



Entanglement entropy and lattice gauge theory



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Bern, Switzerland

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Introduction

What is entanglement?

- Quantum physical implementation of conservation laws

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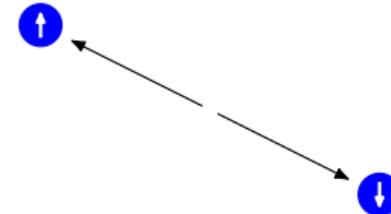
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 - Decay of spin-0 particle: $s = 0$



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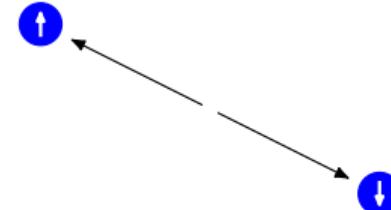


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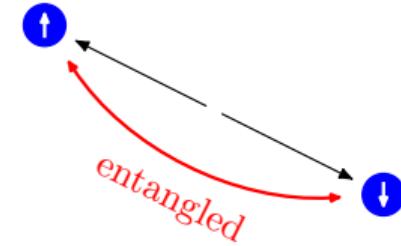


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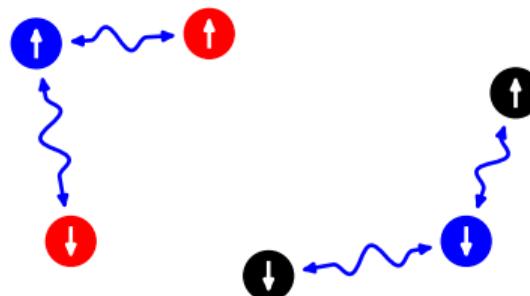
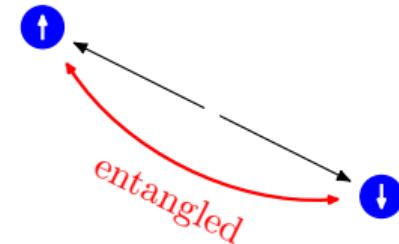


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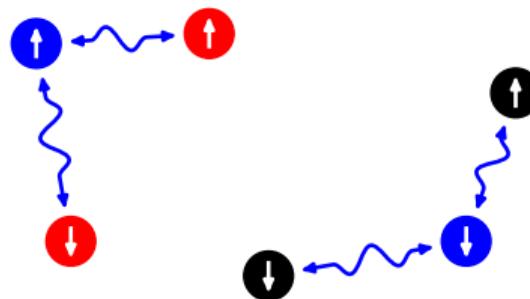
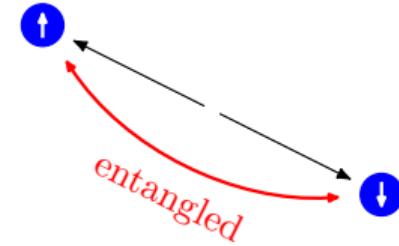


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 - correlations



Introduction

What is entanglement entropy?

■ Preliminaries:

Hilbert space: \mathcal{H} , state vector: $|\psi\rangle \in \mathcal{H}$

Density matrix:

$$\rho = \sum_i p_i |\psi_i\rangle\langle\psi_i| \quad , \quad |\psi_i\rangle \in \mathcal{H} \quad \forall i \quad , \quad \sum_i p_i = 1$$

$$\text{tr}(\rho) = 1$$

pure state: $\rho = |\psi\rangle\langle\psi|$

$$\rightarrow \rho^2 = \rho \text{ (projector)} \rightarrow \text{tr}(\rho^2) = 1$$

mixed state: $\rho = \sum_i p_i |\psi_i\rangle\langle\psi_i|$

$$\rightarrow \rho^2 \neq \rho \text{ (not projector)} \rightarrow \text{tr}(\rho^2) < 1$$

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What is entanglement entropy?

■ Bipartite quantum system: $\mathcal{H}_{AB} = \mathcal{H}_A \otimes \mathcal{H}_B$

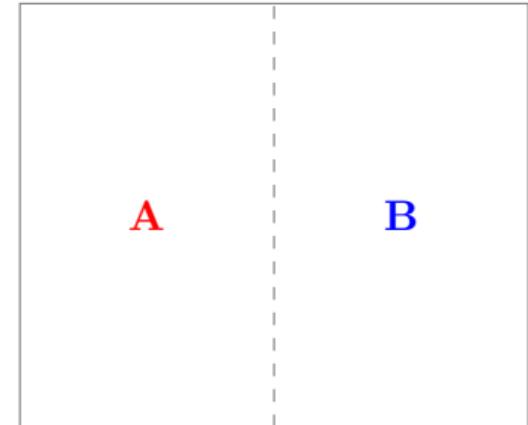
pick pure state: $|\psi\rangle_{AB} \in \mathcal{H}_{AB}$

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$$\rightarrow |\psi\rangle_{AB} = \sum_{mn} a_{mn} |m\rangle_A \otimes |n\rangle_B , \quad \sum_{mn} |a_{mn}|^2 = 1$$

$$\rightarrow \rho_{AB} = |\psi\rangle_{AB}\langle\psi| = \sum_{mnkl} a_{mn} a_{kl}^* |m\rangle_A \langle k| \otimes |n\rangle_B \langle l|$$

(notation: $|\psi\rangle_C\langle\psi| = |\psi\rangle_C \otimes {}_C\langle\psi|$)



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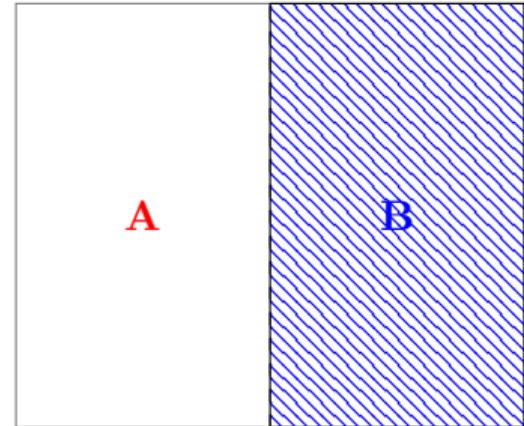
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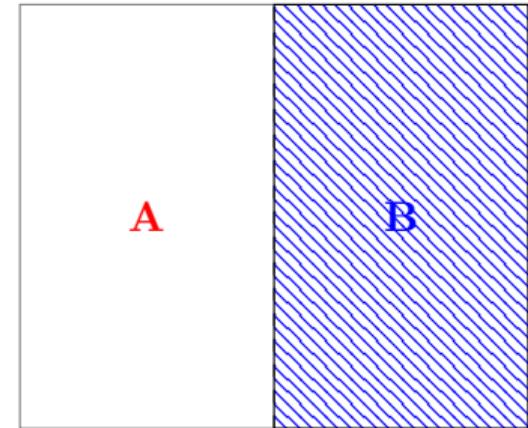
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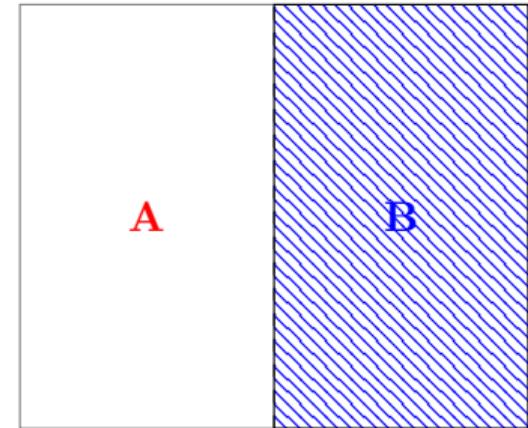
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$\rightarrow |\psi\rangle_{AB} = |\psi\rangle_A \otimes |\psi\rangle_B \implies \text{tr}(\rho_A^2) = 1 \implies$ no entanglement



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- Entanglement measures:

→ Purity: $\text{tr}(\rho_A^2)$

→ Rényi entropies: $H_s(A) = -\frac{1}{s-1} \log \text{tr}(\rho_A^s) , s = 2, 3, \dots$

→ Entanglement entropy: $S_{EE}(A) = -\lim_{s \rightarrow 1} \frac{\partial \log \text{tr}(\rho_A^s)}{\partial s} = \lim_{s \rightarrow 1} \frac{\partial((s-1)H_s(A))}{\partial s} = \lim_{s \rightarrow 1} H_s(A)$

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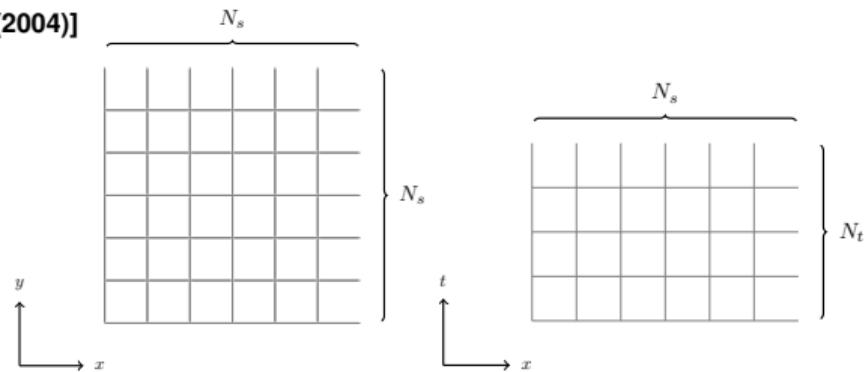
→ Entanglement entropy: $S_{EE}(A) = -\text{tr}(\rho_A \log(\rho_A))$ (Von Neumann entropy)

Entanglement entropy on the lattice

Entanglement entropy on the lattice [P. Calabrese, J. Cardy (2004)]

- SU(N) gauge theory on $N_s^{d-1} \times N_t$ lattice

Partition function: $Z(N_t, N_s) = \int \mathcal{D}[U] e^{-S_G[U]}$



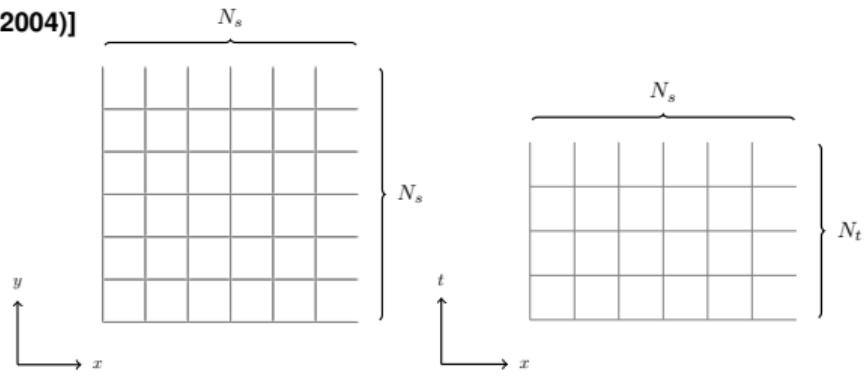
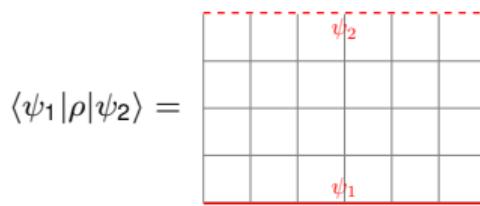
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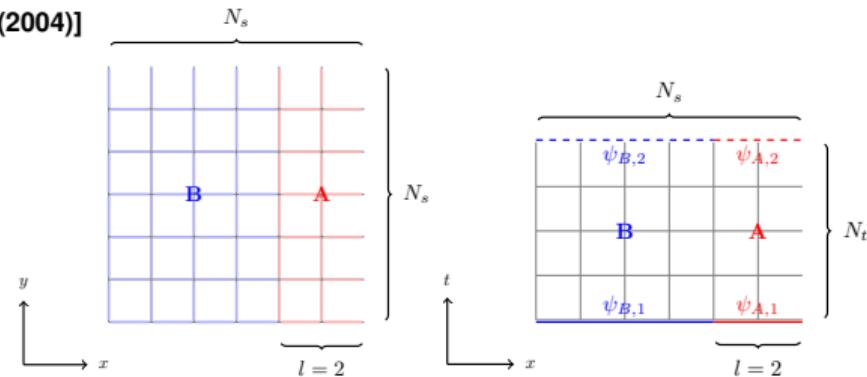
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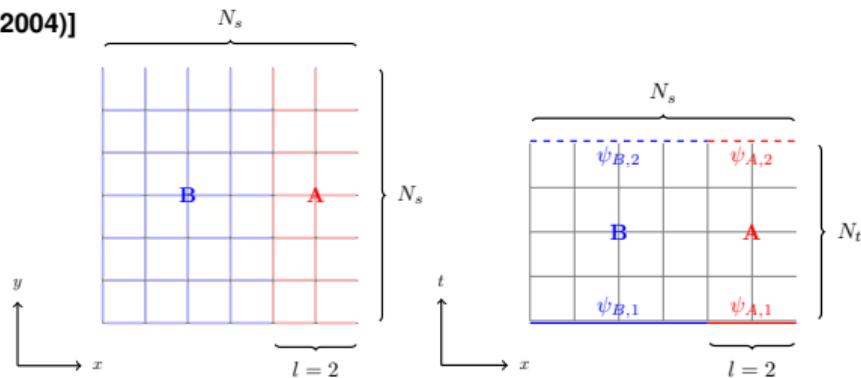
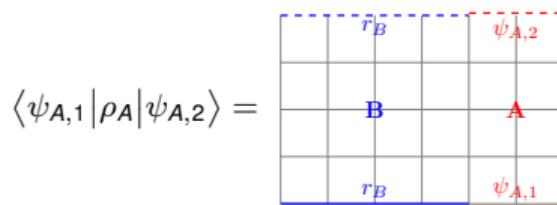
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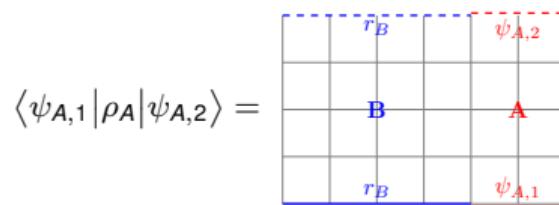
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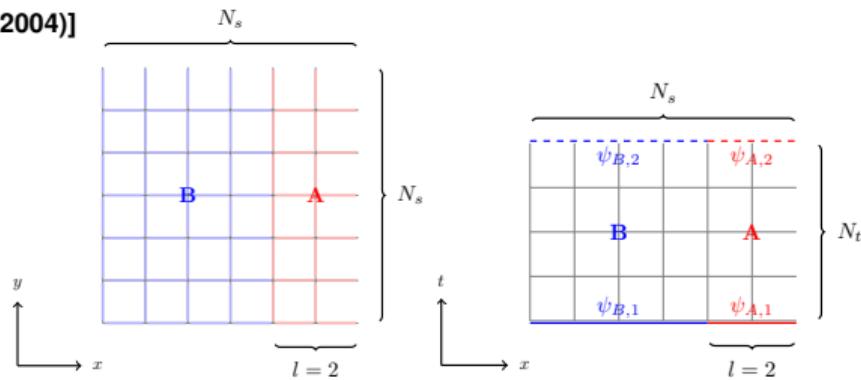
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→ Entanglement entropy:

$$S_{EE} = -\text{tr}_A(\rho_A \log \rho_A) \quad (\text{how ?})$$



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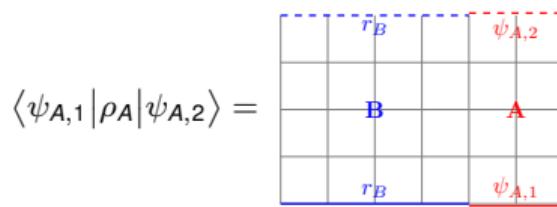
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$\langle \psi_{A,1} | \rho_A | \psi_{A,2} \rangle =$

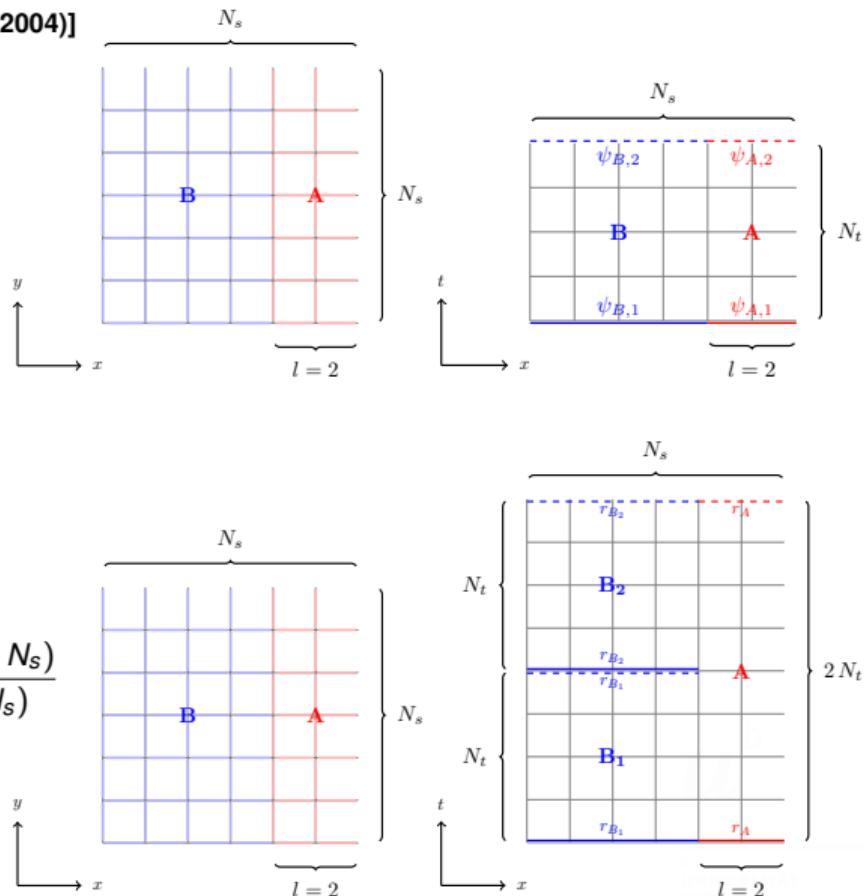
→ Replica method for s-th Rényi entropy:

$$H_s(I, N_t, N_s) = \frac{1}{1-s} \log \text{tr}(\rho_A^s) = \frac{1}{1-s} \log \frac{Z_c(I, s, N_t, N_s)}{Z^s(N_t, N_s)}$$

with "cut partition function" $Z_c(I, s, N_t, N_s)$

$$\rightarrow Z_c(I=0, s, N_t, N_s) = Z^s(N_t, N_s) \quad \forall s \in \mathbb{N}$$

$$\rightarrow Z_c(I=N_s, s, N_t, N_s) = Z(s N_t, N_s) \quad \forall s \in \mathbb{N}$$



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→ Entanglement entropy (EE):

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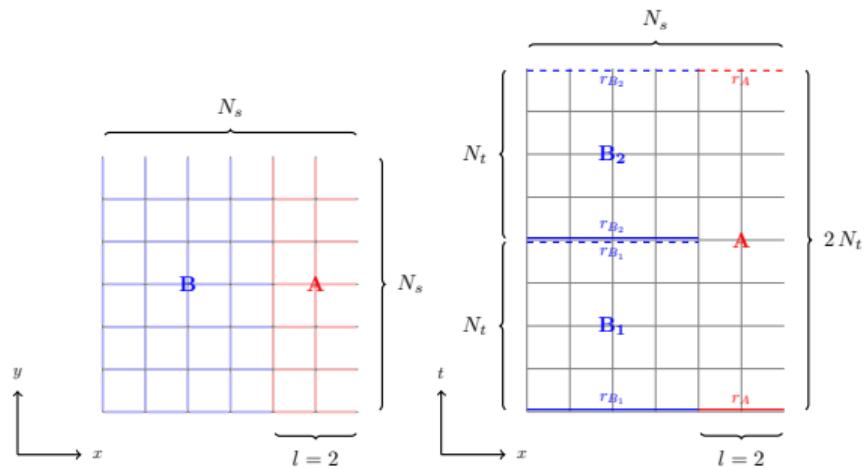
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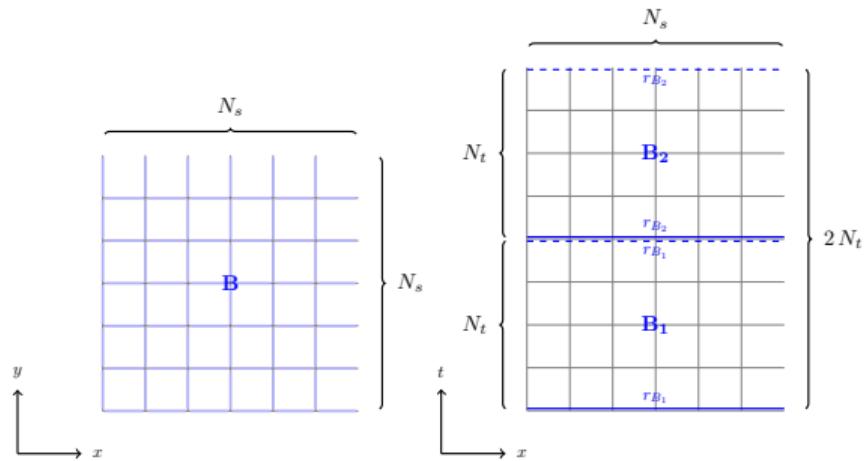
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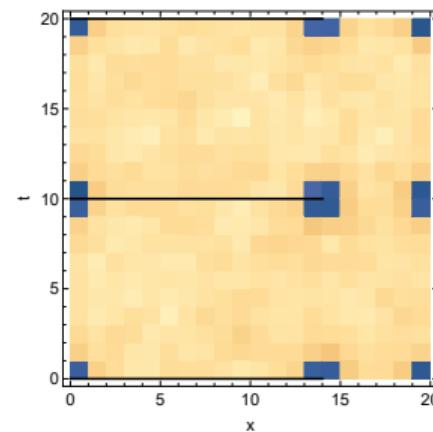
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Issue: UV-divergent piece

$$\frac{S_{EE}}{|\partial A|} = \frac{C_0}{\epsilon^2} - \frac{C}{l^q} + (\text{finite})$$



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→ Instead of EE, measure discrete derivative w.r.t. $I > 0$:

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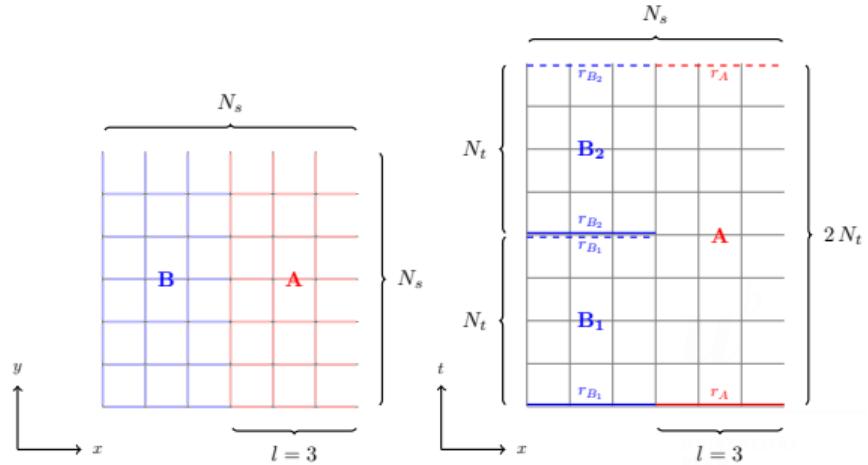
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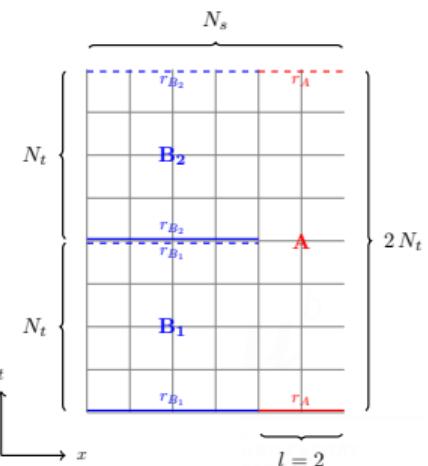
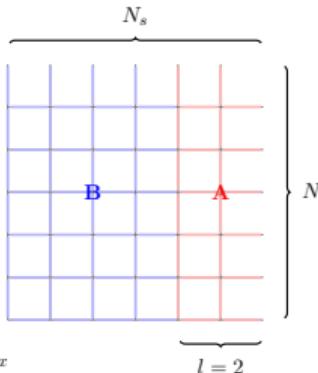
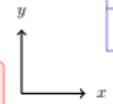
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- Approach from literature: [P. V. Buividovich, M. I. Polikarpov (2008)], [Y. Nakagawa et al. (2009)]

→ interpolating partition function:

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Entanglement entropy on the lattice

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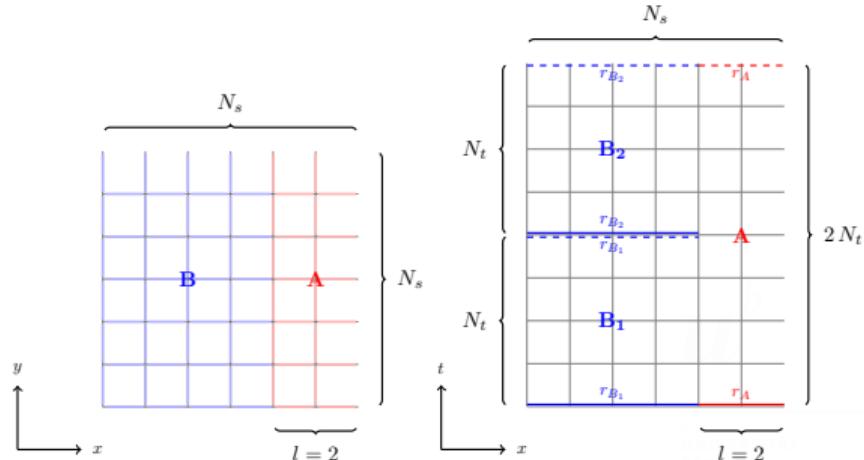
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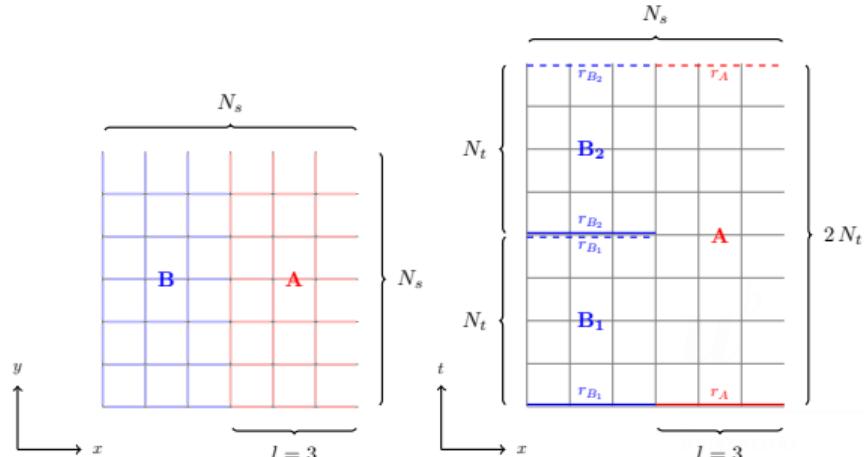
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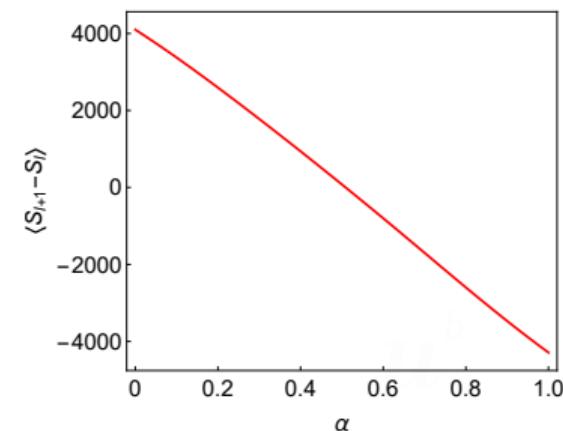
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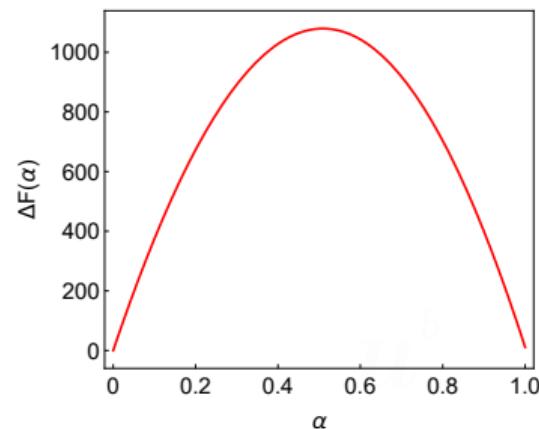
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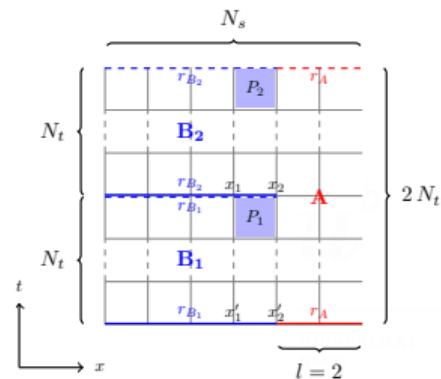
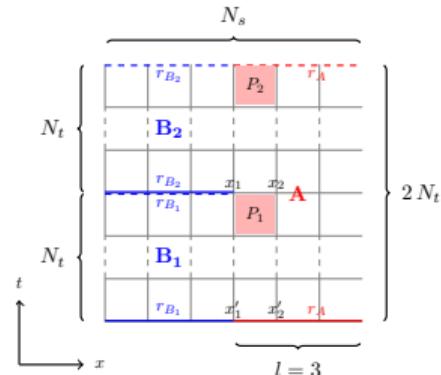
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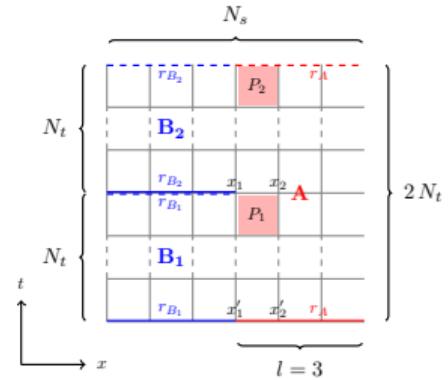
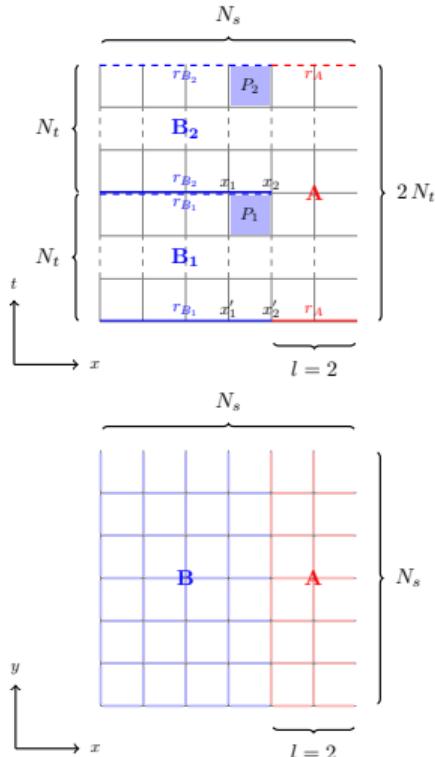
→ $Z_I^*(\alpha)$ imposes simultaneously BC_A and BC_B on plaquettes P_1, P_2 if $\alpha \neq 0, 1$.



Entangling surface deformation method

How can we avoid (huge) free energy barriers?

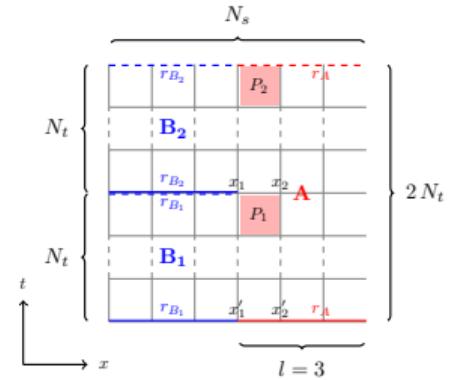
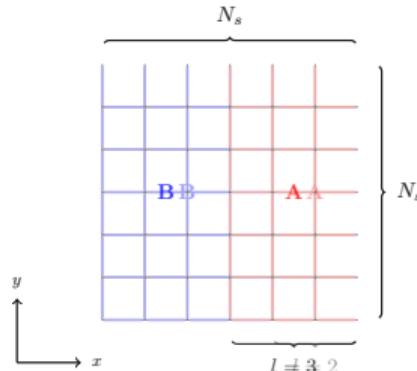
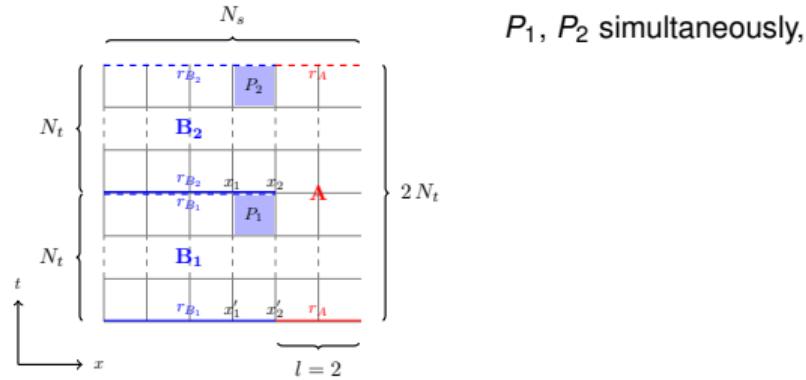
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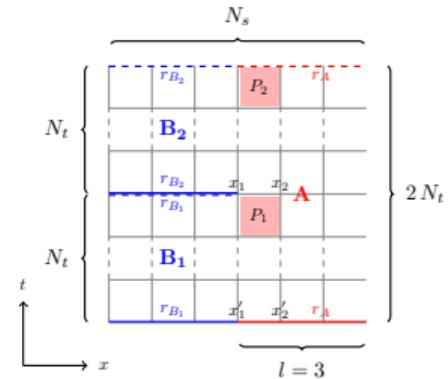
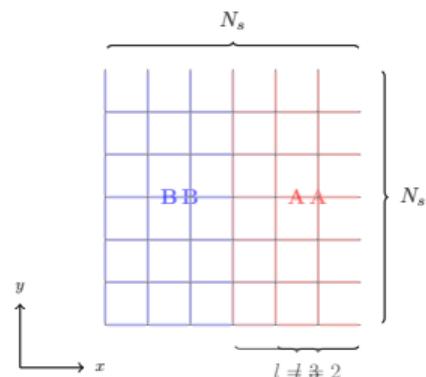
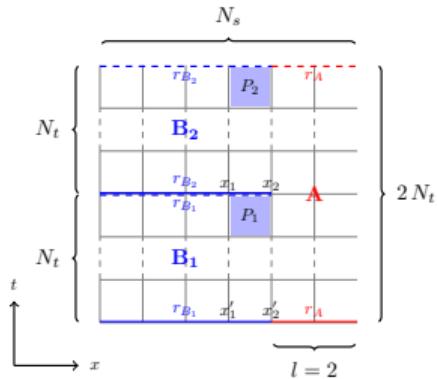
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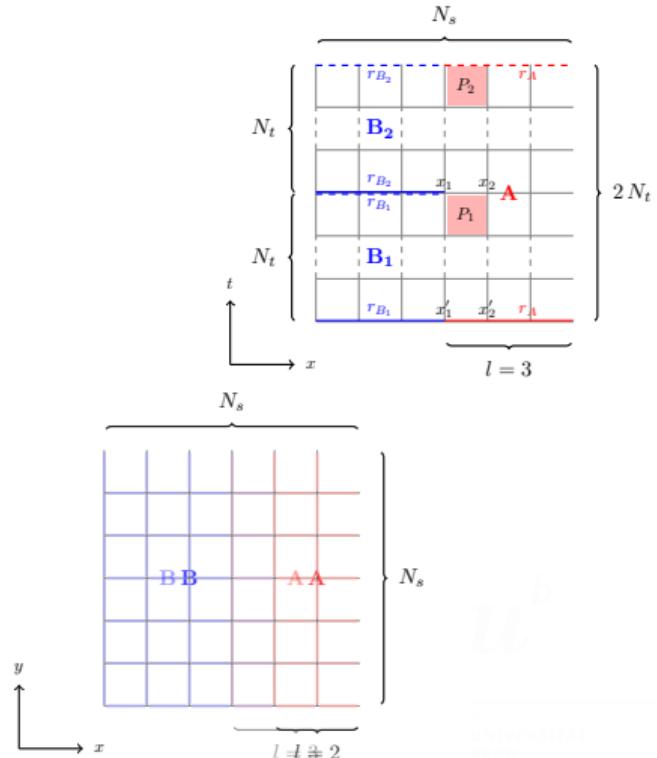
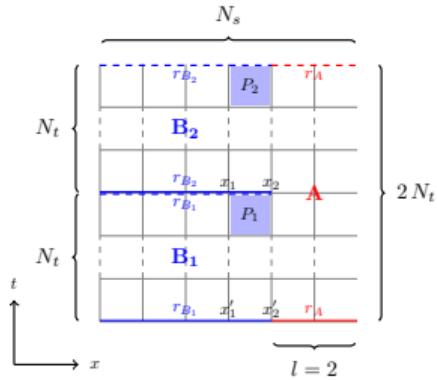
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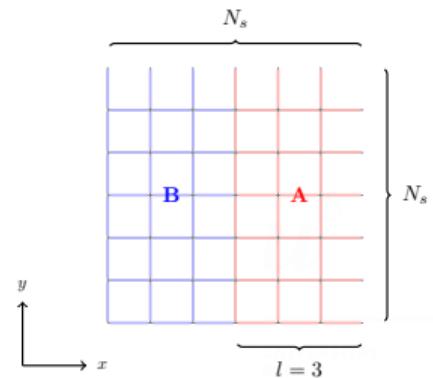
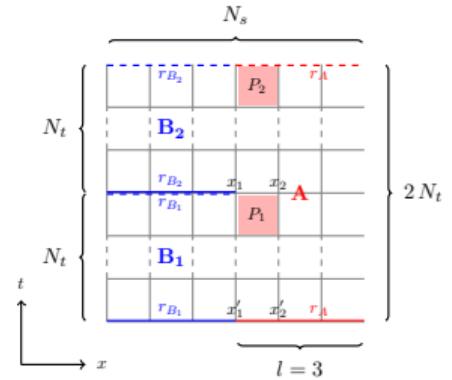
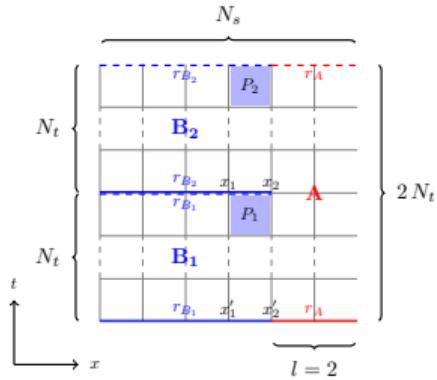
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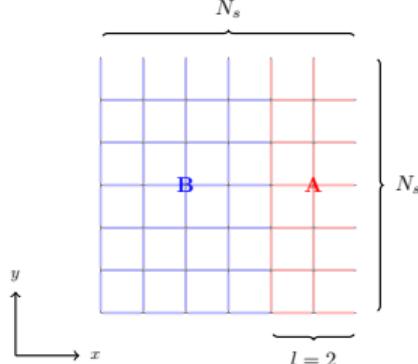
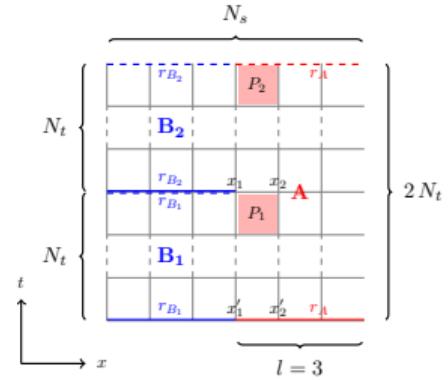
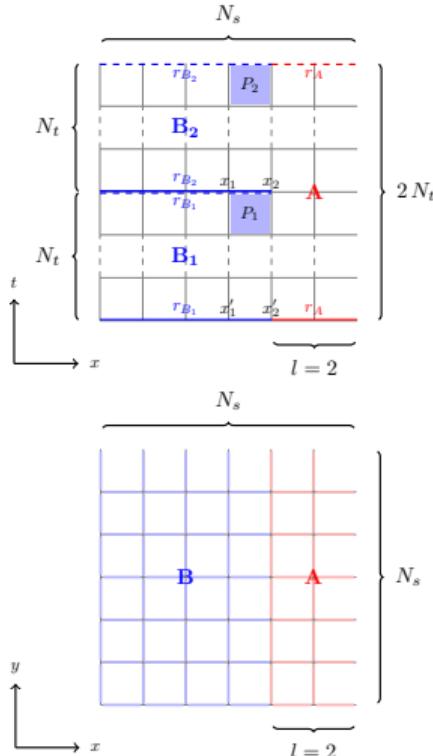
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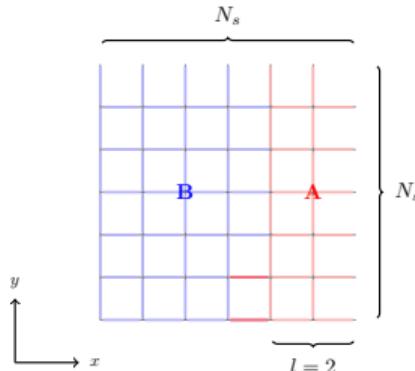
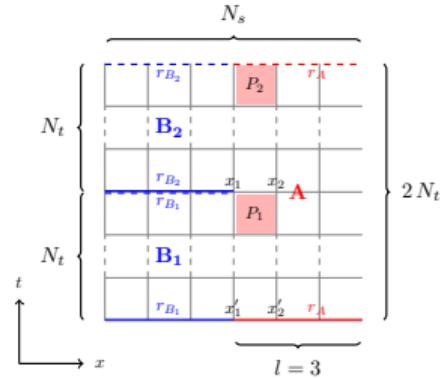
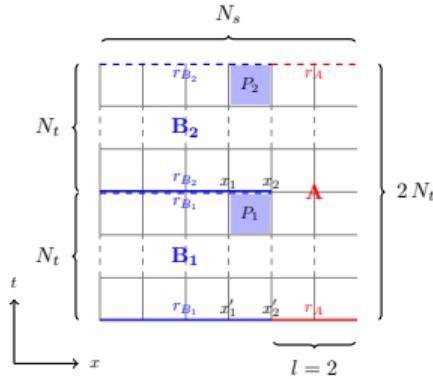
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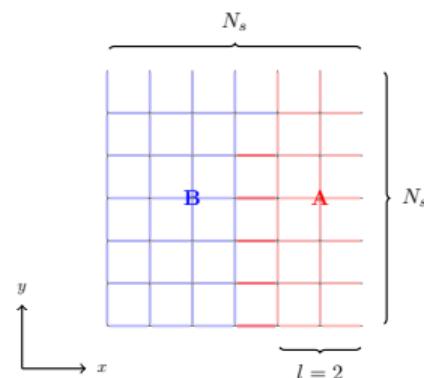
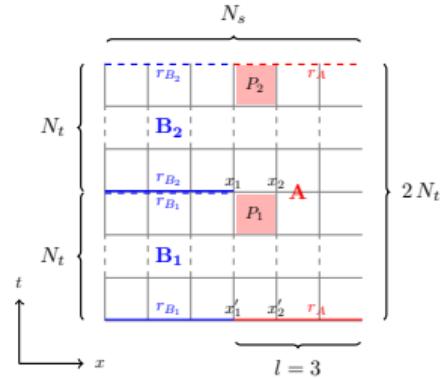
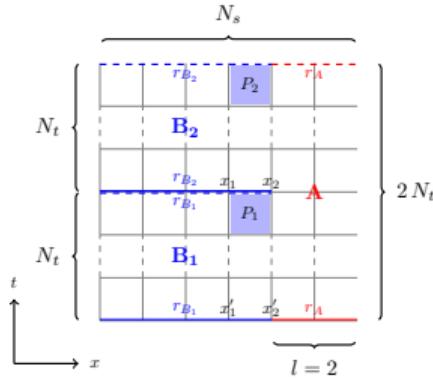
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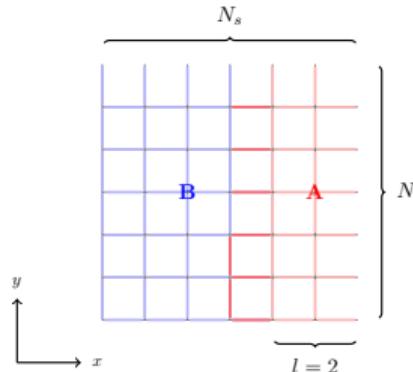
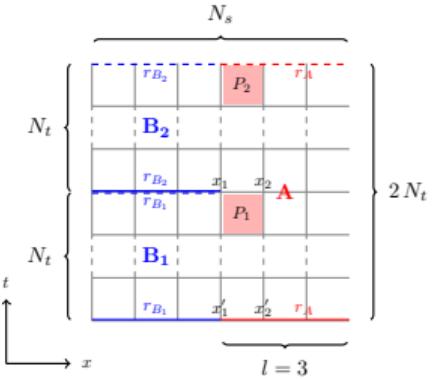
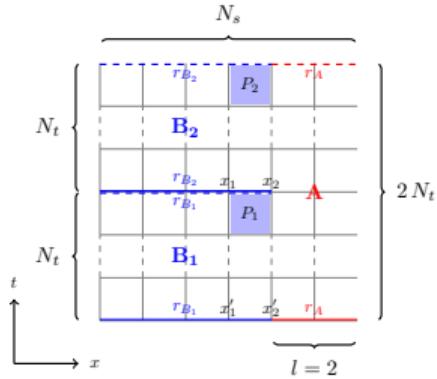
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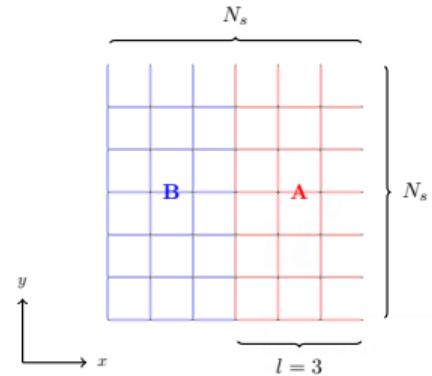
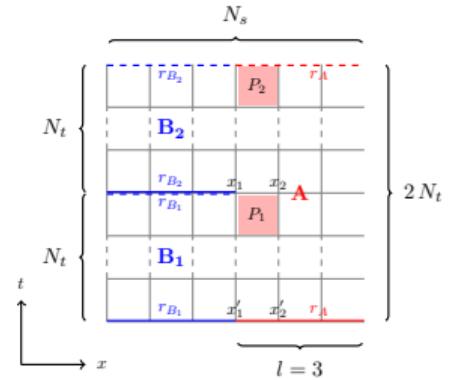
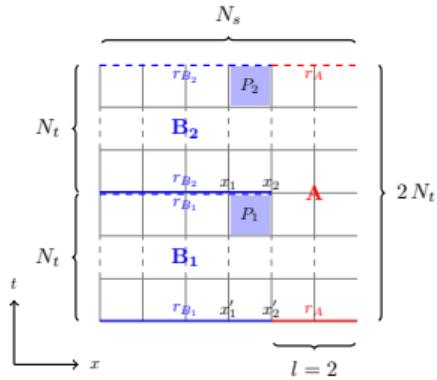
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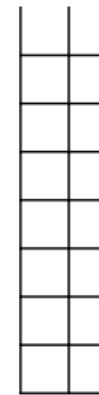
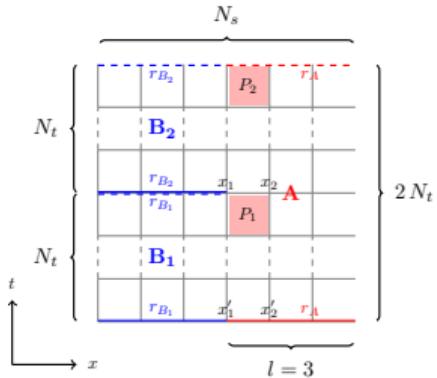
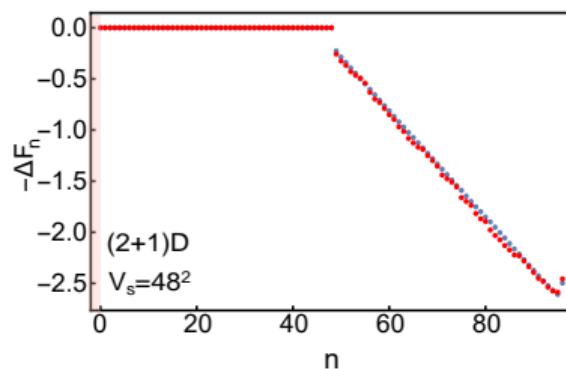
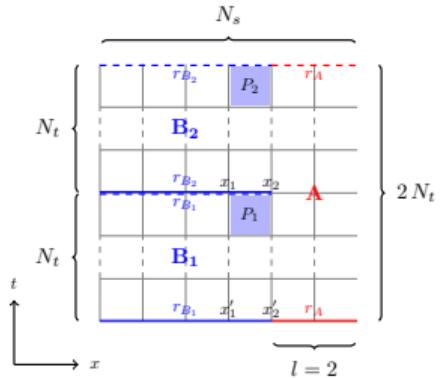
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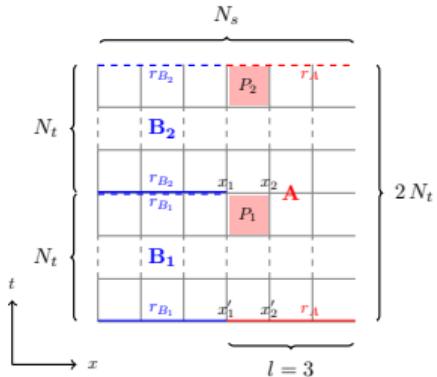
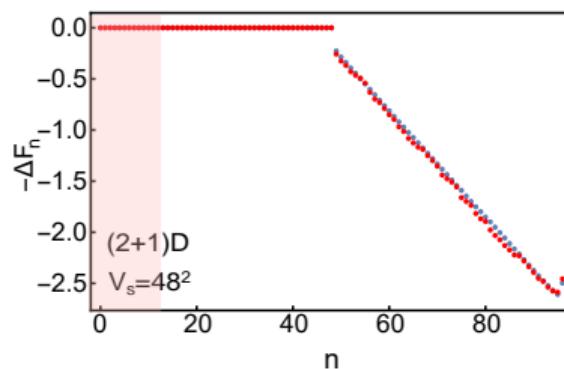
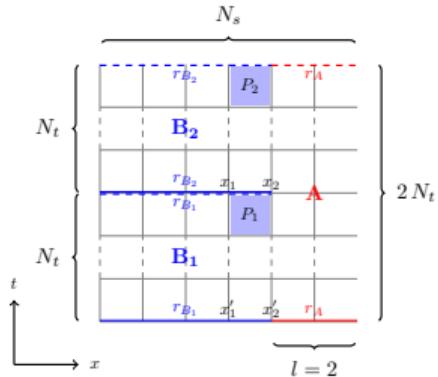
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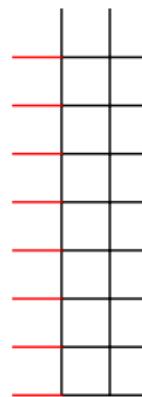
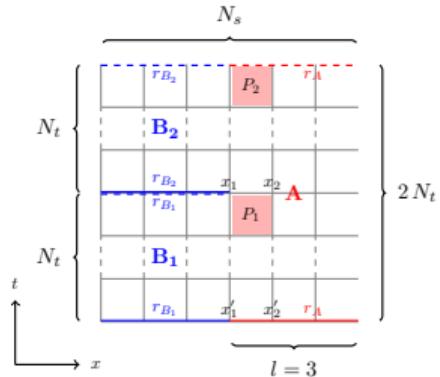
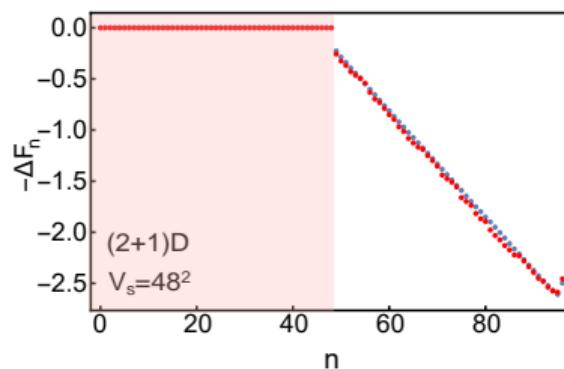
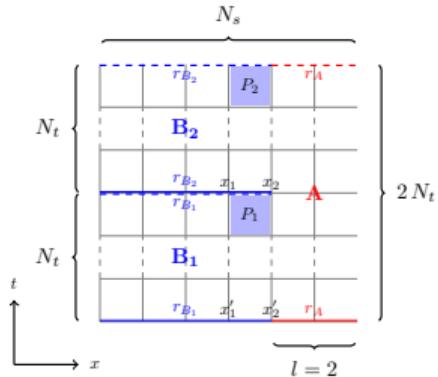
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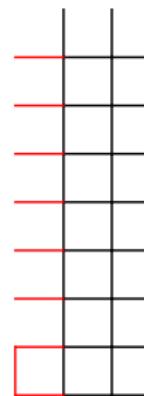
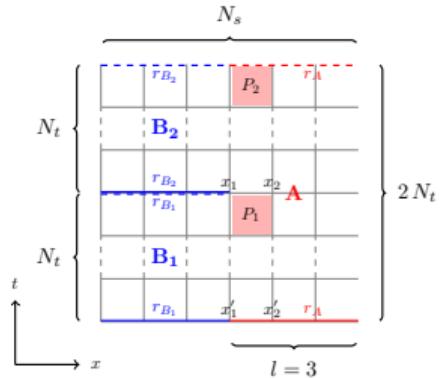
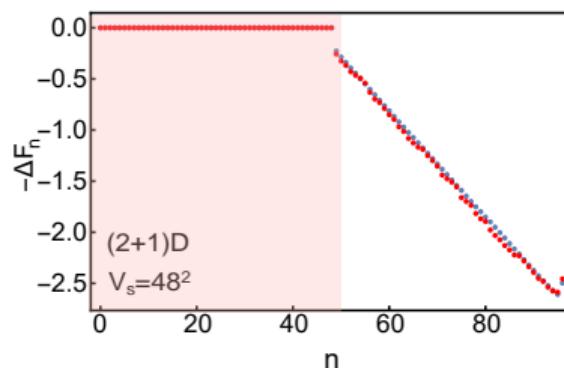
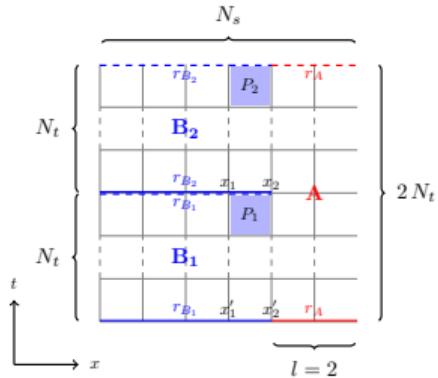
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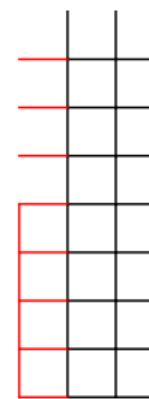
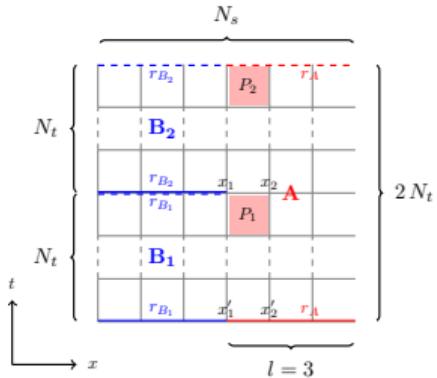
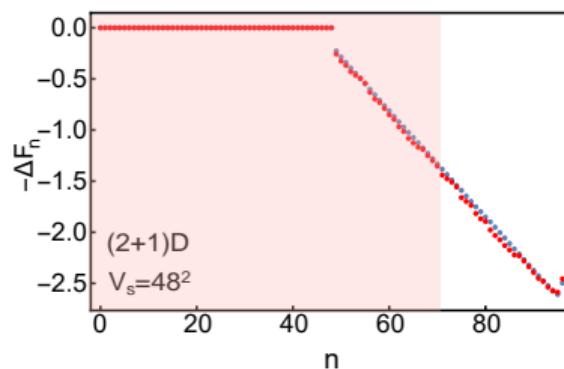
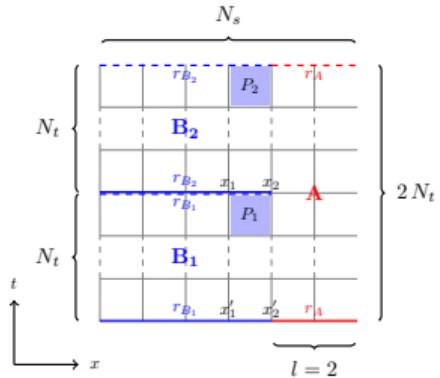
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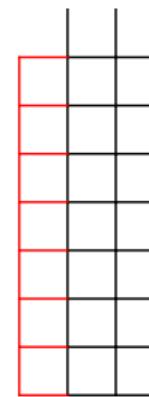
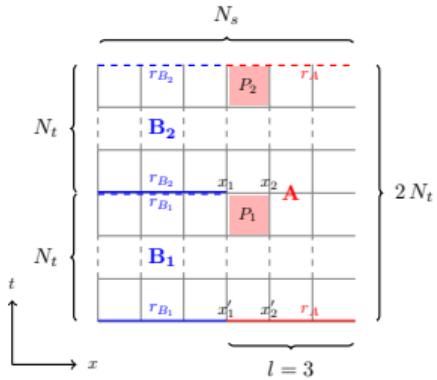
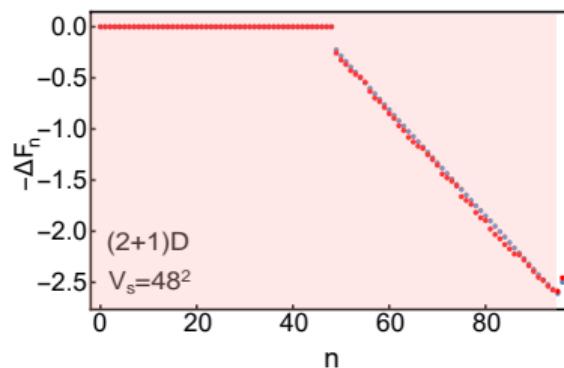
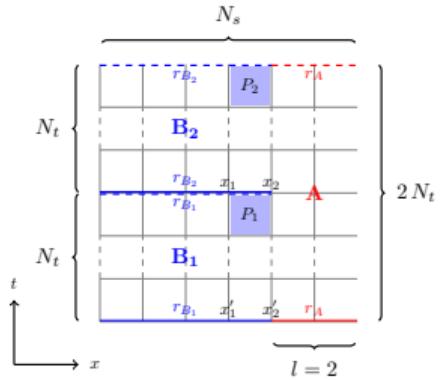
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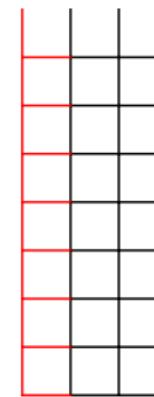
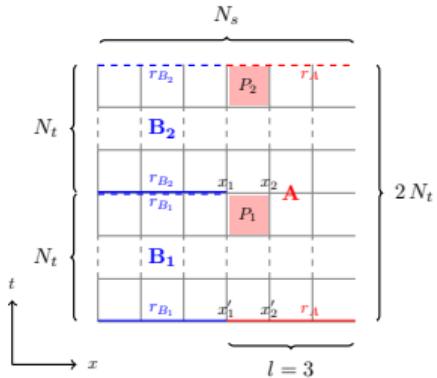
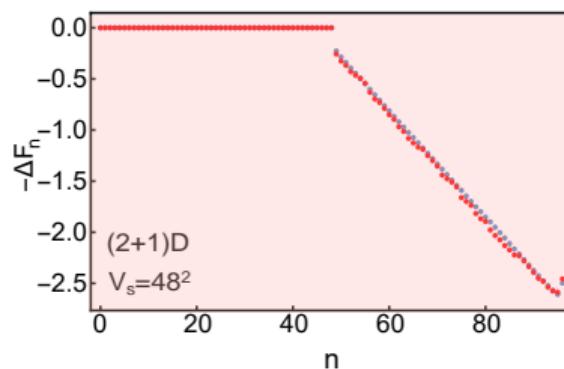
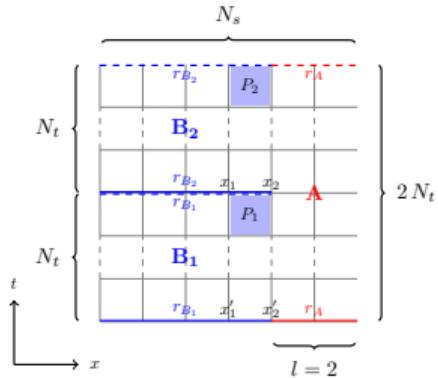
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Entangling surface deformation method

How can we avoid (huge) free energy barriers?

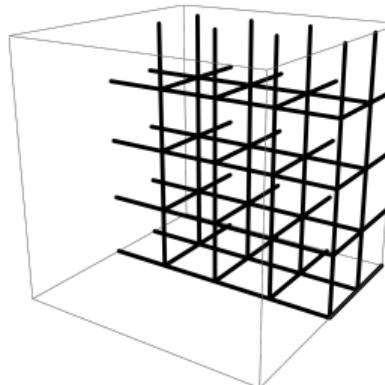
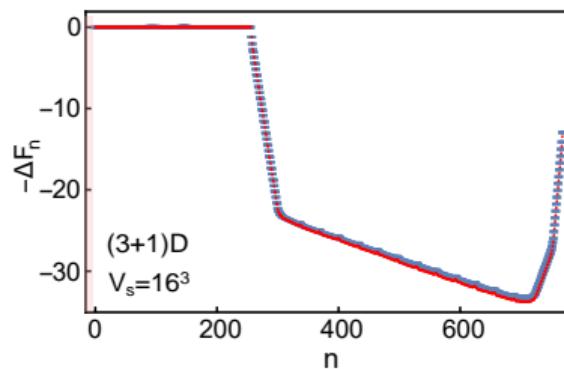
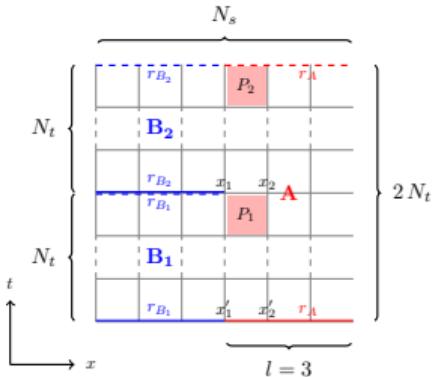
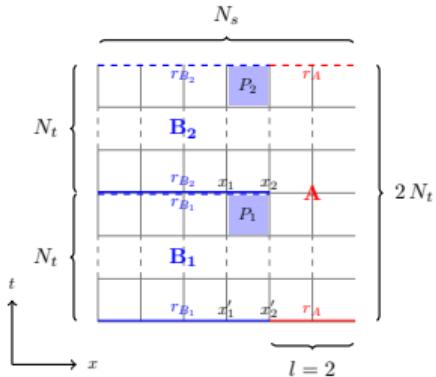
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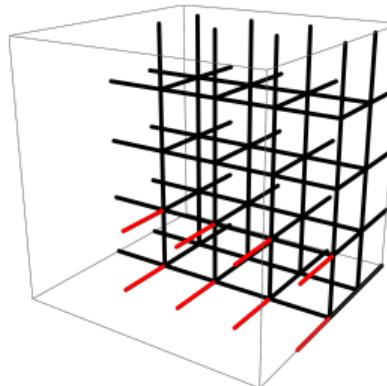
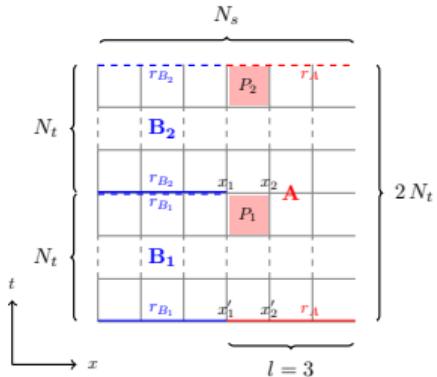
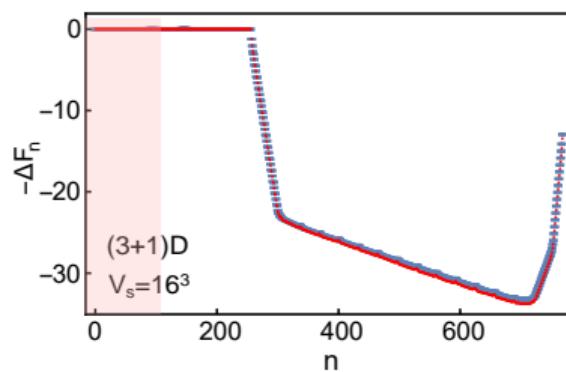
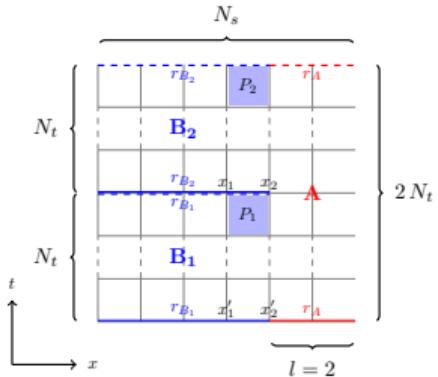
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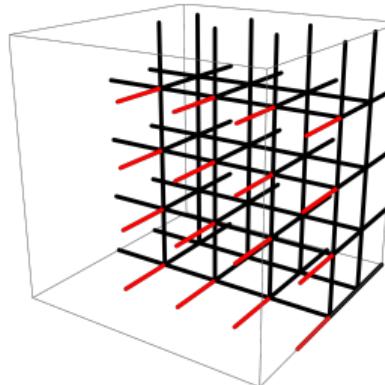
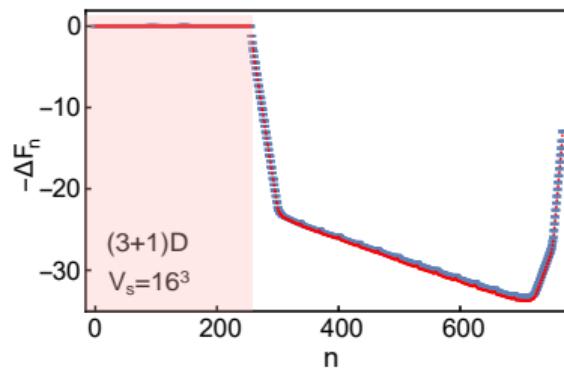
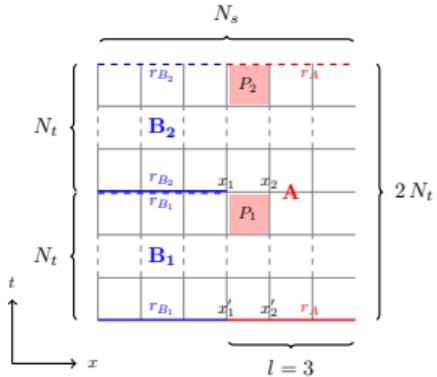
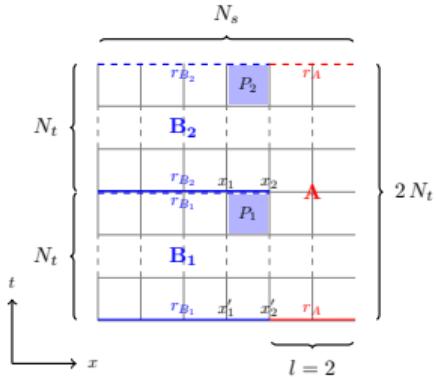
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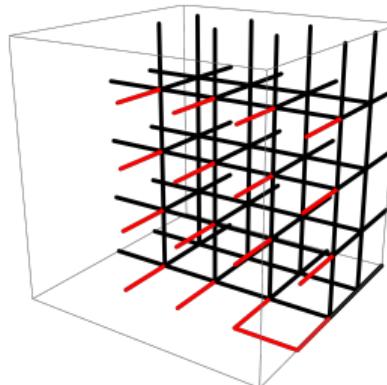
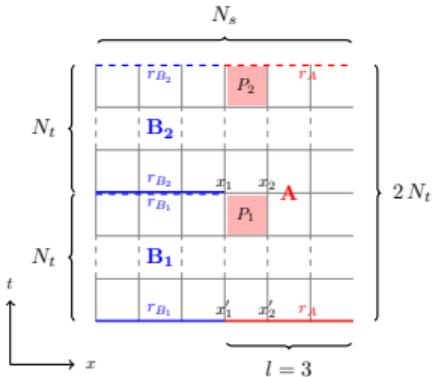
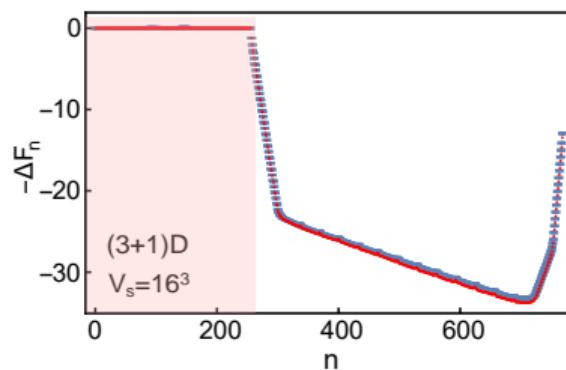
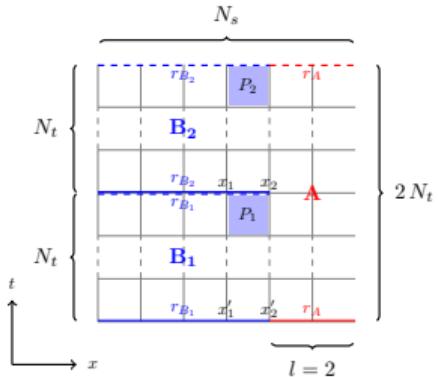
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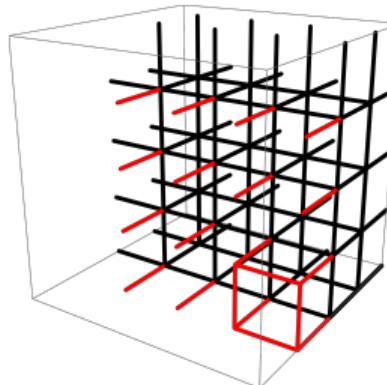
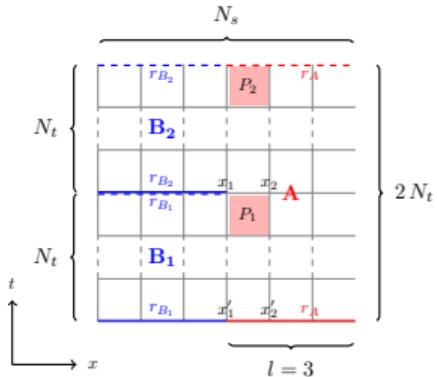
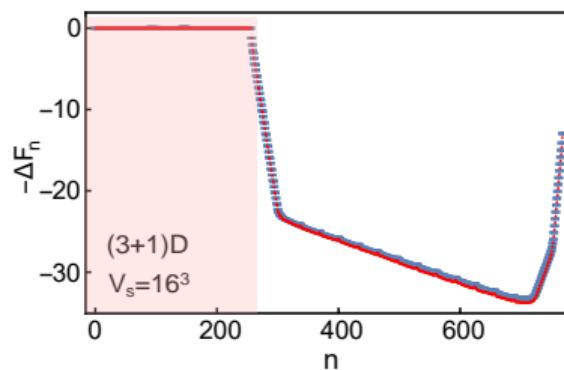
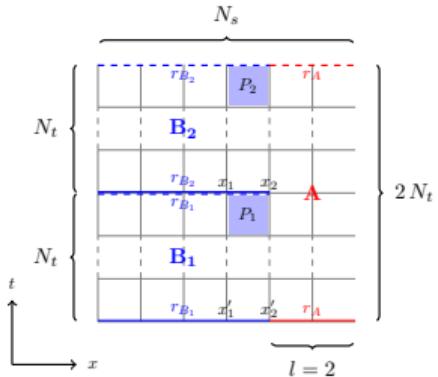
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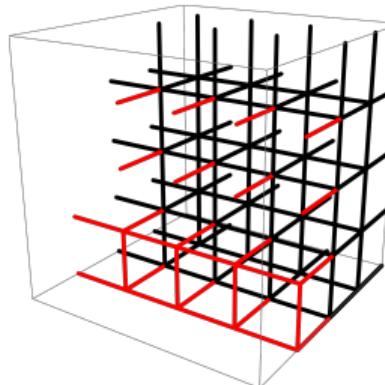
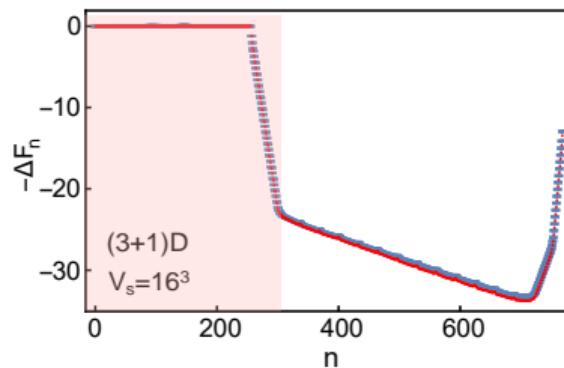
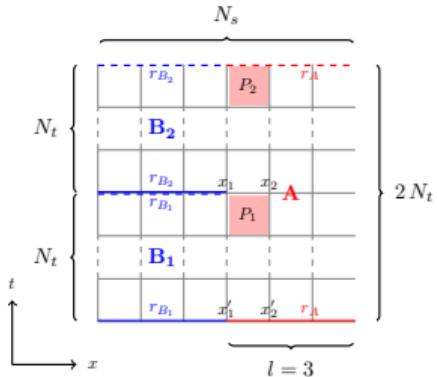
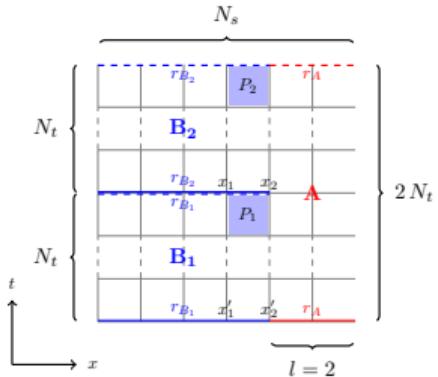
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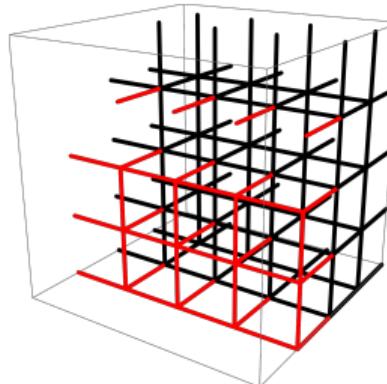
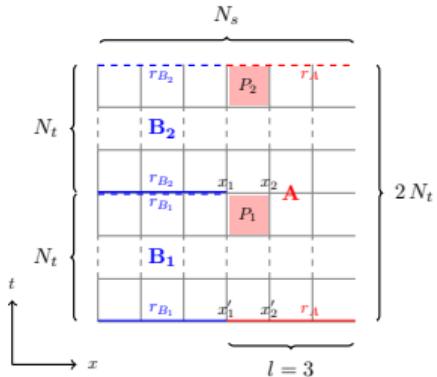
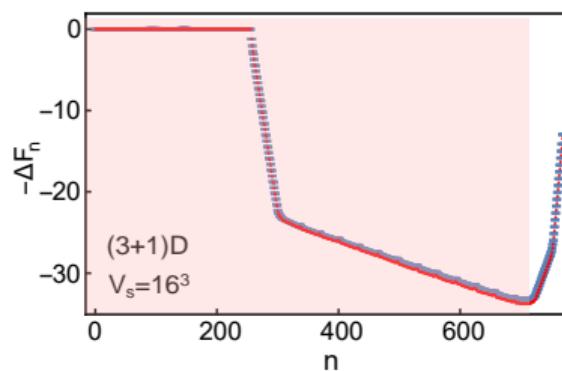
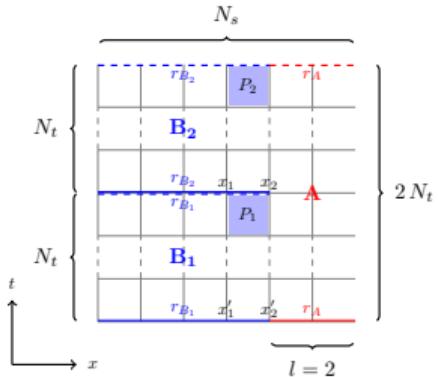
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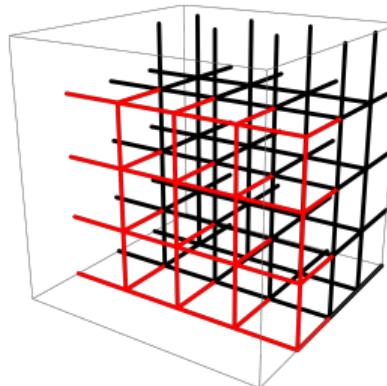
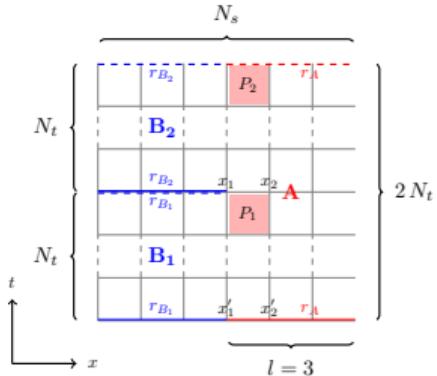
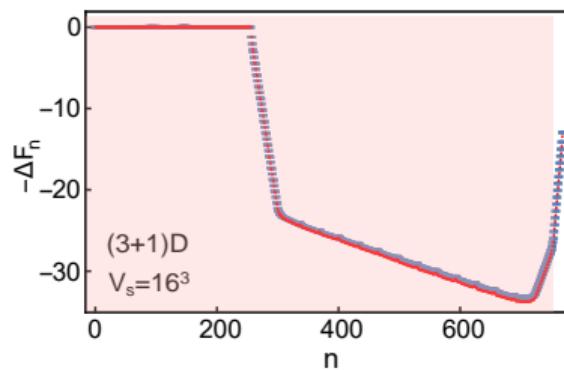
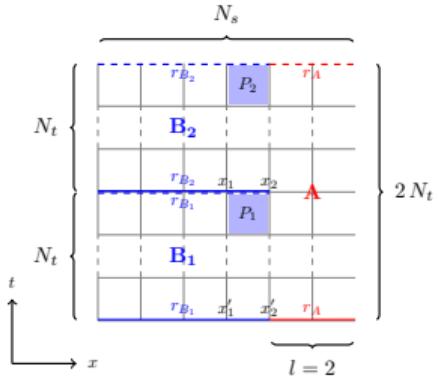
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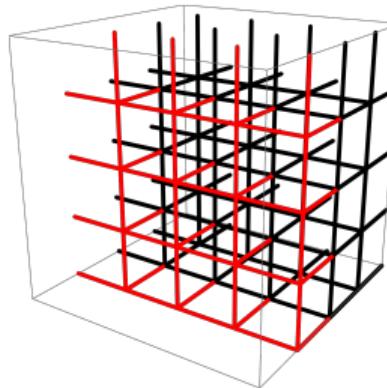
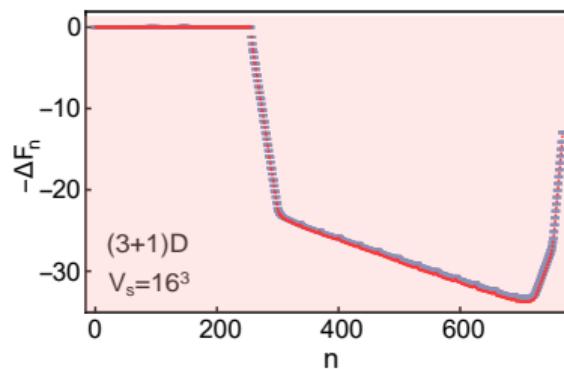
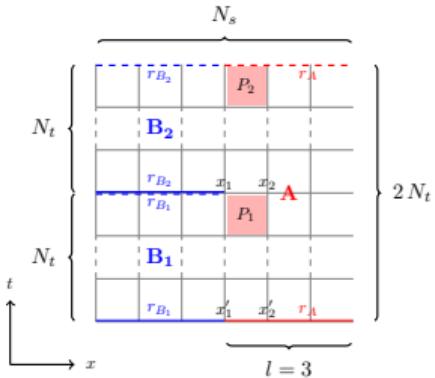
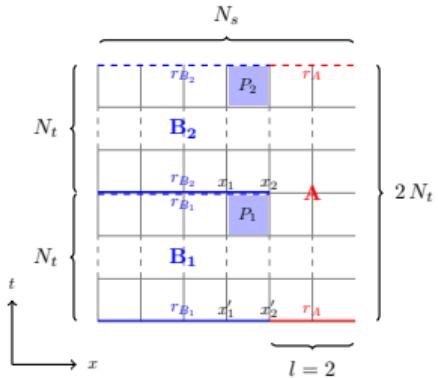
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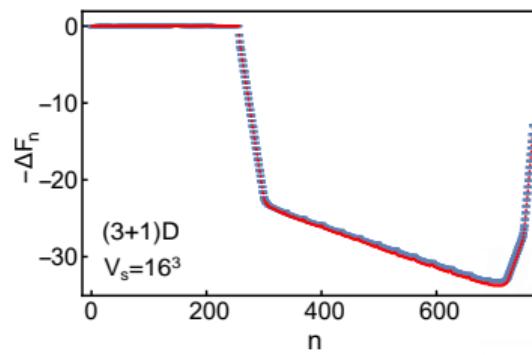
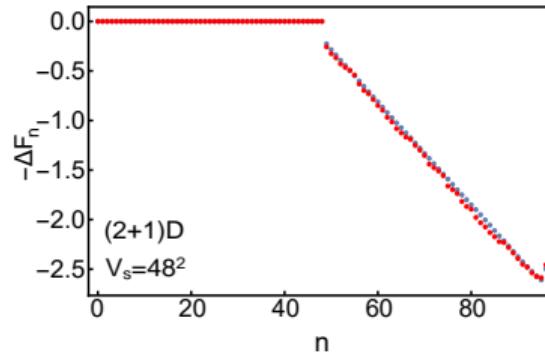
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Entangling surface deformation method

Free-energy plateau

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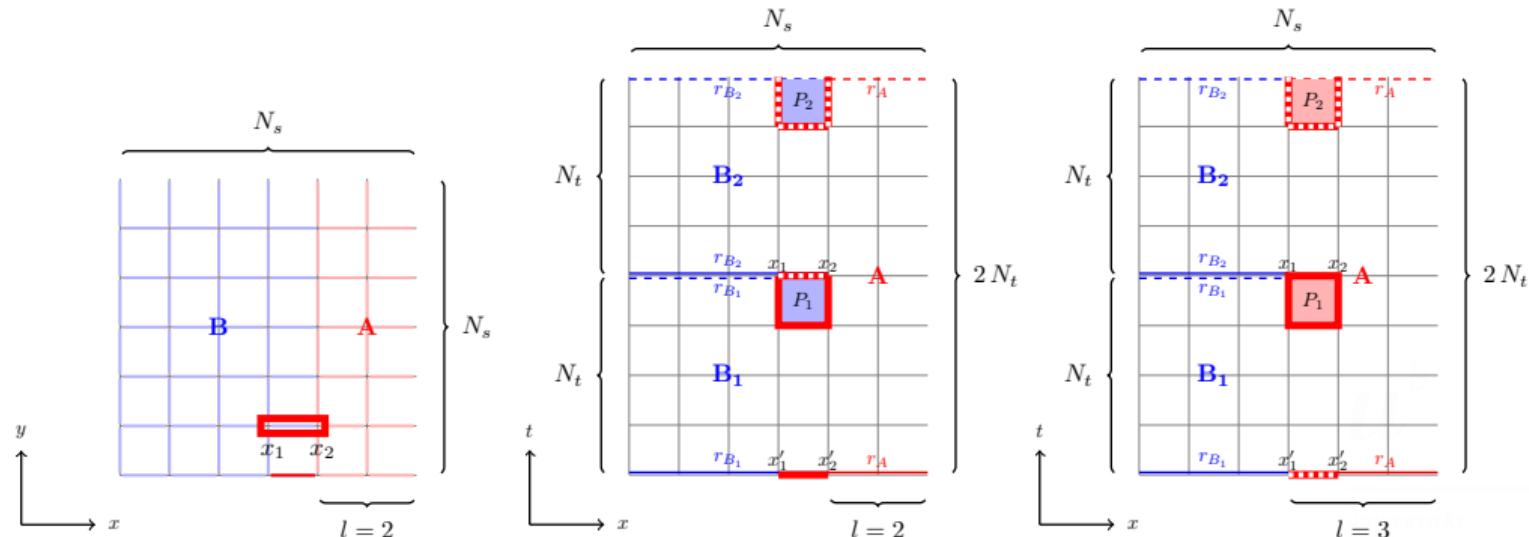


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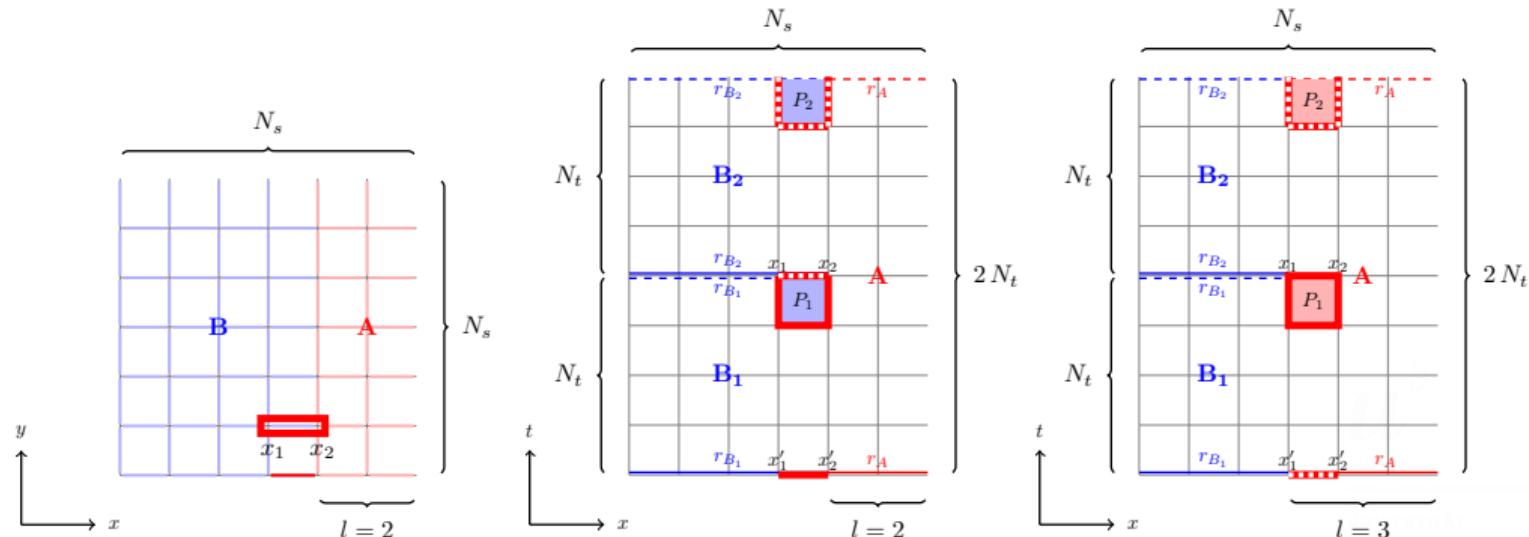
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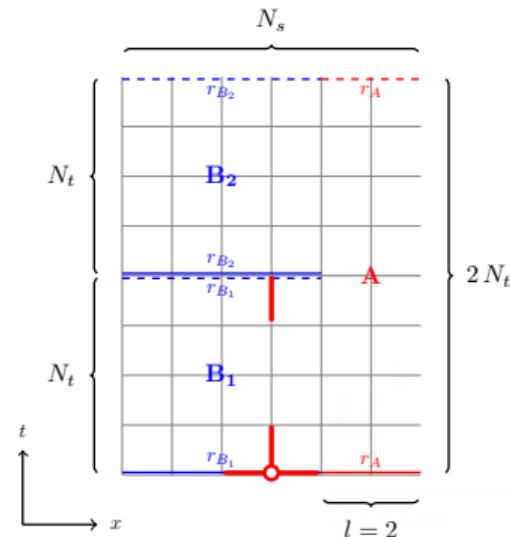
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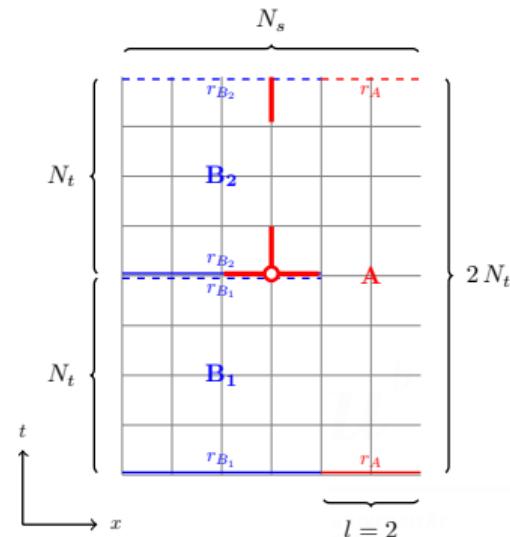
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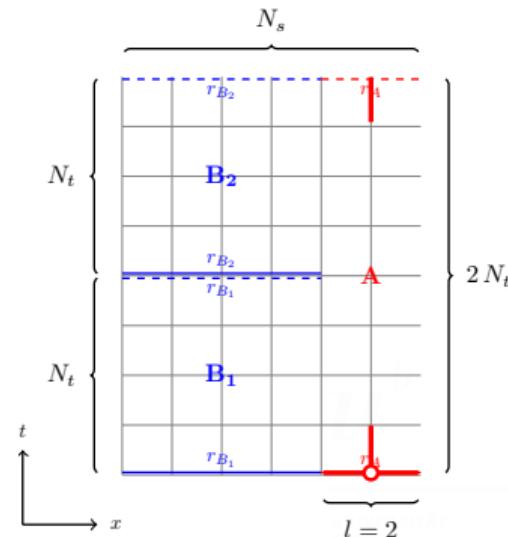
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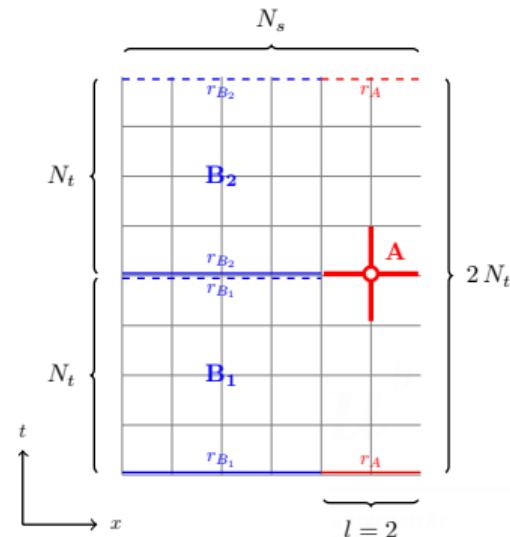
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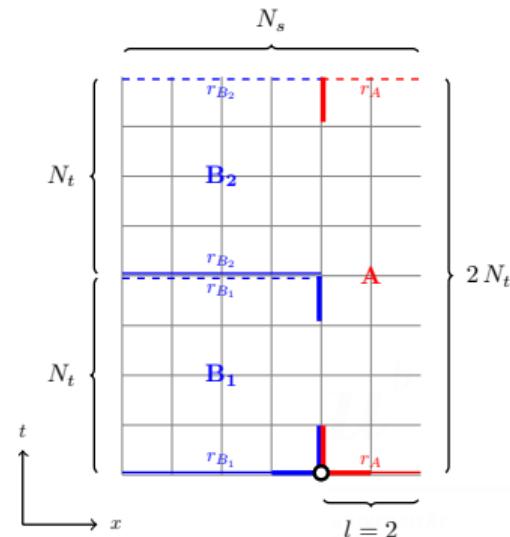
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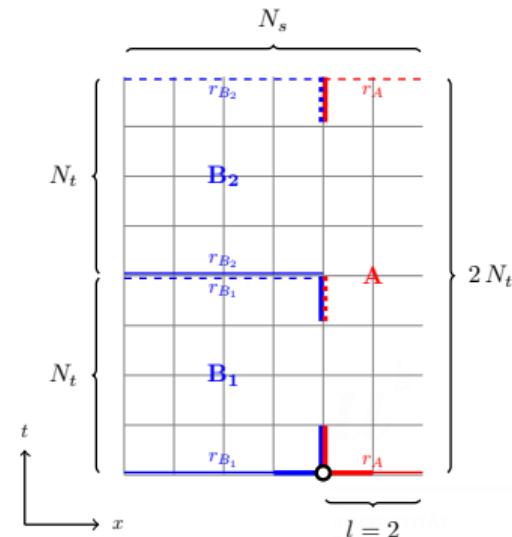
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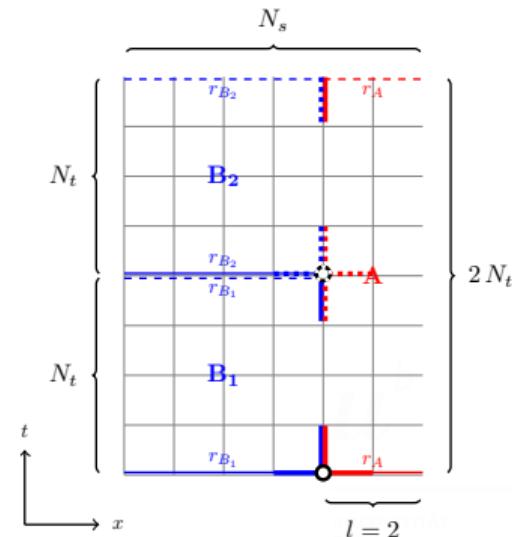
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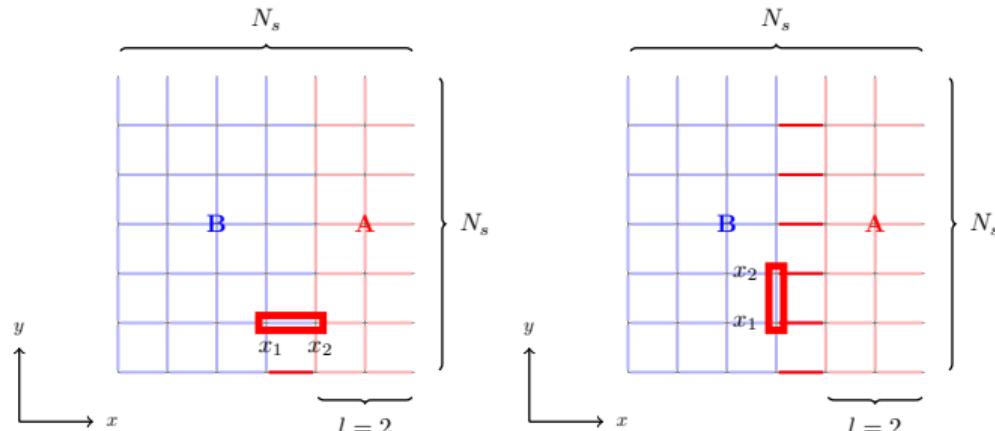
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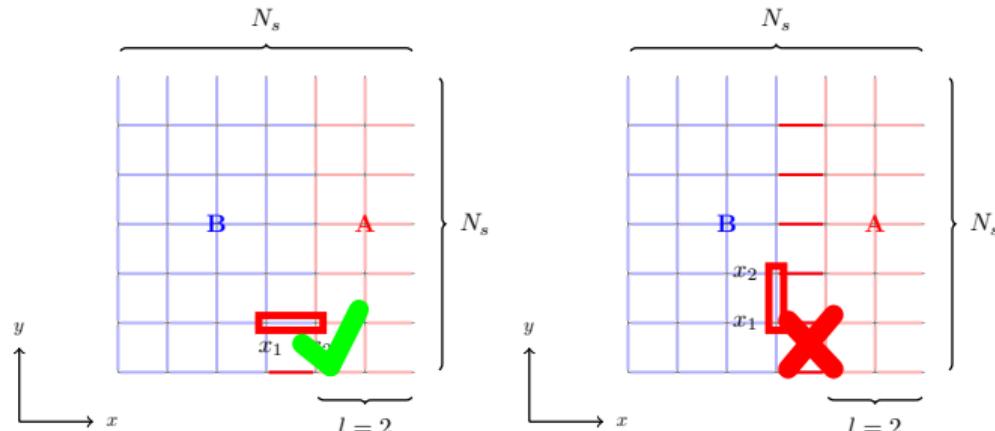
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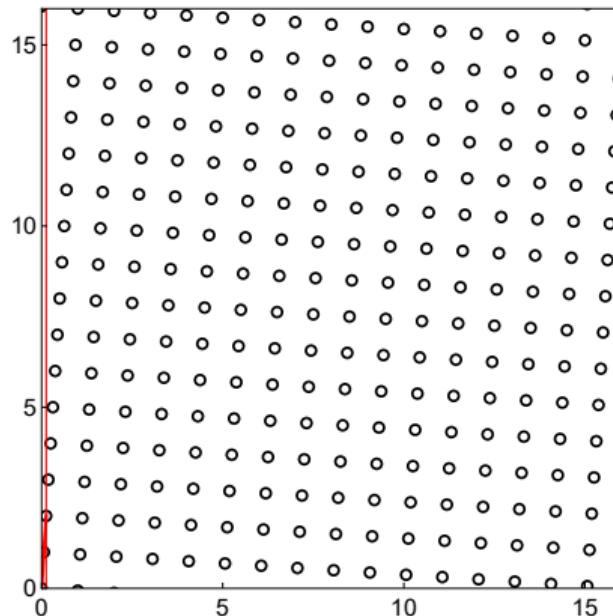
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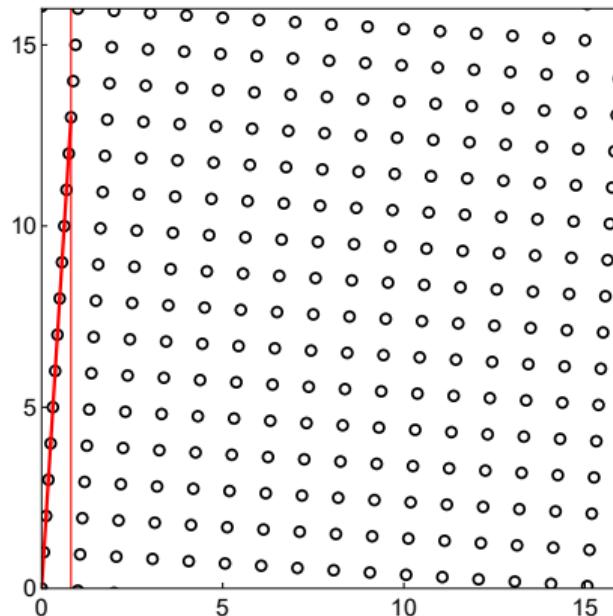
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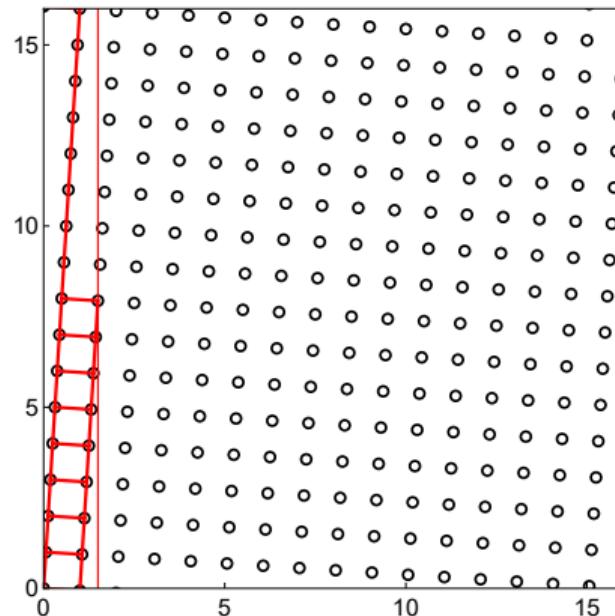
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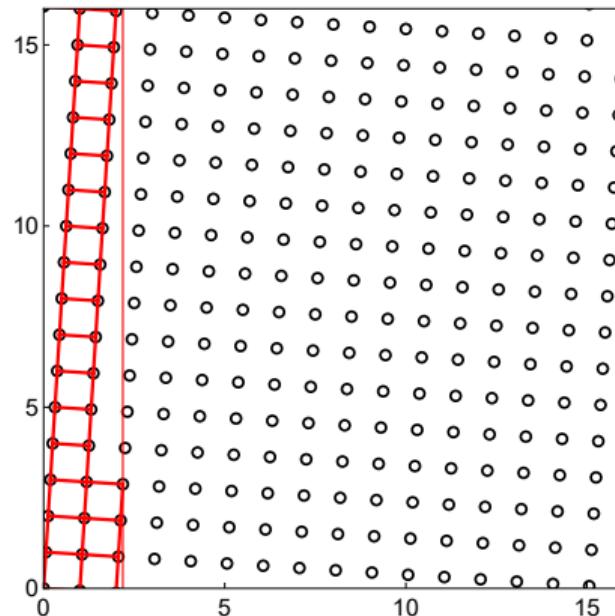
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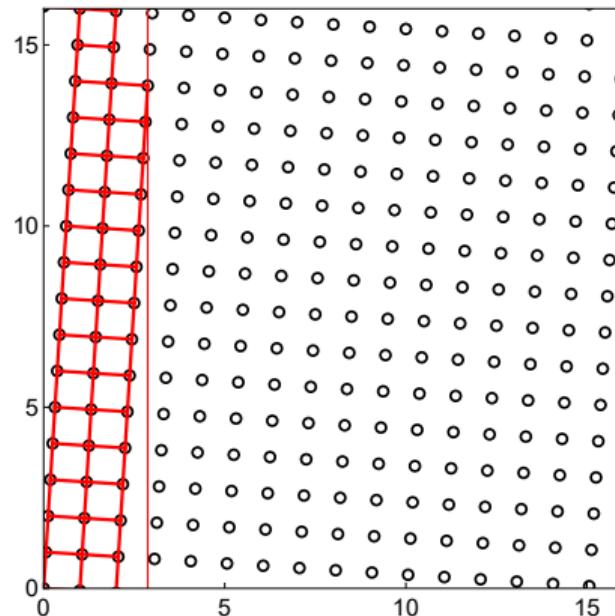
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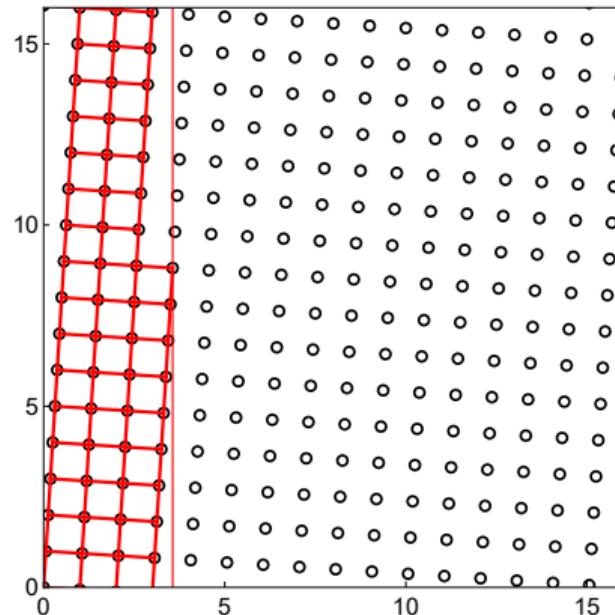
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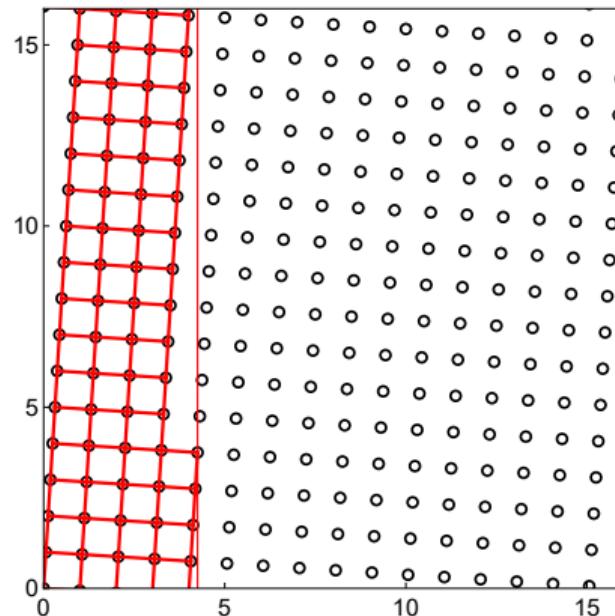
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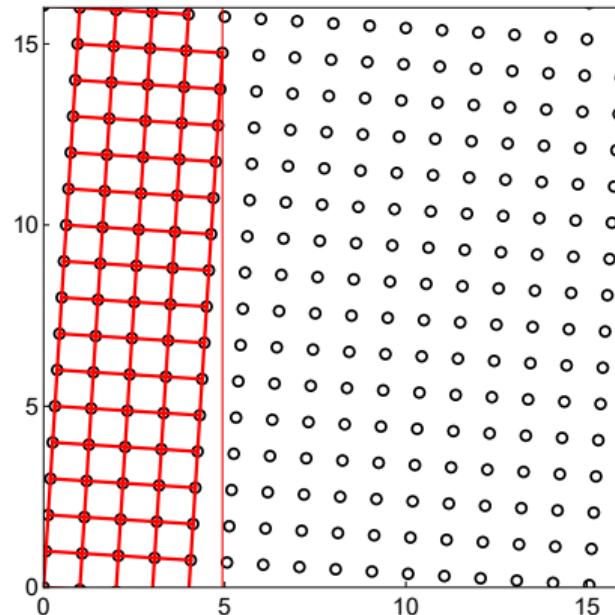
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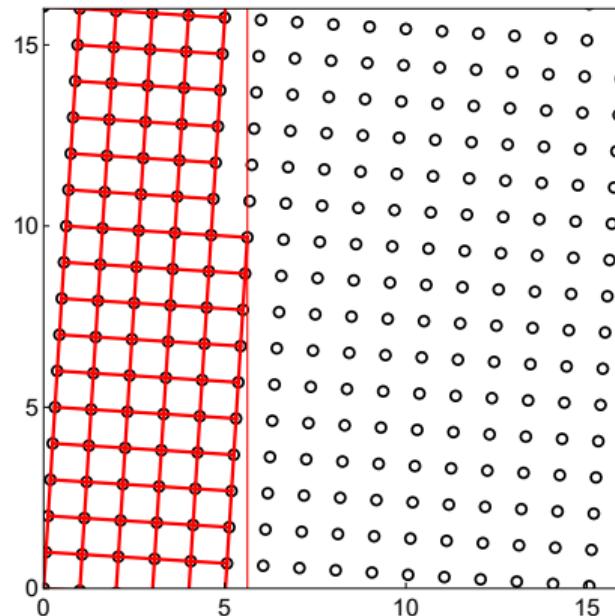
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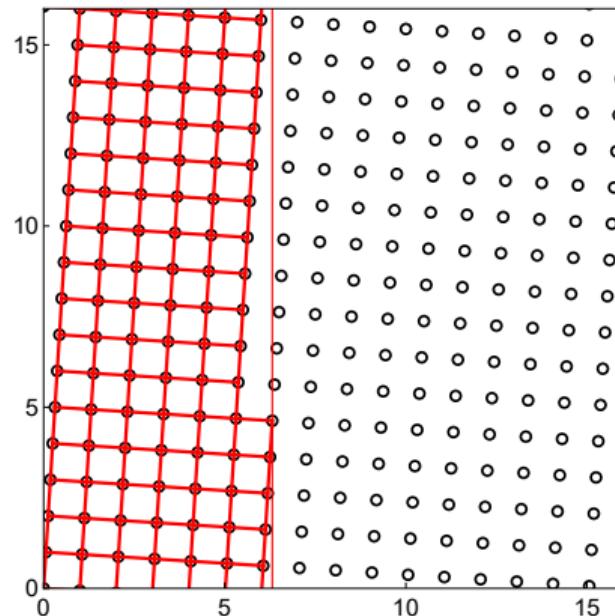
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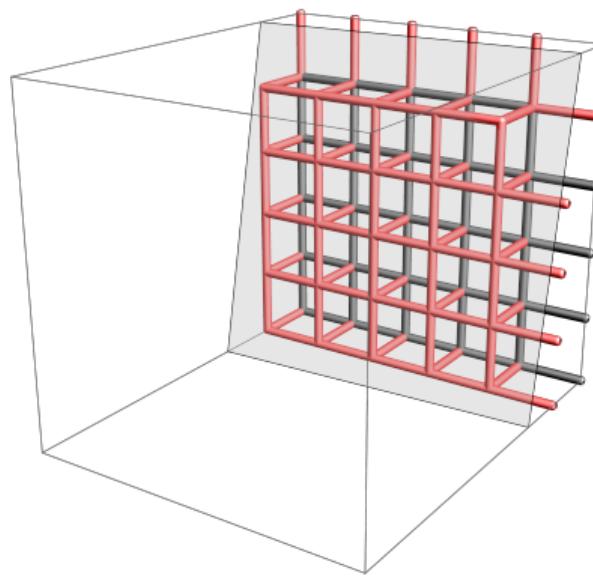
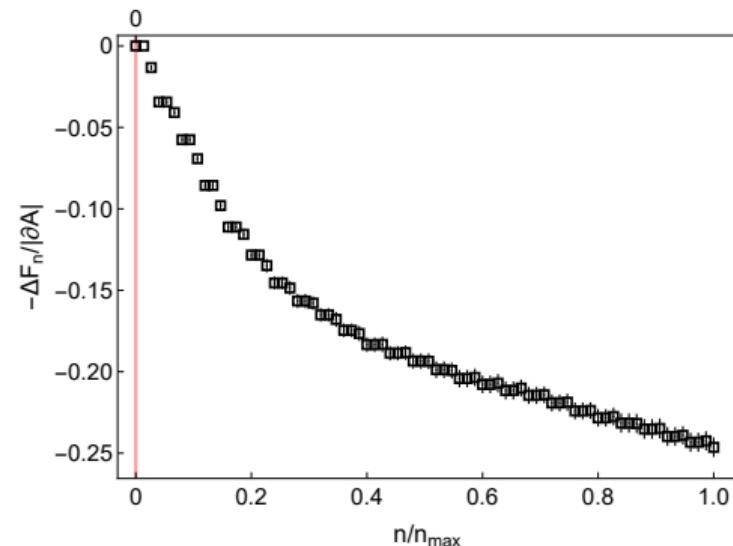


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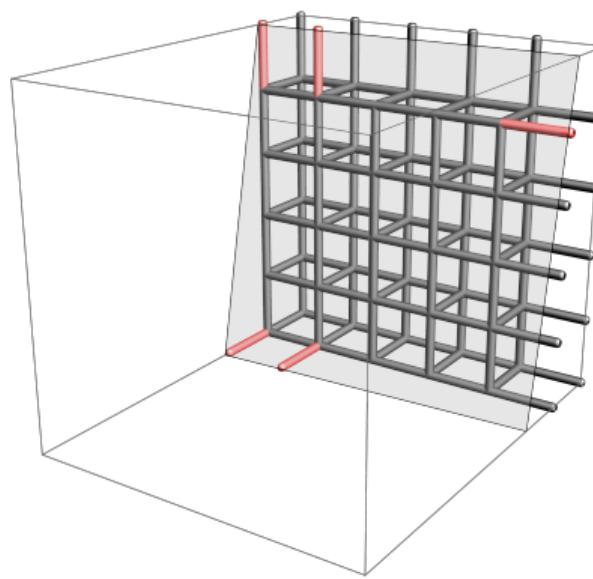
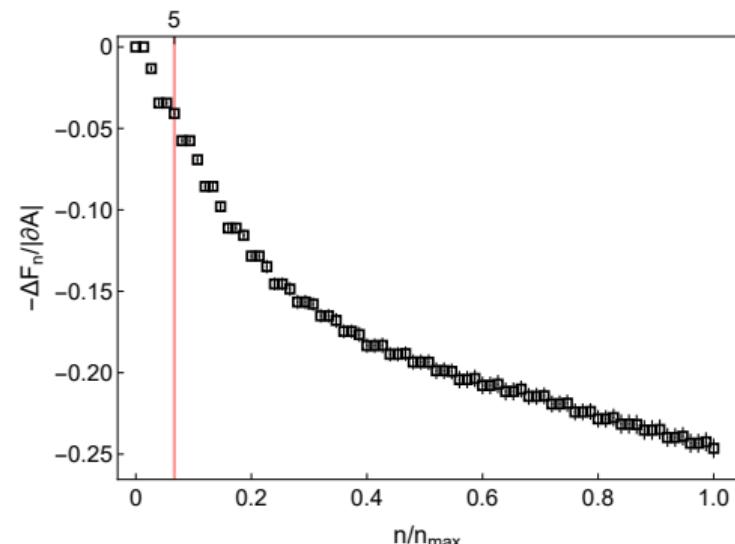


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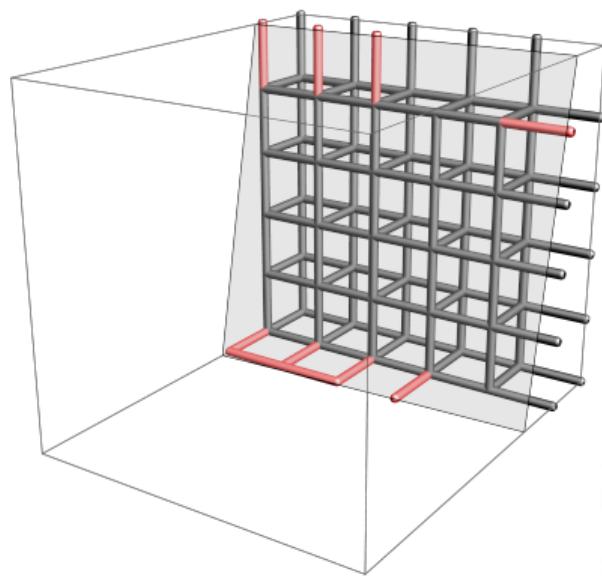
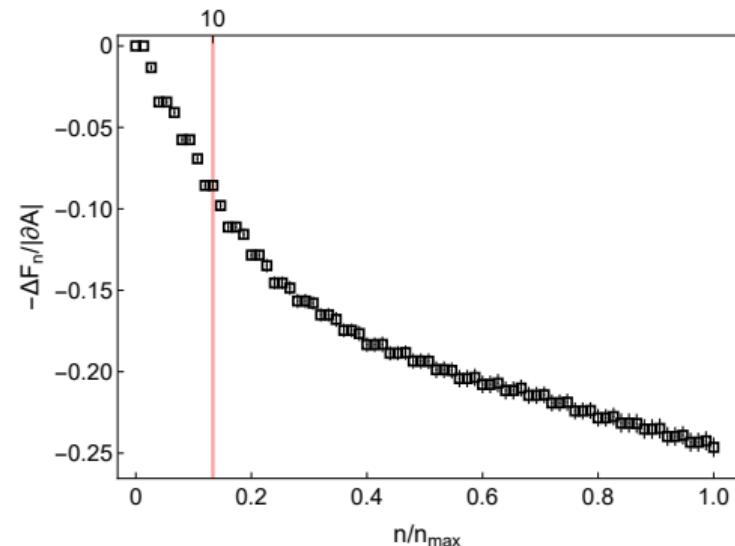


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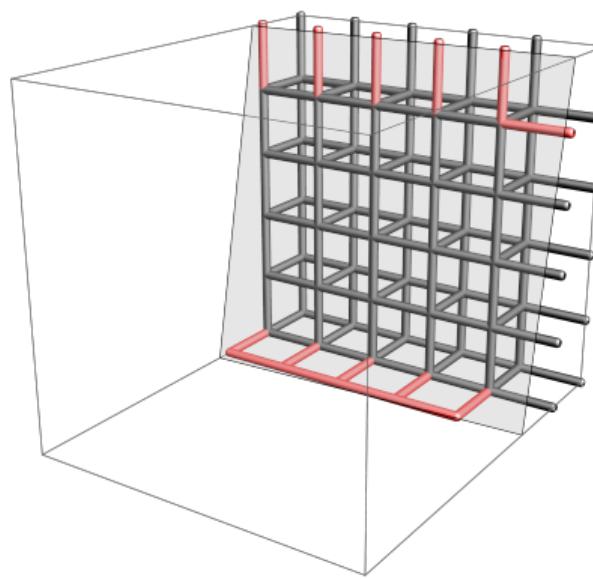
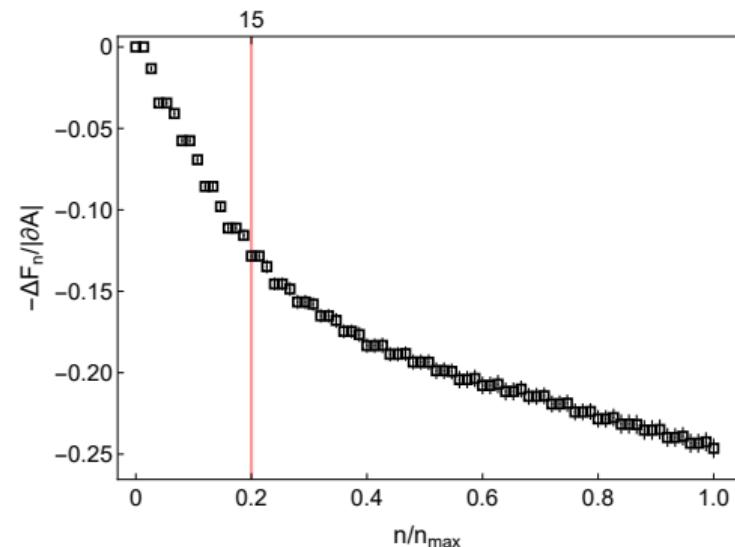


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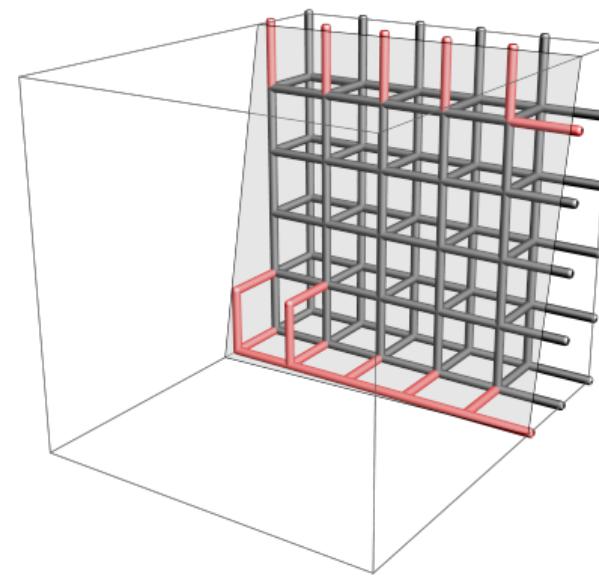
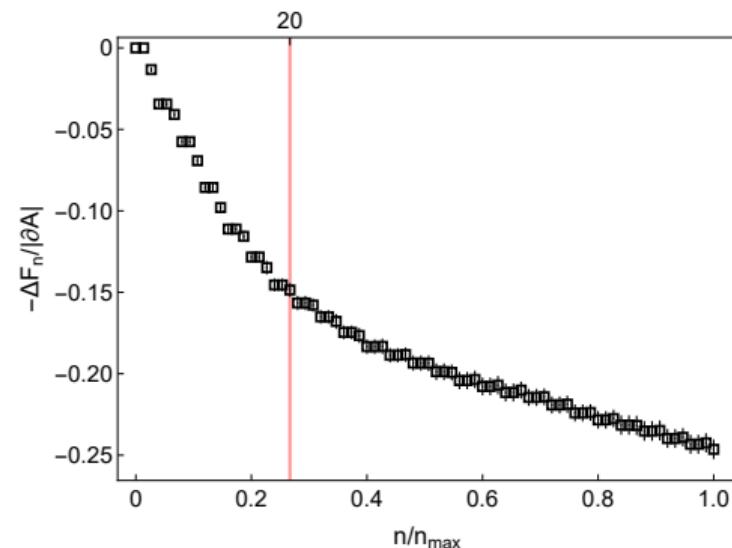


Entangling surface deformation method

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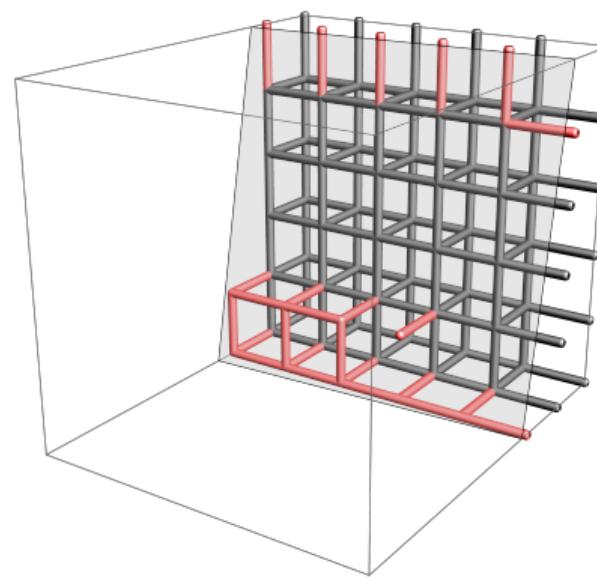
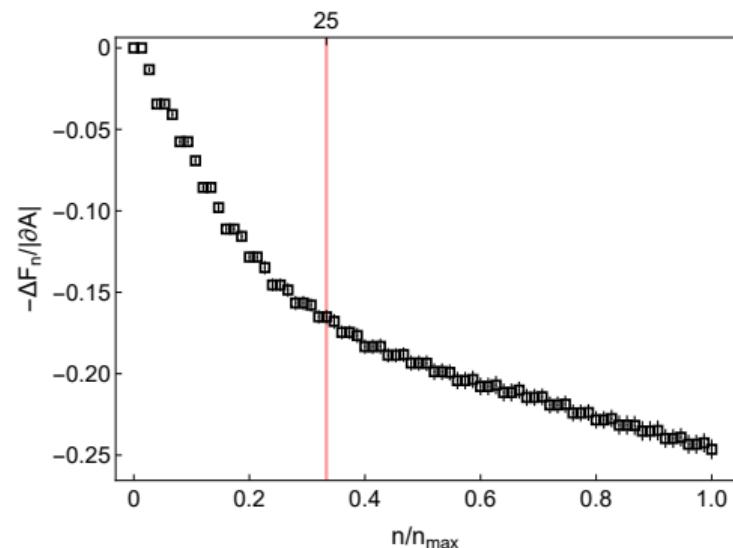


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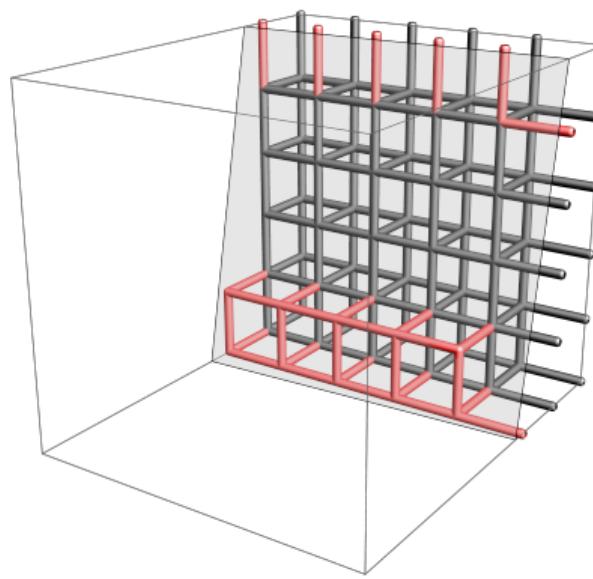
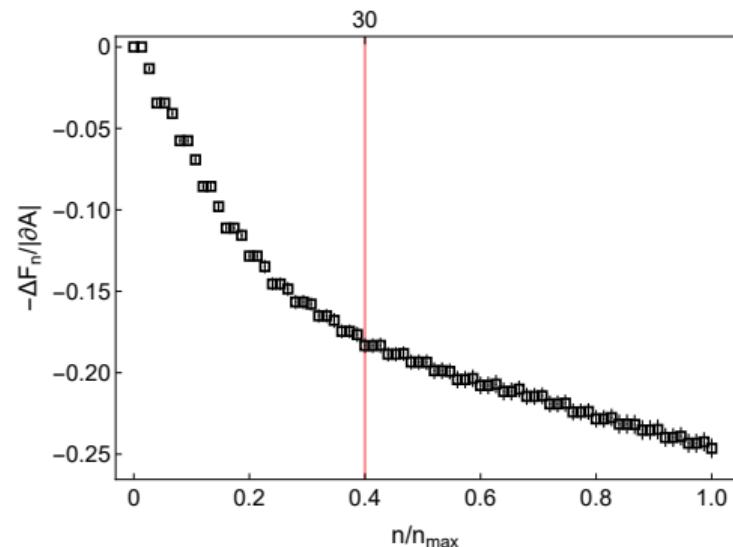


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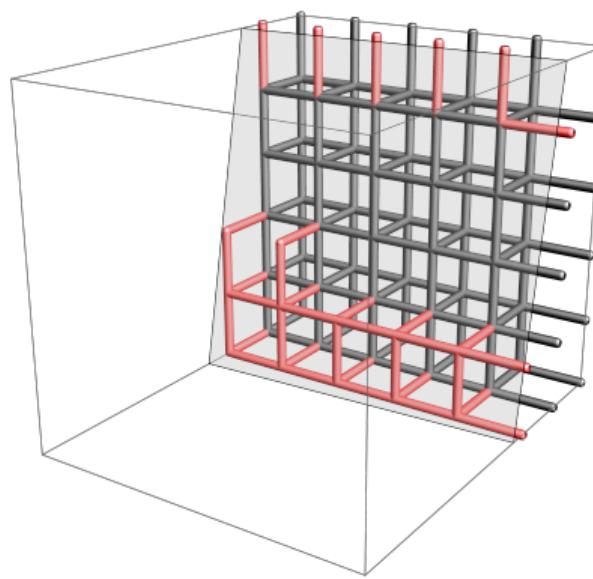
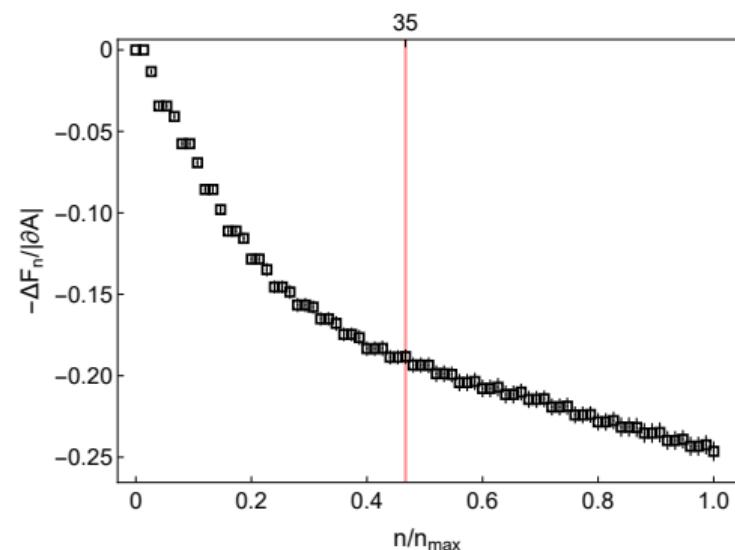


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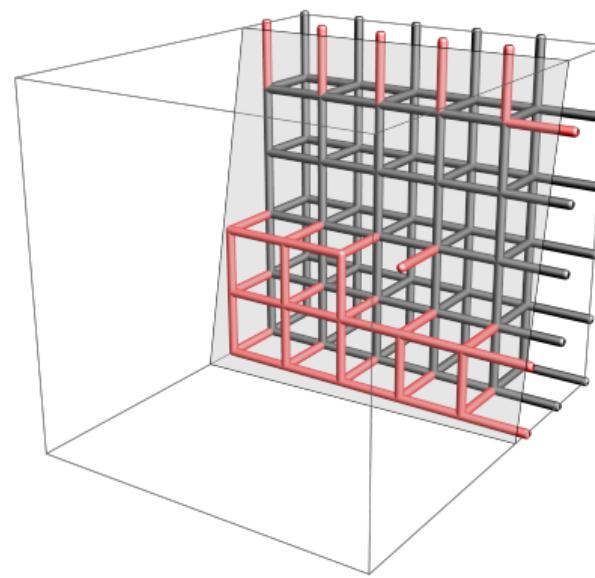
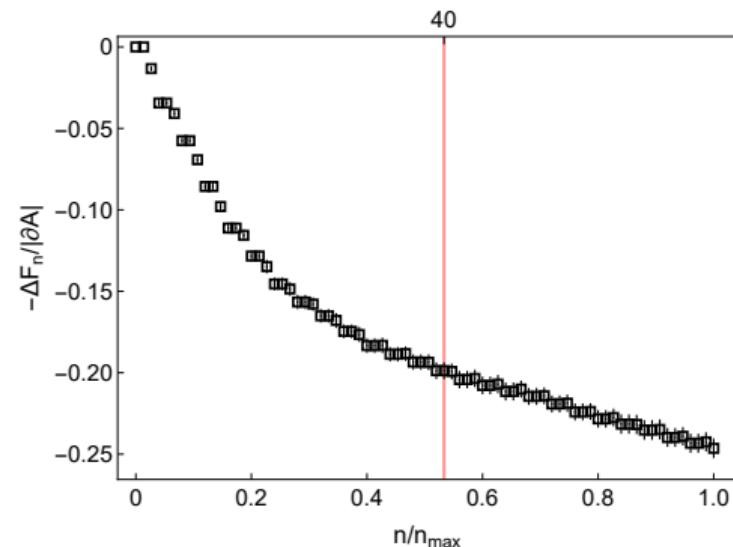


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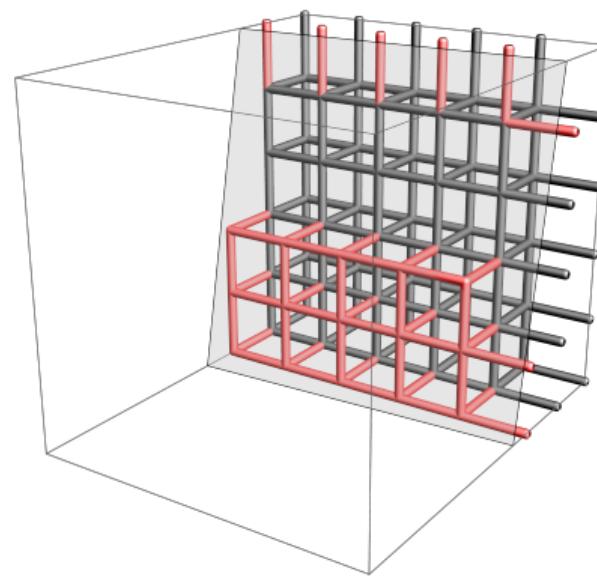
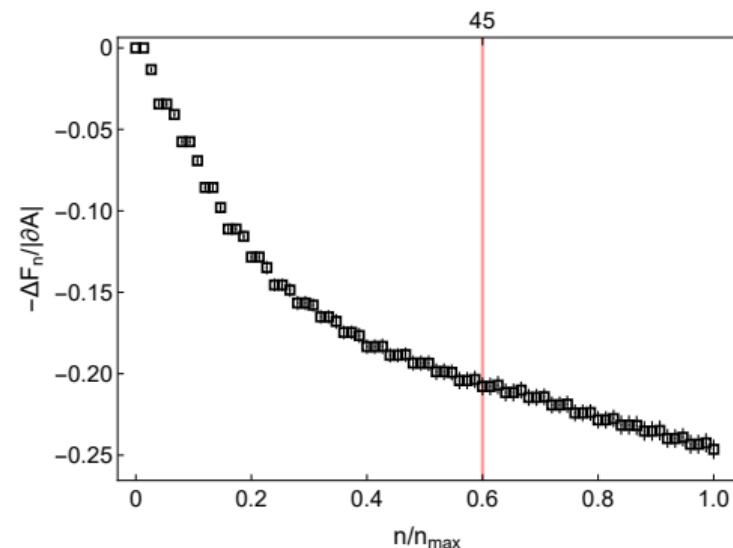


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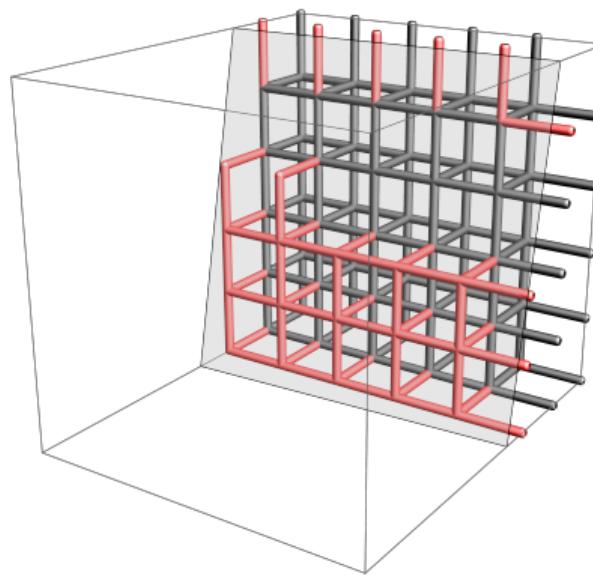
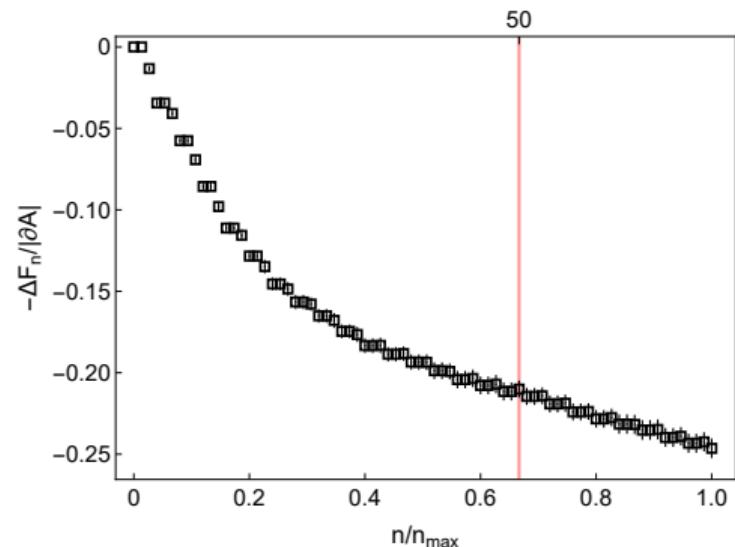


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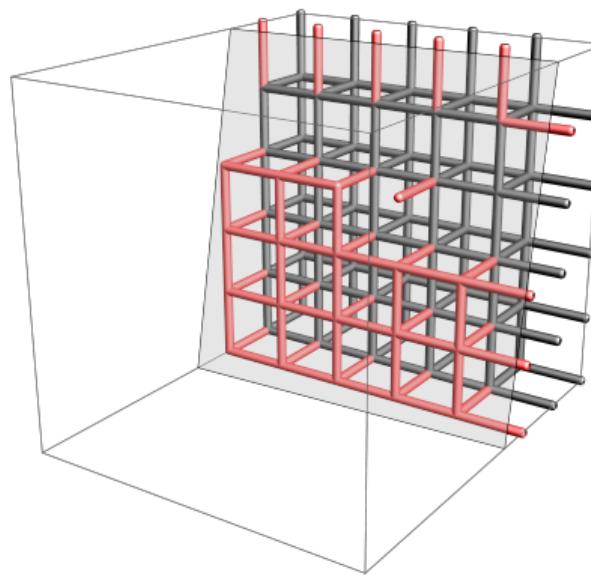
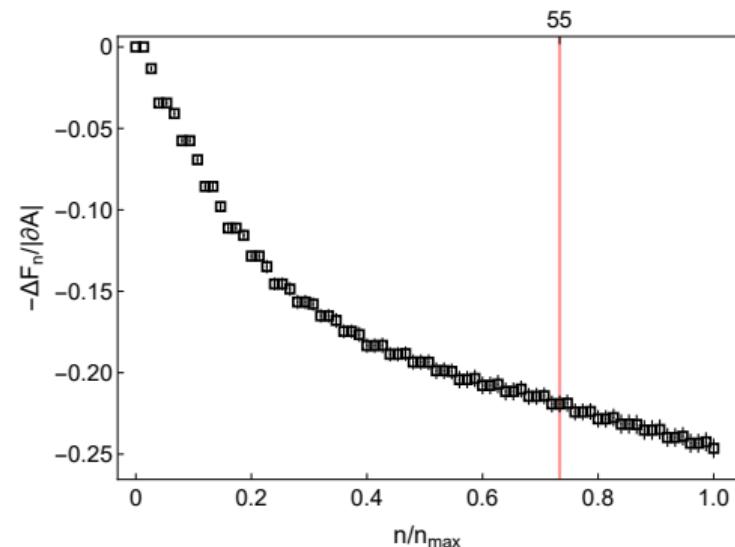


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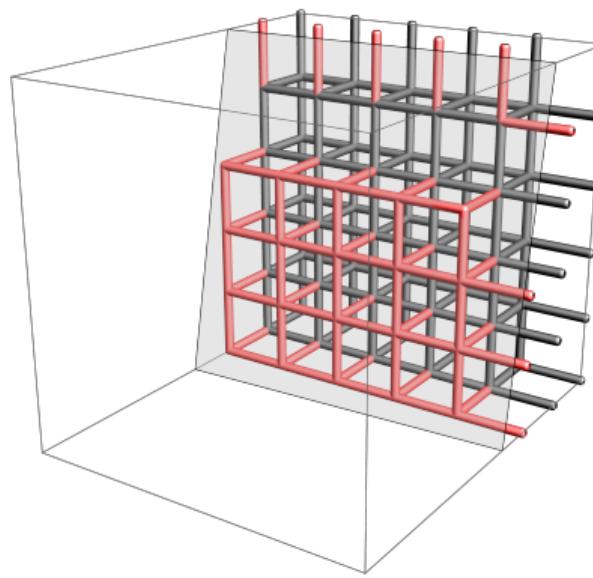
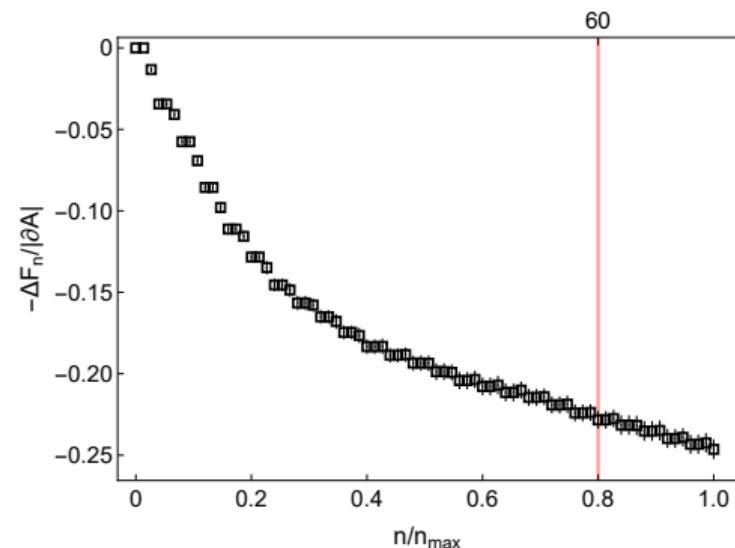


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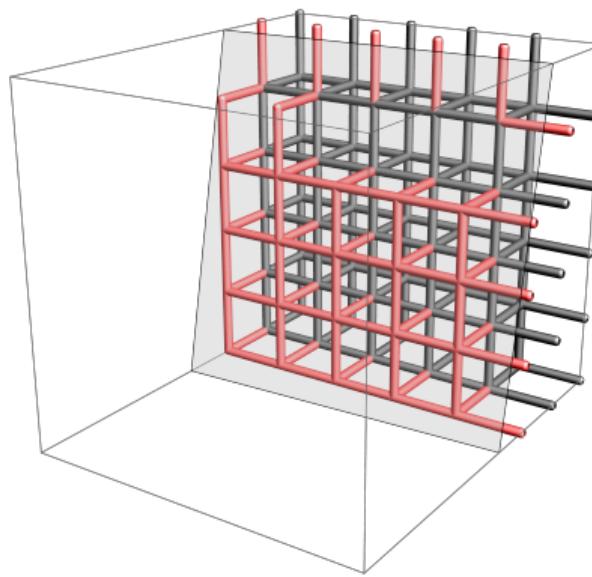
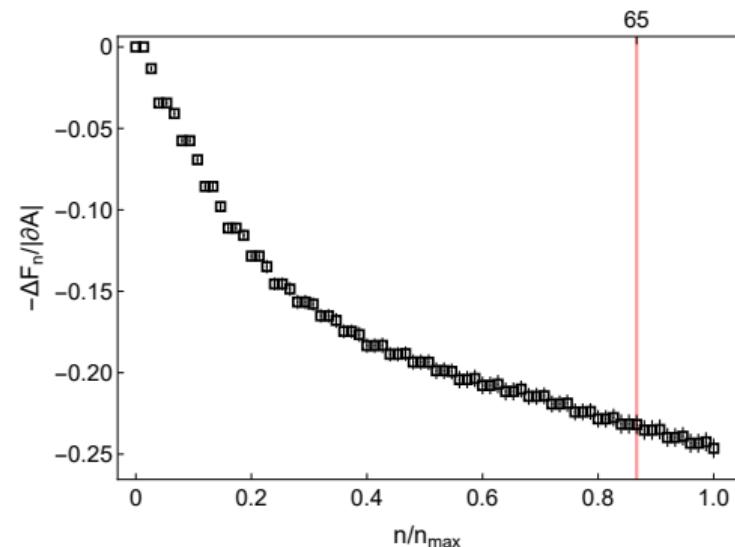


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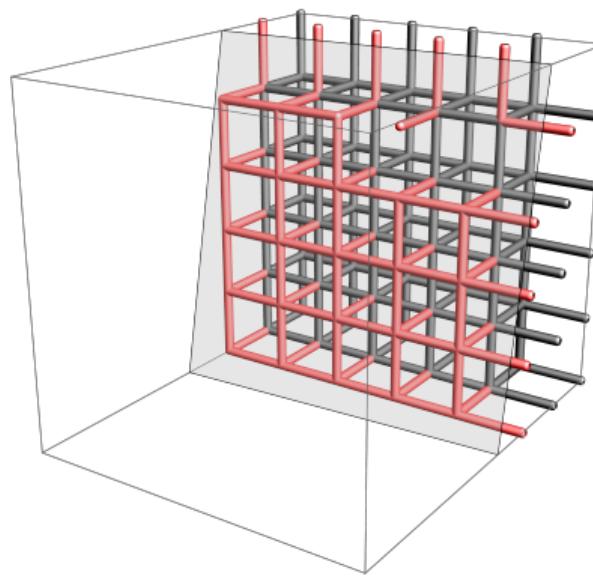
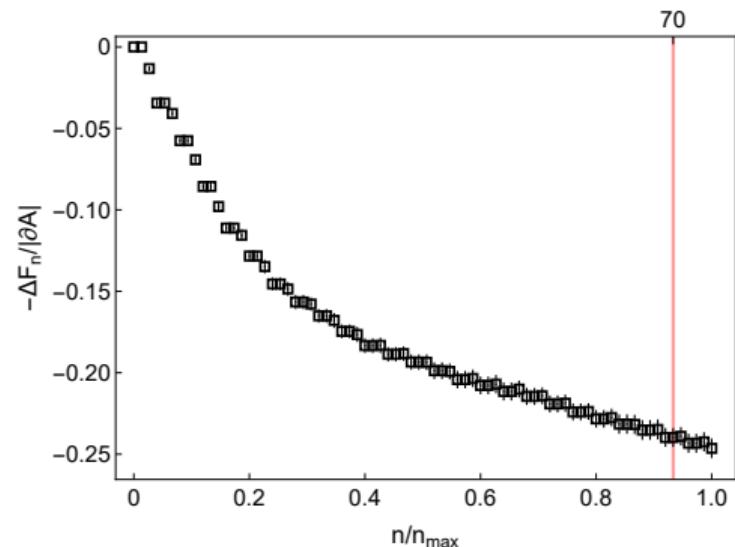


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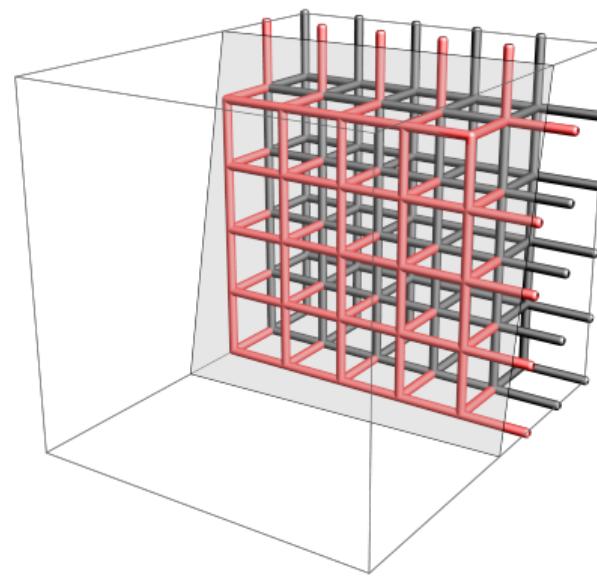
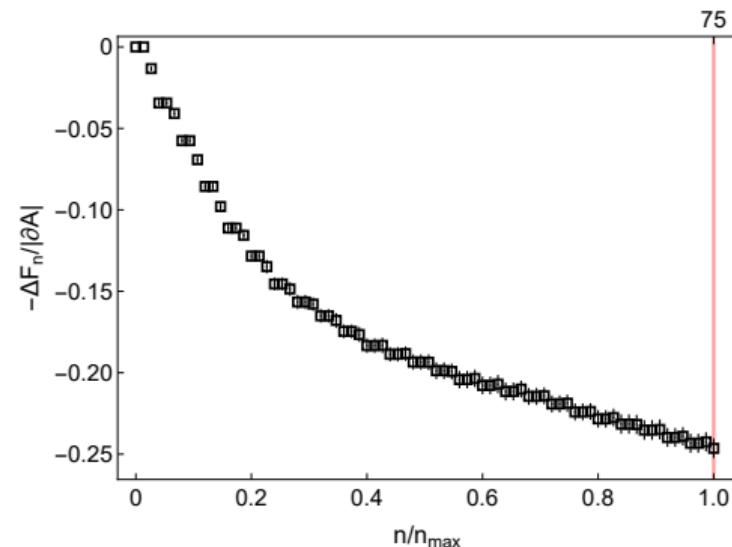


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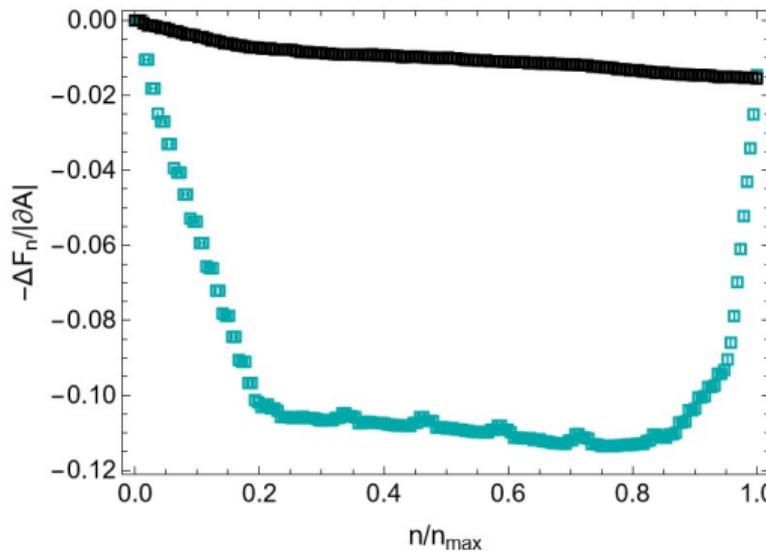
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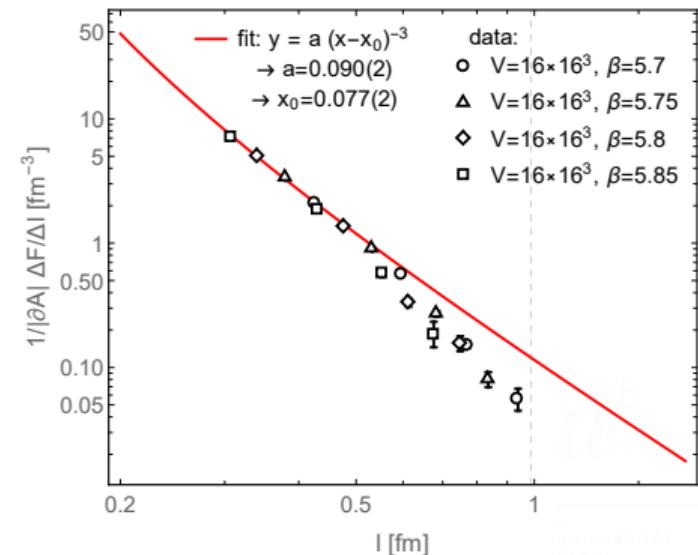
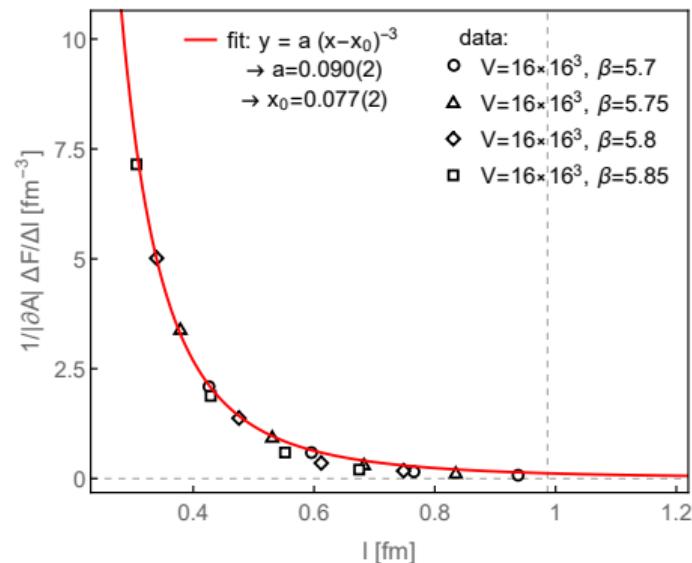
comparison of boundary update methods: [non-tilted lattice](#) \longleftrightarrow tilted lattice



Results

Results in 4D [arXiv.2211.00425]

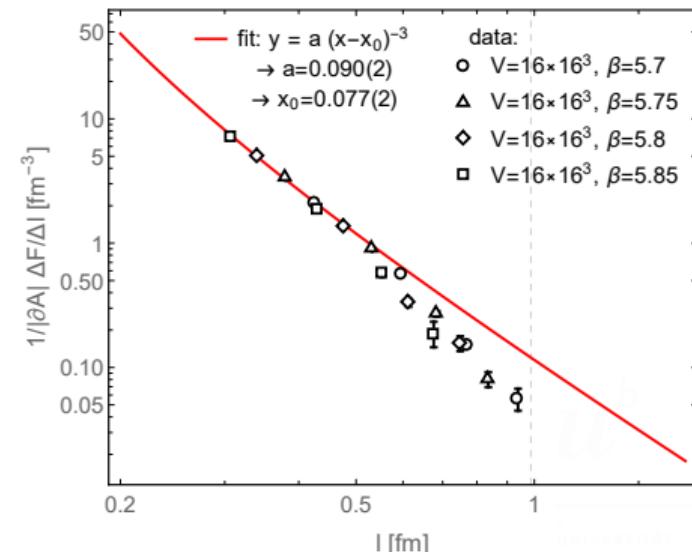
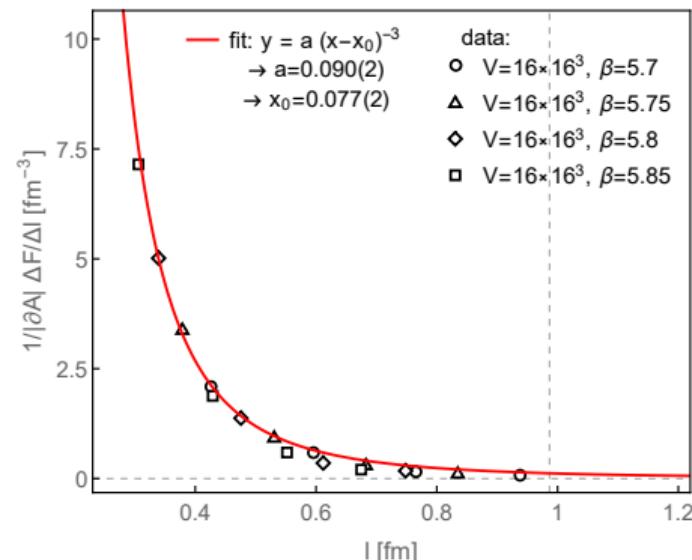
- Entanglement entropy change as function of entangling region width l for SU(3) on $N_s^3 \times 2 \cdot N_t$ lattice with $N_s = N_t = 16$, $\beta \in \{5.7, 5.75, 5.8, 5.85\}$.



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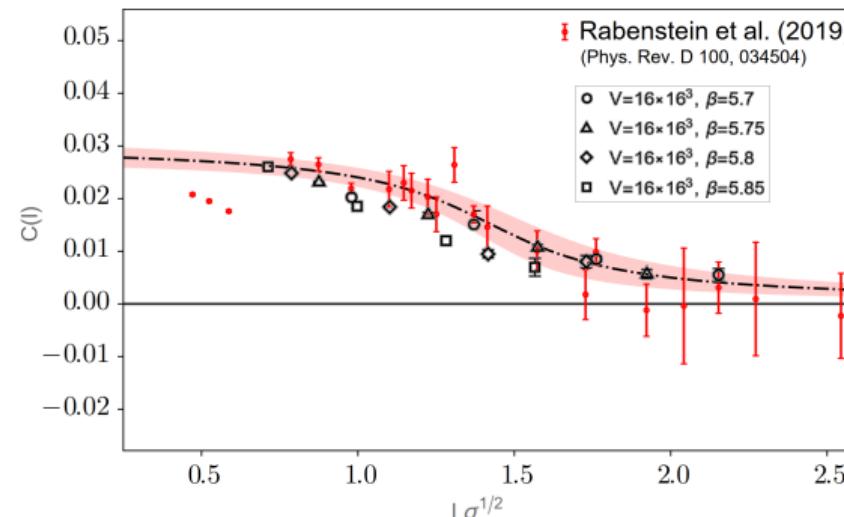
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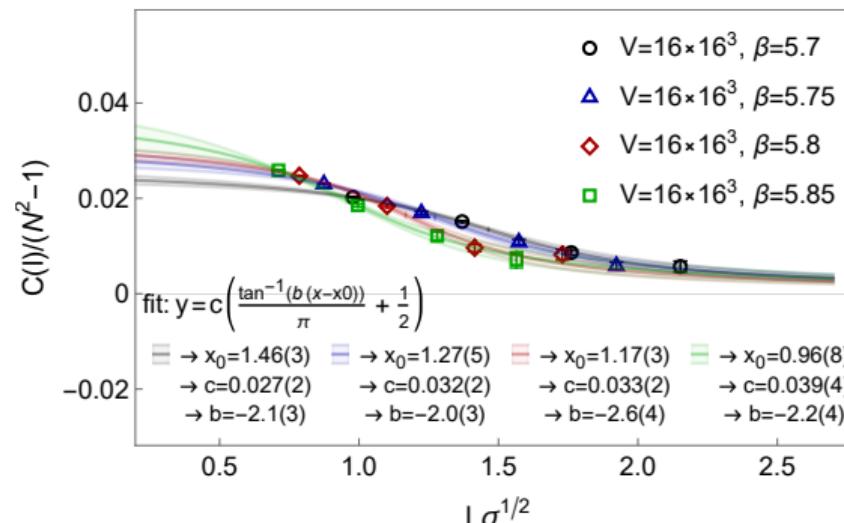
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Relation to thermal entropy [arXiv:2304.08949]

■ Thermodynamic entropy on the lattice

Lattice free energy: $F_L(N_t, V, N) = -\log(Z(N_t, V, N)) = N_t F(T(N_t), V, N)$

(spatial lattice volume V , temporal lattice size N_t , some charge N)

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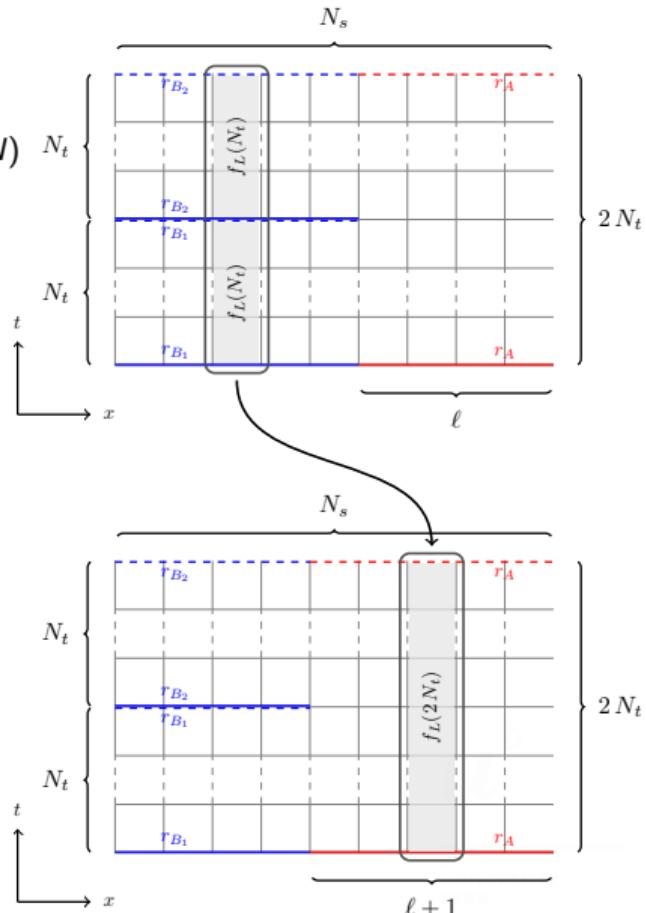
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Lattice entanglement entropy density at finite T , large ℓ :

$$\frac{1}{N_s^{(d-2)}} \left. \frac{\partial S_{EE}(\ell', N_t, N_s)}{\partial \ell'} \right|_{\ell'=\ell+1/2} \approx f_L(2 N_t) - 2 f_L(N_t)$$



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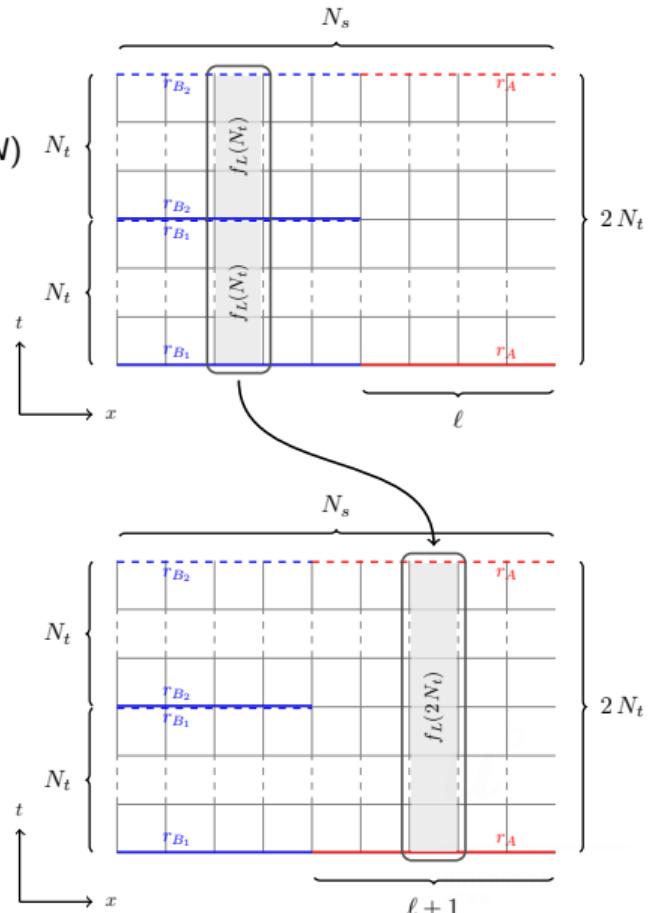
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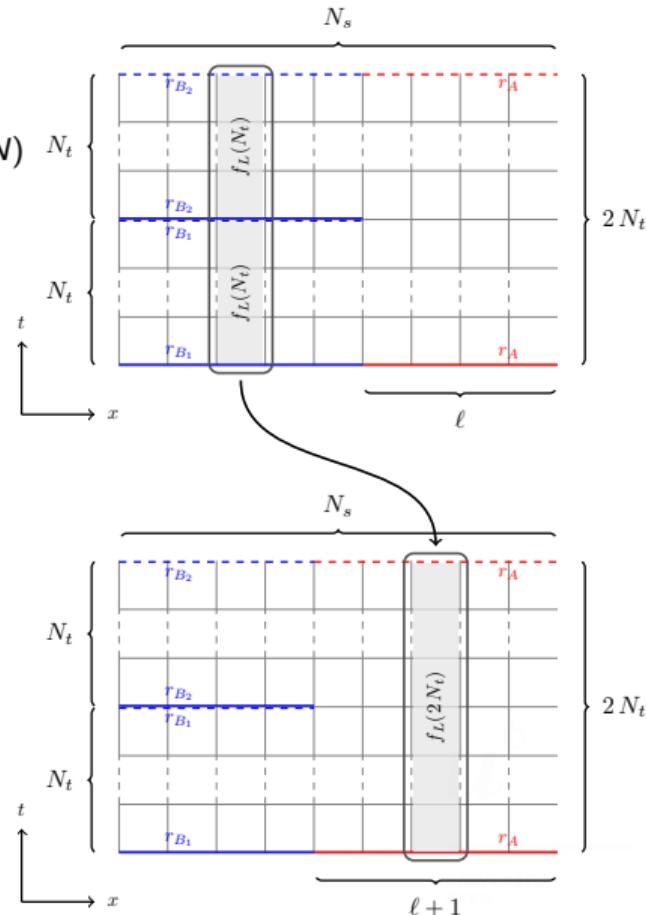
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thermal entropy density of region A : $s_{\text{th}, A}$



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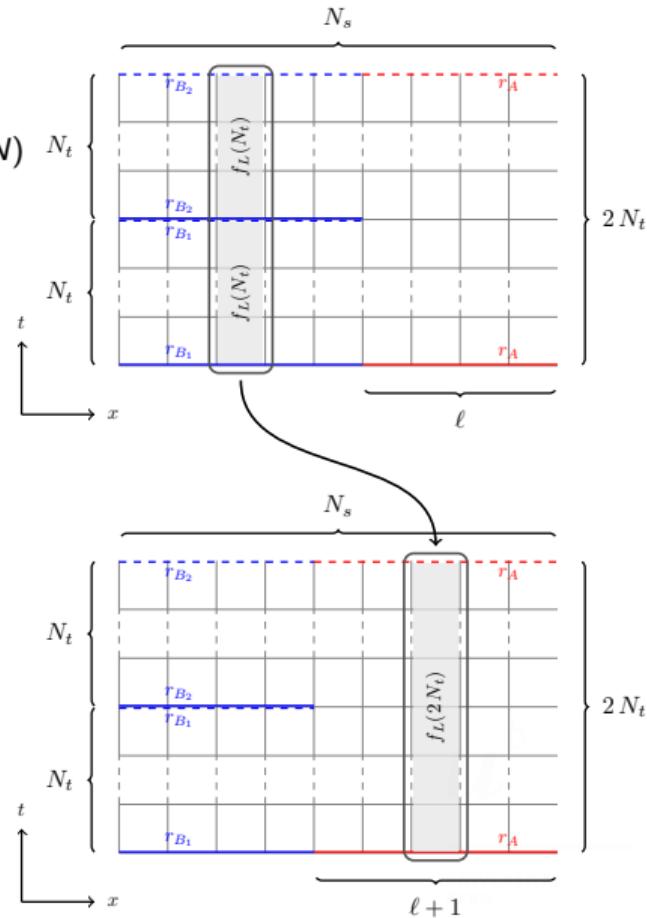
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→ Replica trick can be used to determine thermal entropy

(without having to know F_L or other intergration constants)



Results

Relation to thermal entropy [arXiv:2304.08949]

- Test in (2+1) dimensions at high temperature:

→ holography: $s_{\text{th},A} \propto T^{7/3}$ (Bekenstein-Hawking entropy) for $T/T_c \gg 1$

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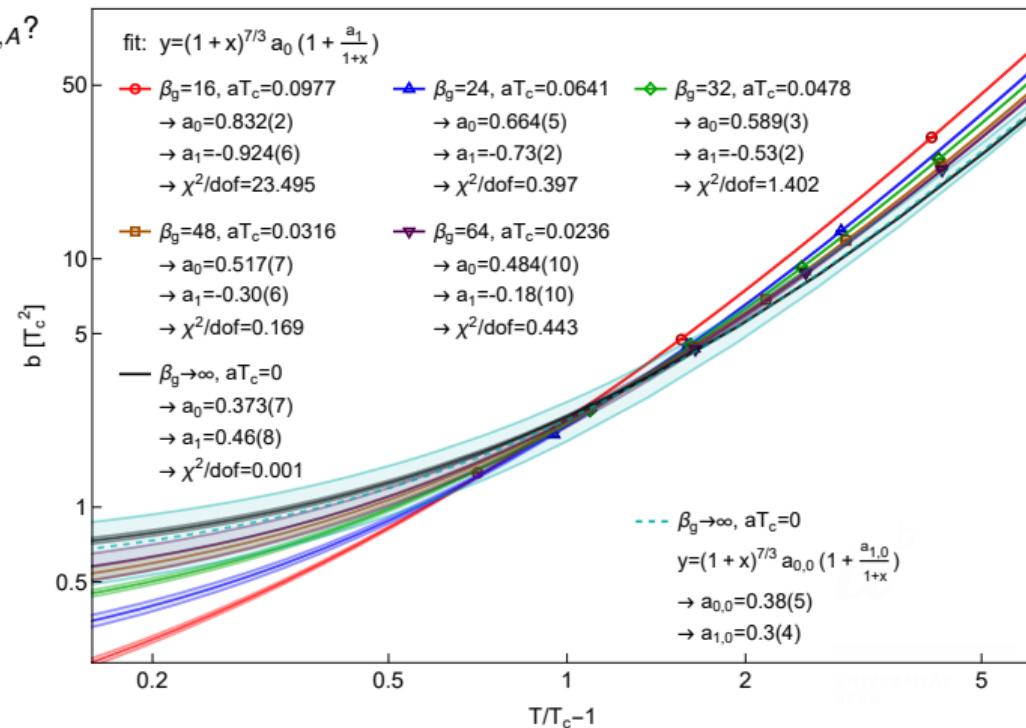
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Conclusions & outlook

Conclusions

- New method to determine entanglement measures (Rényi and entropies) in $SU(N)$ lattice gauge theories.
 - No more free energy barriers when using "tilted lattice".
 - significant error reduction possible.
 - Comparison with literature results promising.
 - Replica trick can be used to compute thermal entropy.

Outlook

- Application to further cases:
 - $SU(N)$, $N = 2, 3, 4, 5, \dots, ?$, $d = 3, 4$, $T = 0, T \neq 0$
 - different entangling region shapes; alternative entropy measures?
 - "metric reconstruction" (holography) for $SU(2), SU(3)$?
- improved simulation algorithm for "tilted lattice".

Thank you!