm 0.006 M_{\odot}$$

- 1) **binary orbit** \rightarrow relations btw M_{NS} & M_{WD}
- 2) **WD in the direction of pulses (!)**

Shapiro delay of light

1 & 2 \rightarrow separate determination of M_{NS} & M_{WD}



The **heaviest** NS known: **PSR J0740+6620**

NS – WD (white dwarf) binary [v1) Cromartie+ (2019); v2) Fonseca+ (2021)]

spin period = 2.89 ms; orbital period = 4.77 days

inclination angle = **87.56 ± 0.17** ; edge on

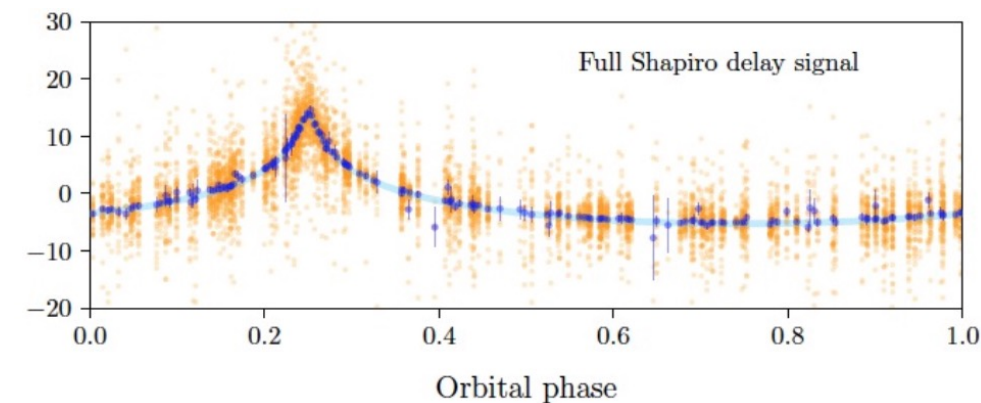
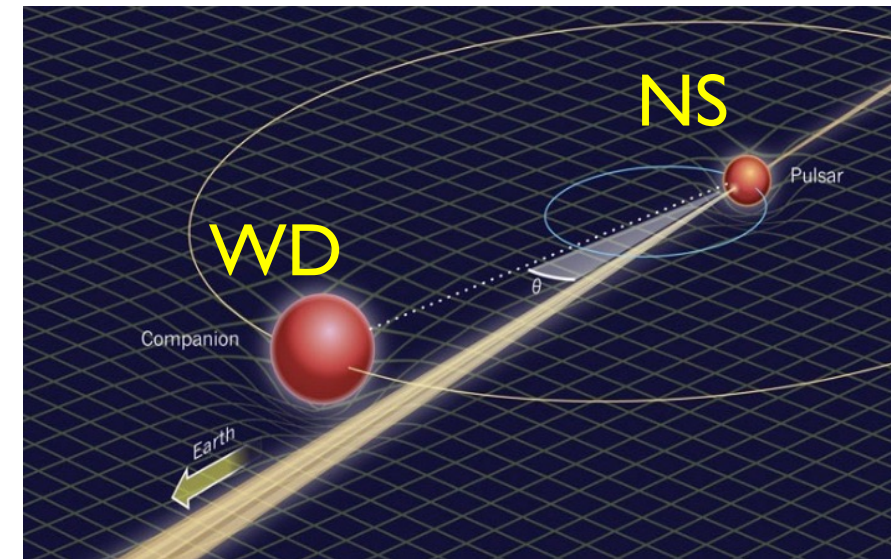
$$M_{\text{NS}} = \mathbf{2.08 \pm 0.07 M_{\odot}}$$

$$M_{\text{WD}} = 0.253 \pm 0.005 M_{\odot}$$

distance from the earth : 1.14 kpc

Also, a target of NICER

M & R (see below)



Mass-Radius (M-R) relations

Tolman-Oppenheimer-Volkoff (TOV) eq.

Einstein eq. **QCD EoS**

$$G_{\mu\nu} = 8\pi G T_{\mu\nu}$$

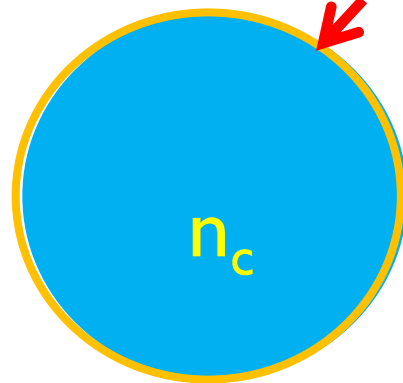
$$\left\{ \begin{array}{l} \frac{dP(r)}{dr} = - \frac{GM(r)\varepsilon(r)}{r^2} \left(1 + \frac{P}{\varepsilon}\right) \left(1 + \frac{4\pi r^3 P}{M(r)}\right) \left(1 - \frac{2GM(r)}{r}\right)^{-1} \\ \frac{dM(r)}{dr} = 4\pi r^2 \varepsilon(r) \end{array} \right.$$

GR effects (> 1)

static
& spherical sym.

$r \rightarrow$ integrated

$$n_c = n_B(r=0)$$

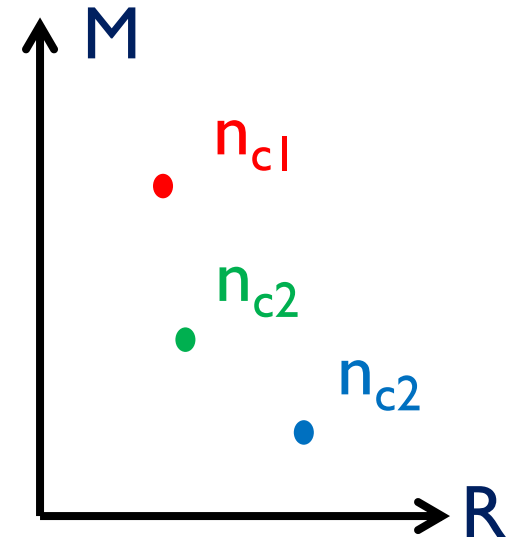


(def. of R)

$$M(n_c)$$

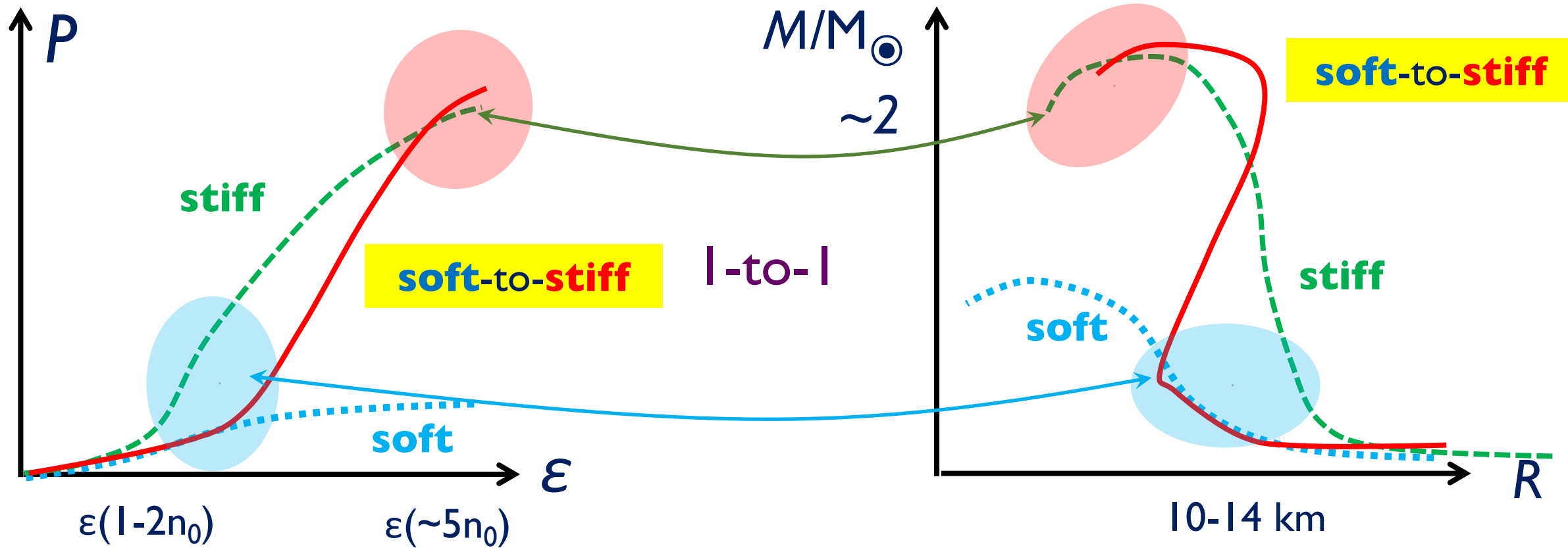
$$R(n_c)$$

at given n_c



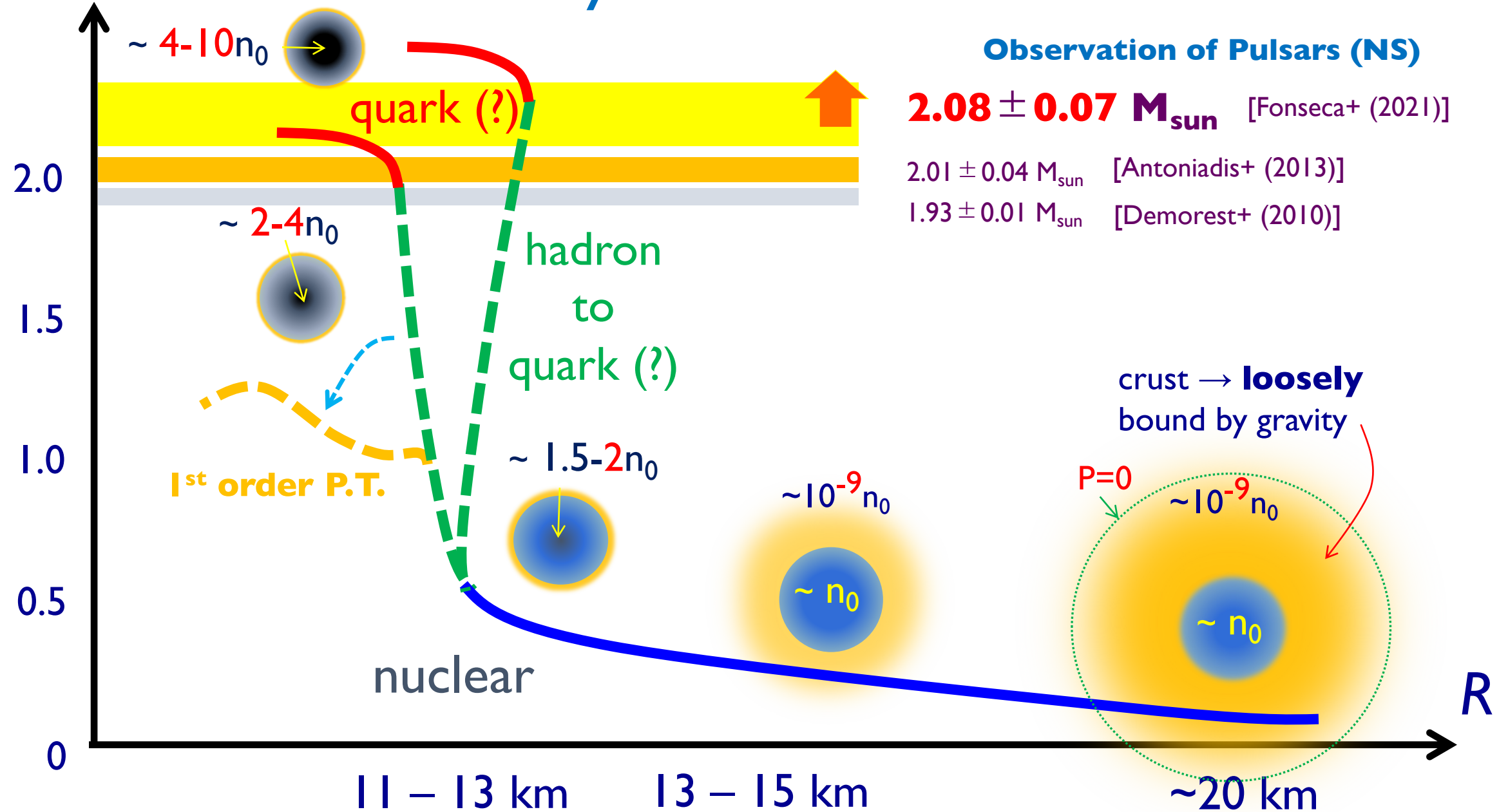
EoS & Neutron Star M-R relation

Ref) Lattimer & Prakash (2001)



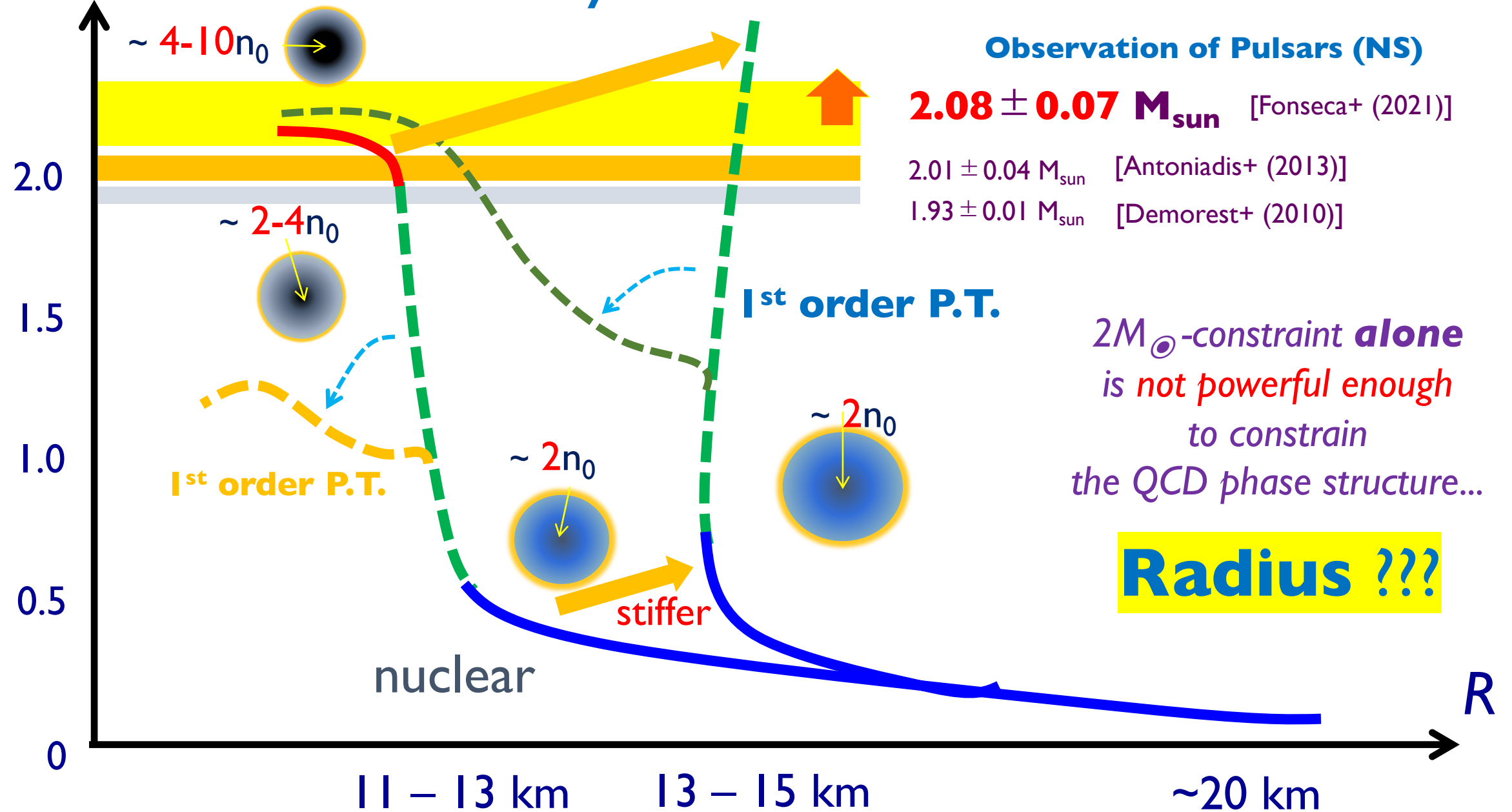
Density vs $M-R$ curves

Ref) Lattimer & Prakash (2001)



Density vs $M-R$ curves

Ref) Lattimer & Prakash (2001)



Observation of Pulsars (NS)

$2.08 \pm 0.07 M_{\text{sun}}$ [Fonseca+ (2021)]

$2.01 \pm 0.04 M_{\text{sun}}$ [Antoniadis+ (2013)]

$1.93 \pm 0.01 M_{\text{sun}}$ [Demorest+ (2010)]

1st order P.T.

1st order P.T.

stiffer

$2M_{\odot}$ -constraint **alone** is *not powerful enough* to constrain the QCD phase structure...

Radius ???

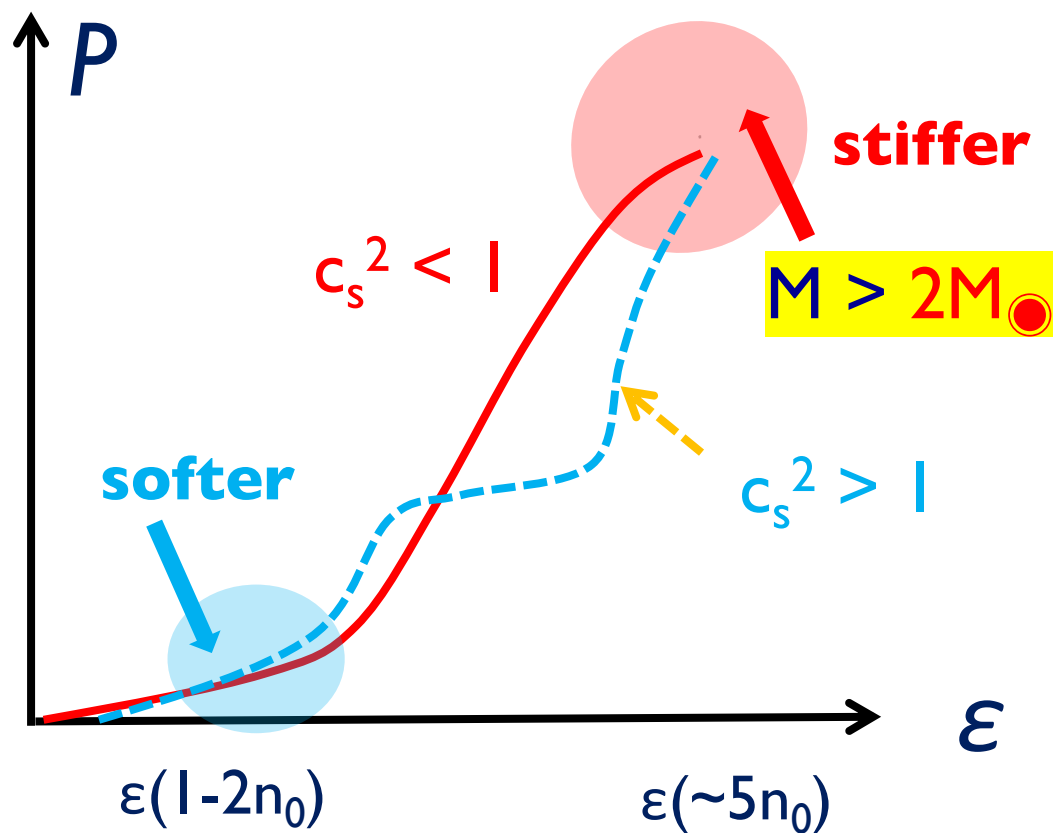
11 – 13 km

13 – 15 km

~20 km

Correlating **low** ($< \sim 2n_0$) & **high** ($> \sim 5n_0$) density EOS

speed of sound: $c_s^2 = dP/d\varepsilon < 1$ (causality)



The constraints are **tighter** if

- **high** density EOS is **stiffer**
- **low** density EOS is **softer**

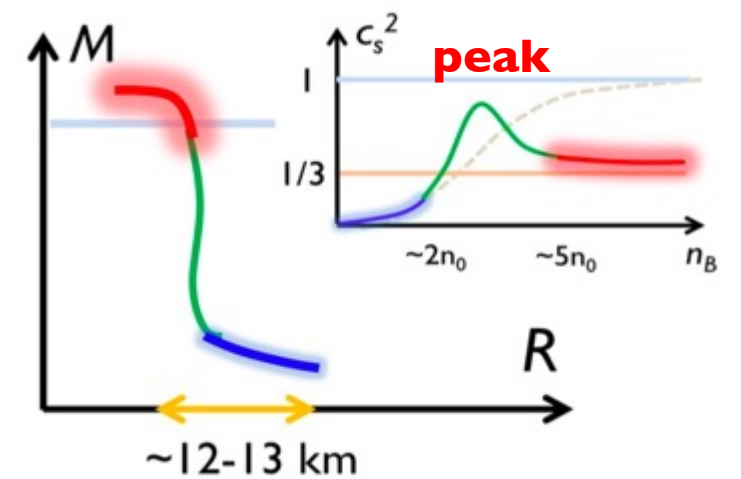
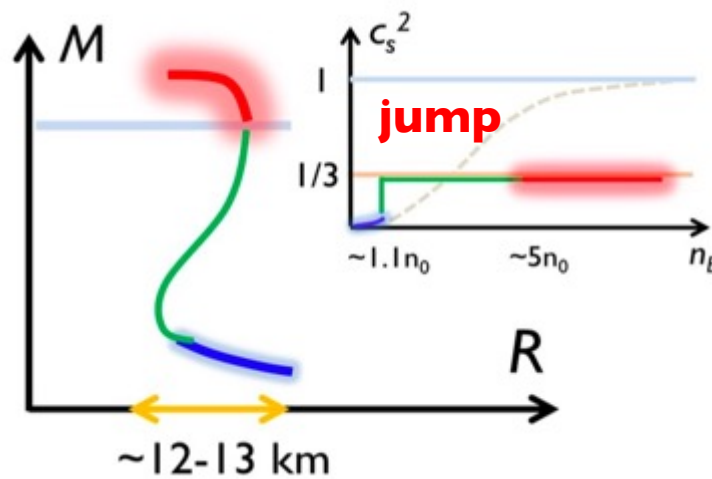
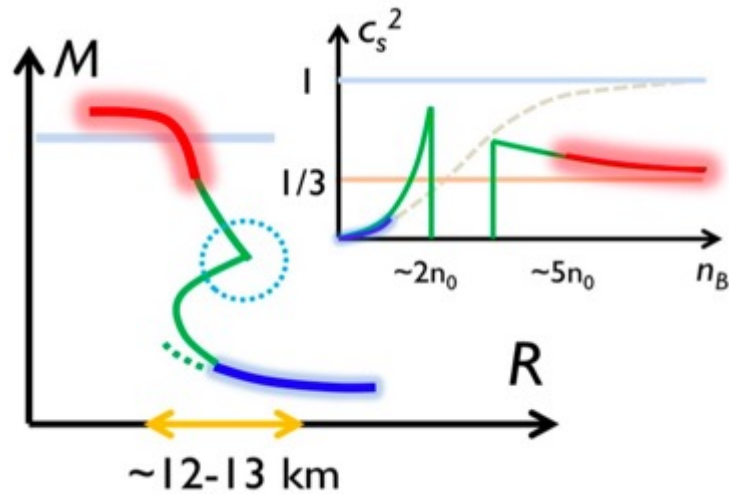
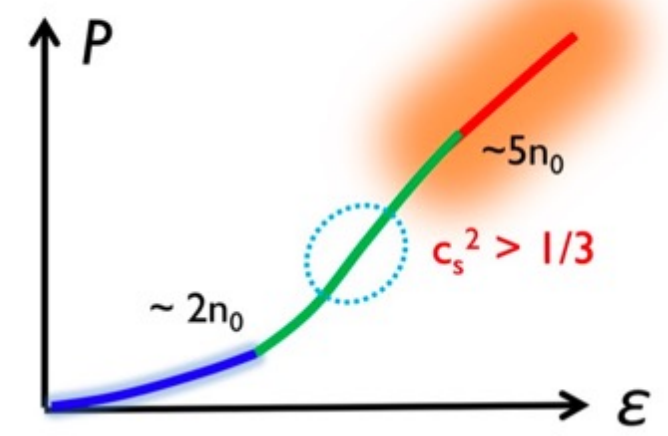
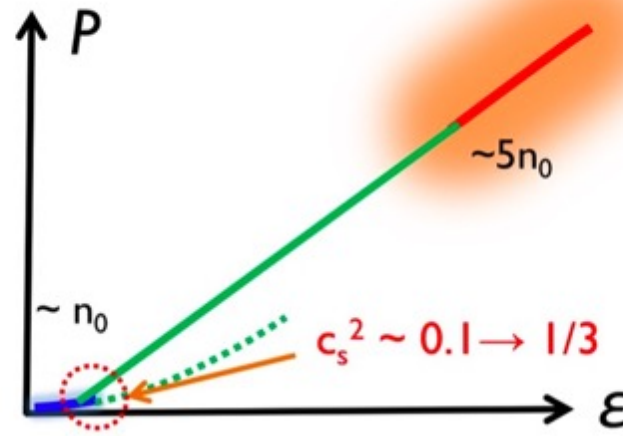
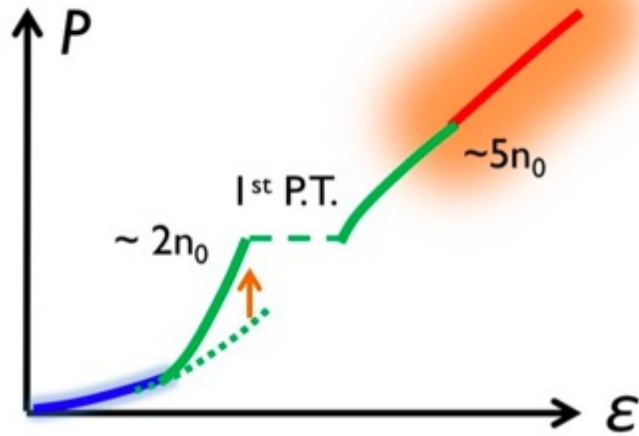
“**soft-to-stiff**” combination is **theoretically challenging**

(but *likely* from observations & nuclear physics)

Several possible scenarios

[e.g. TK '21, mini-review]

→ topics in **Lect. 3**



Key questions to be addressed

- **low density:**
($n_B < \sim 2n_0$)
 - **observational** constraints for **$R_{1.4}$** ?
 - **domain** of **applicability**? ($1.1-2.0n_0$?)
 - **precision** of low density **calculations**?
 - **high density:**
($n_B > \sim 5n_0$)
 - **mechanisms** for **stiff EOS**?
 - effective **d.o.f**? Hadrons or quarks or..?
 - **when** matter becomes **weakly coupled**?
 - **Inbetween:**
 - hadron-to-quark **phase transitions**?
 - **mechanism** for stiffening? **peaks in c_s^2** ?
- } **see below**
- } **Lect.3**

R_{1.4} & low density EOS ($< \sim 2n_0$)

Soft vs **stiff** low density EOS ($n_B < \sim 2n_0$)

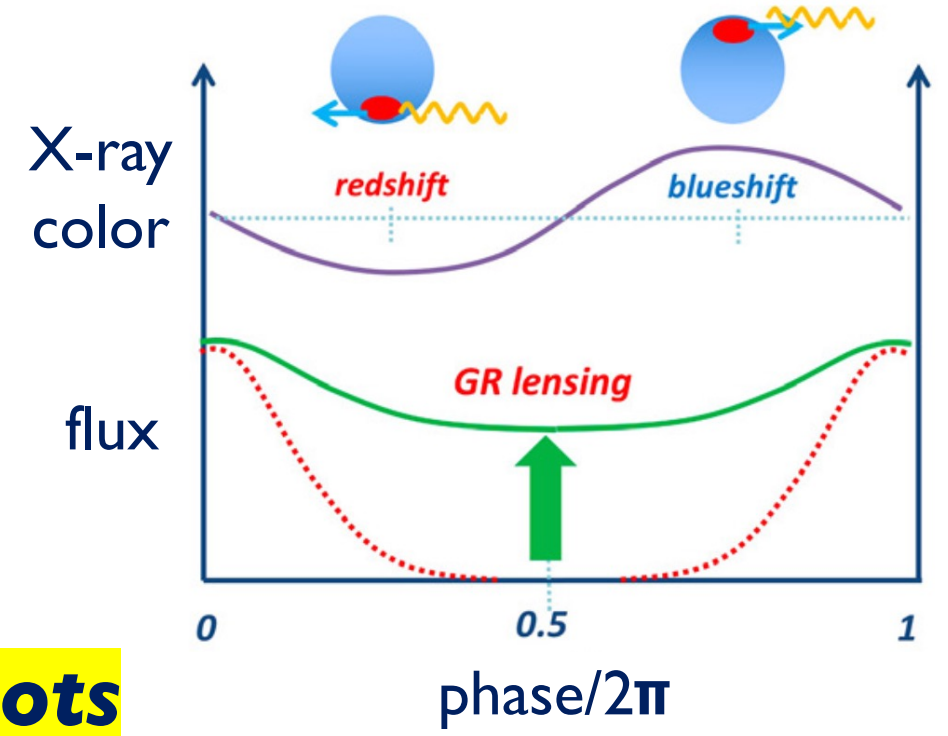
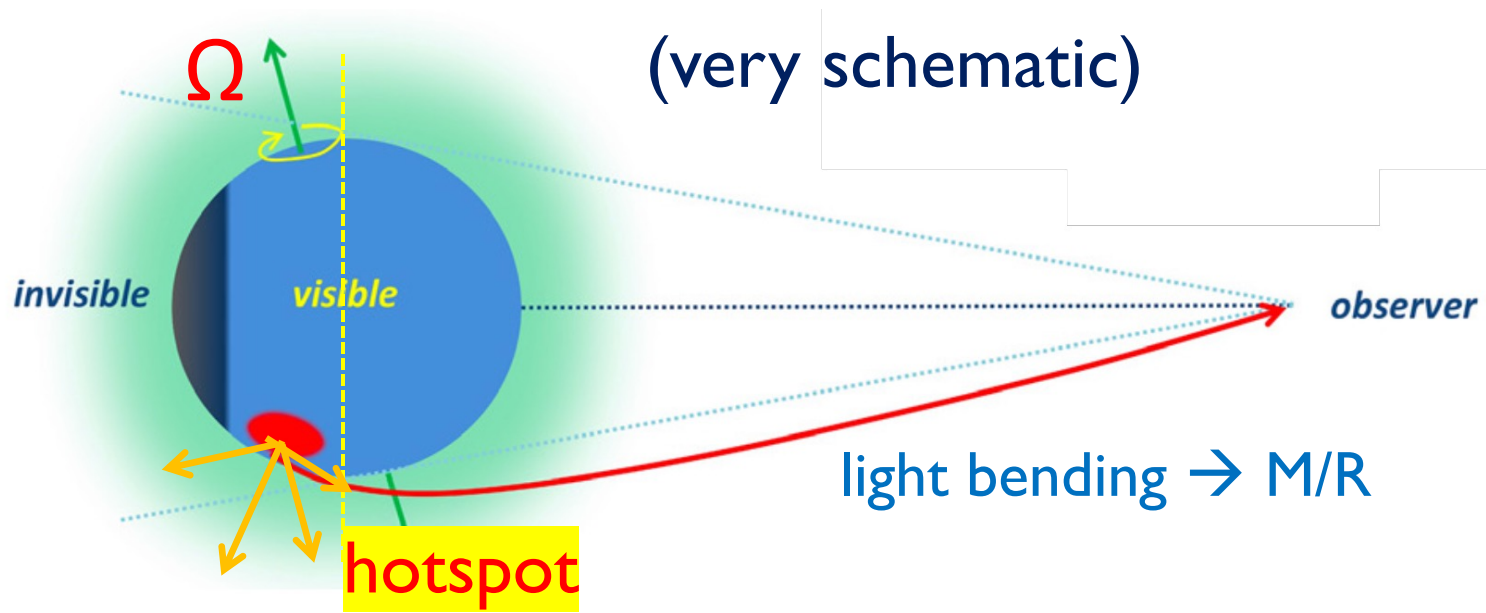
sources of info

NS observations : $R_{1.4} \rightarrow$ EOS at **1-3** n_0
 [Y_p for β -equilibrium]

Nuclear physics : EOS at **< ~ 2** n_0
 [theory for all Y_p , exp. for $Y_p \sim 0.5$]

Heavy Ion experiments : “EOS” for **2-5(?)** n_0
 [but $Y_p \sim 0.5$ & $T \sim 20-100$ MeV,
 considerable extrapolations needed]

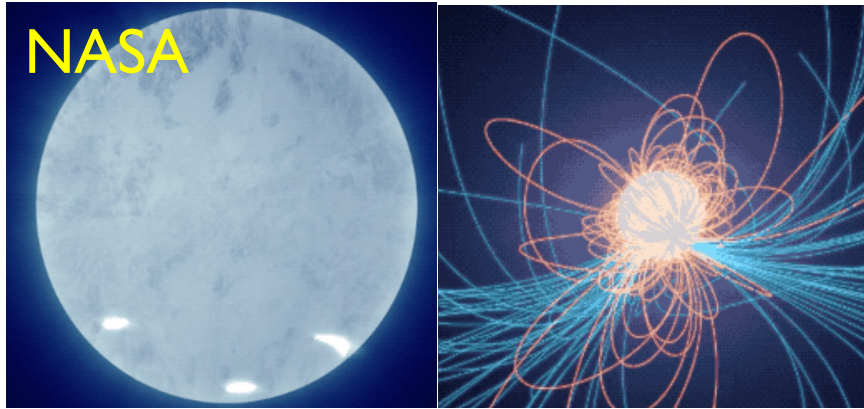
NICER (Neutron star Interior Composition Explorer, 2017-)



idea: **follow hot spots**

period		Ω	(pulse period)
Doppler shifted spectra	\longrightarrow	$R\Omega$	(surface velocity)
GR lensing	(in principle)	M/R	(red shift)

NICER analyses for R (&M)



...reality is more complicated

hot spots →

not simple dipole types
shapes? numbers of spots?

[examined by two teams]

PSR J0030+0451

$$M = 1.44^{+0.15}_{-0.14} M_{\odot}$$

Miller+ '19

$$R = 13.02^{+1.24}_{-1.06} \text{ km}$$

$$M = 1.34^{+0.15}_{-0.16} M_{\odot}$$

Riley+ '19

$$R = 12.71^{+1.14}_{-1.19} \text{ km}$$

PSR J0740+6620

$$M = 2.08^{+0.07}_{-0.07} M_{\odot}$$

NICER + XMM Newton

$$R = 13.7^{+2.6}_{-1.5} \text{ km}$$

Miller+ '21

$$R = 12.39^{+1.30}_{-0.98} \text{ km}$$

Riley+ '21

Nuclear calculations (**microscopic**)

1) Prepare **NN + 3N** forces + ...

- a) meson exchange models (traditional) [e.g., Illinois, Bonn, Argonne, Nijmegen,...]
- b) Chiral EFT (ChEFT) for π & N [e.g., Weinberg, Epelbaum, Meissner, Schwenk,...]
- c) Lattice QCD for NN, NY, YY int. [e.g., HAL collaboration,...]

2) Use the microscopic forces in **many-body methods**

- a) variational [e.g., Pandharipande, Takano, Togashi, ...] (w. **soft nucleons** $p^2/m_N < m_\pi$)
- b) Quantum Monte-Carlo [e.g., Carlson, Gandolfi,...]
- c) Hartree-Fock + many-body perturbation [e.g., Schwenk, Drischler, ...]

Advantage: can check how **systematic uncertainties** propagate;

→ the methods can predict **the domains of applicability**.

Domain of nuclear methods: rough estimate

$$\varepsilon(n_B) \quad \begin{array}{l} \text{2-body int.} \sim n_B^2 \text{ (contact)} \ \& \ n_B^{4/3} \text{ (long-range)} \\ \text{3-body int.} \sim n_B^3 \text{ (contact)} \ \& \ n_B^{5/3} \text{ (long-range)} \end{array}$$

Interactions **dominate** over kin. $E(\text{NR}) \sim n_B^{5/3}/m_N$

e.g.1) Akmal-Pandharipande-Ravenhall EoS (**APR**)

PNM n_B	2-body int.		3-body int.	
	$\langle v_{ij}^\pi \rangle$	$\langle v_{ij}^R \rangle$	$\langle V_{ijk}^{2\pi} \rangle$	$\langle V_{ijk}^R \rangle$
n_0	-4.1	-29.9	1.2	4.5
2 n_0	-25.1	-36.4	-17.4	30.6
3 n_0	-35.7	-44.7	-34.1	78.0
4 n_0	-52.2	-41.1	-76.9	160.3

↓ grow rapidly!

e.g.2) parameterized **pure neutron** matter EoS

[Gandolfi+ '09]

fit to Quantum Monte-Carlo

~kin. + 2-body

~3-body

$$\varepsilon = n_0 \left[(12 \pm 1 \text{ MeV}) \left(\frac{n_B}{n_0} \right)^{1.45 \pm 0.05} + (4 \pm 2 \text{ MeV}) \left(\frac{n_B}{n_0} \right)^{3.3 \pm 0.3} \right]$$

$n_B \sim 2n_0$ is the upper bound

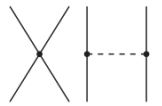
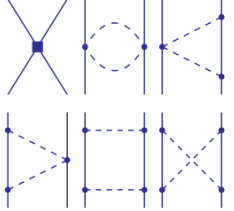
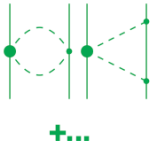
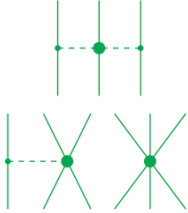
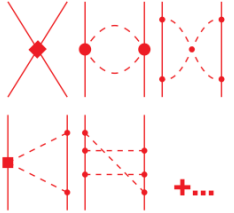

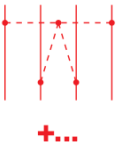
Chiral EFT [Weinberg 1990-]

chiral symmetry

- 1) **pions must** appear together with **derivatives** or **explicit sym. breaking terms**
 - 2) **Short distance** nucleon **contact** terms
 - 3) other heavy d.o.fs : “**integrated out**”
- **systematic organization** of terms in Lagrangian (!)

advantages especially in organizing **many-body forces**

expansions of Q/Λ_χ , m_π/Λ_χ

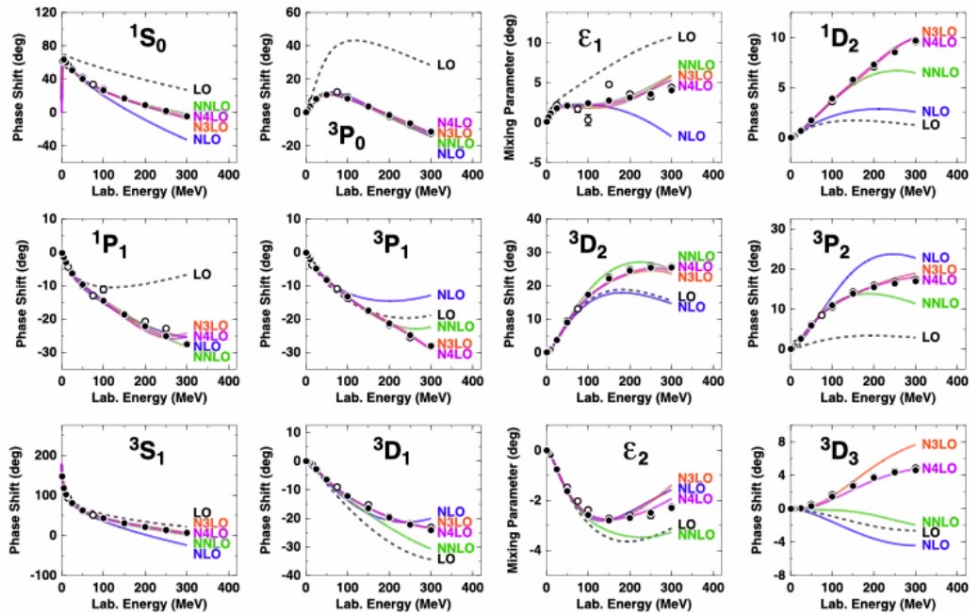
	2N Force	3N Force	4N Force
LO $(Q/\Lambda_\chi)^0$			
NLO $(Q/\Lambda_\chi)^2$			
NNLO $(Q/\Lambda_\chi)^3$			
N ³ LO $(Q/\Lambda_\chi)^4$			

Low energy constants

2N forces (well determined)

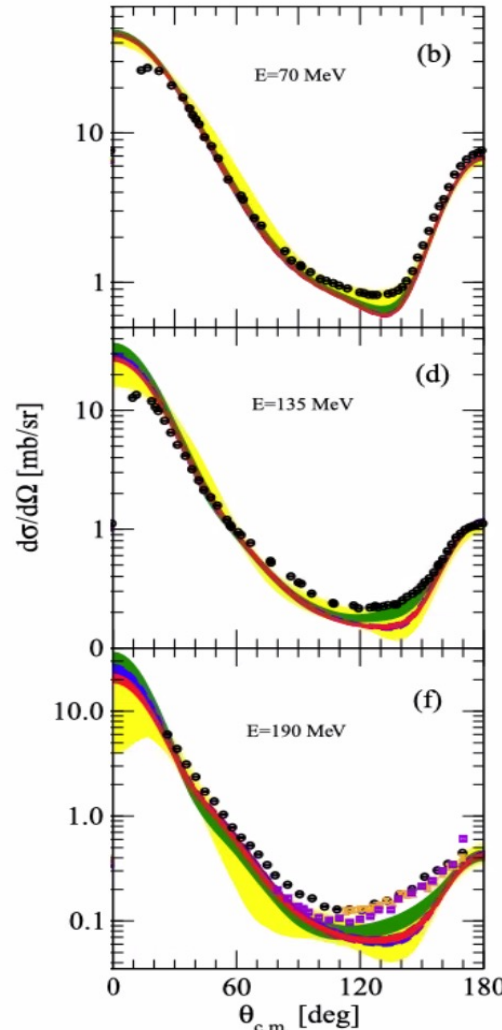
3N forces (progressing)

Neutron-proton scattering phase shifts



excellent fits (~6000 data)

~40 LECs (N⁴LO, 2NF)

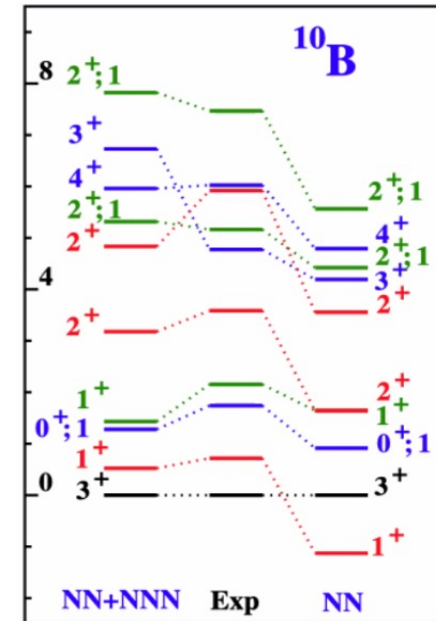


d(np) – p scatt. [Sekiguchi+ '17]

2NF to N⁴LO
miss some contributions

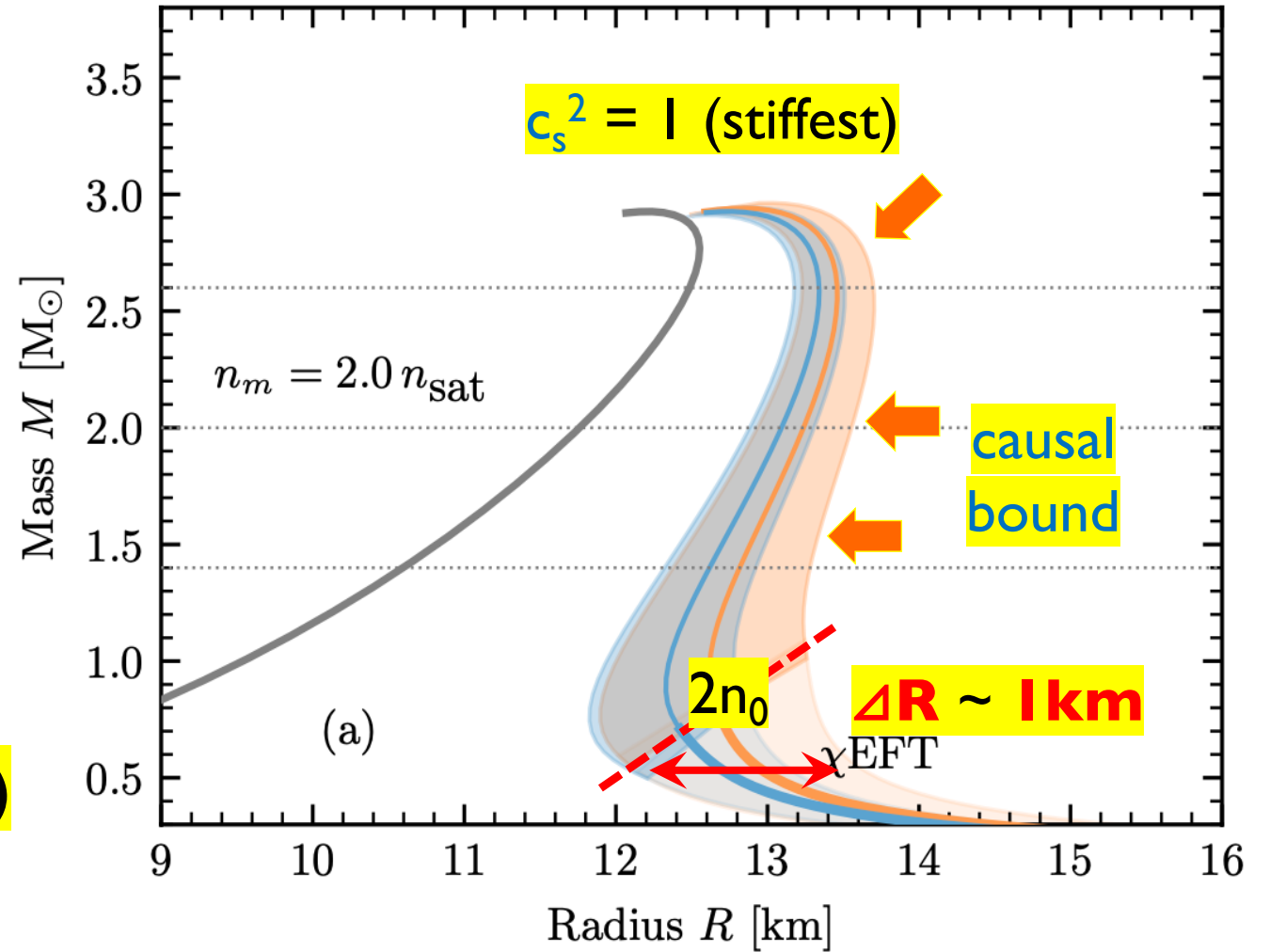
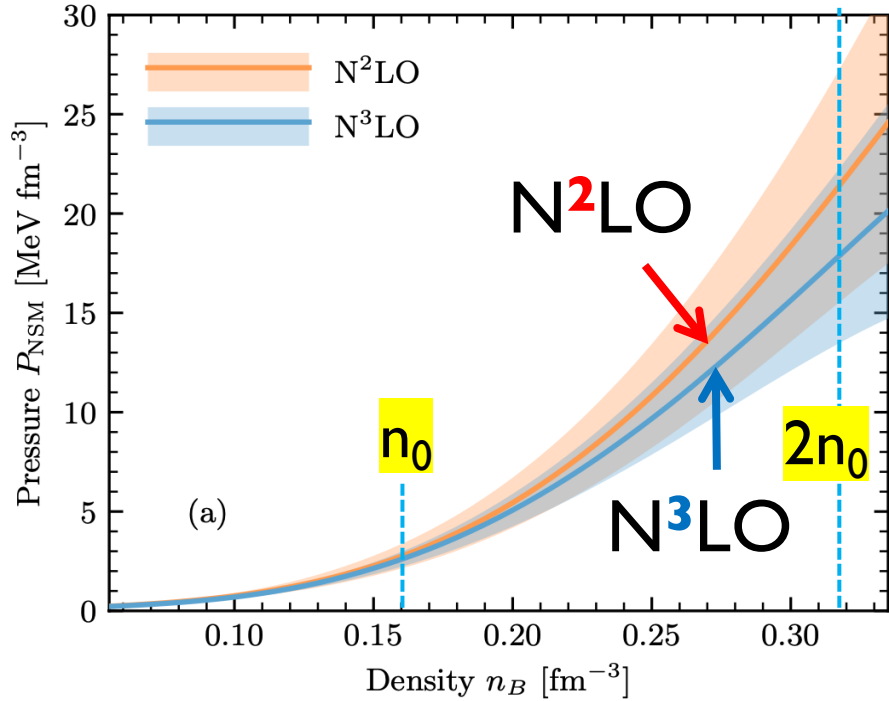
spectra of light nuclei

significant improvement with 3NF



EOS with N^3 LO ChEFT band

[Drischler+ '21]



EOS: $P = P_{\chi|2n_0} + c_s^2 (\epsilon - \epsilon_{\chi|2n_0})$

if we trust ChEFT to $2n_0$, the causality $\rightarrow R_{1.4} < \sim 12.9 \text{ km}$

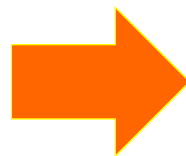
NICER + XMM + GW + nuclear physics (+ “ $c_s^2 < 1$ ”)

(see Lect.2)

NICER + XMM Newton

$$R_{2.08} = 13.7^{+2.6}_{-1.5} \text{ km}$$

Miller+ '21



+ GW + nuclear ($0.5n_0$)

[Miller+ '21]

$$R_{2.08} = 12.35^{+0.75}_{-0.75} \text{ km}$$

$$R_{1.4} = 12.45^{+0.65}_{-0.65} \text{ km}$$

reduction of errors (!)

NICER + XMM Newton

$$R_{2.08} = 12.39^{+1.30}_{-0.98} \text{ km}$$

Riley+ '21



+ GW + ChEFT ($0.5-1.1n_0$)

[Raajimakers+ '21]

$$R_{1.4} = 12.33^{+0.76}_{-0.81} \text{ km}$$

reduction of errors (!)

$$R_{2.08} \sim R_{1.4} (!)$$

Summary of Lecture I

- glancing at NS properties; M, R, Y_p, B, T, \dots
- EOS and M-R relations
- EOS at low density; NICER $R_{1.4}$ and ChEFT

$R_{2.08} \sim R_{1.4} (!)$ hints for soft-to-stiff EOS

Lect. 2 : Gravitational waves & NS-NS mergers

Lect. 3 : From hadrons to quarks in NS