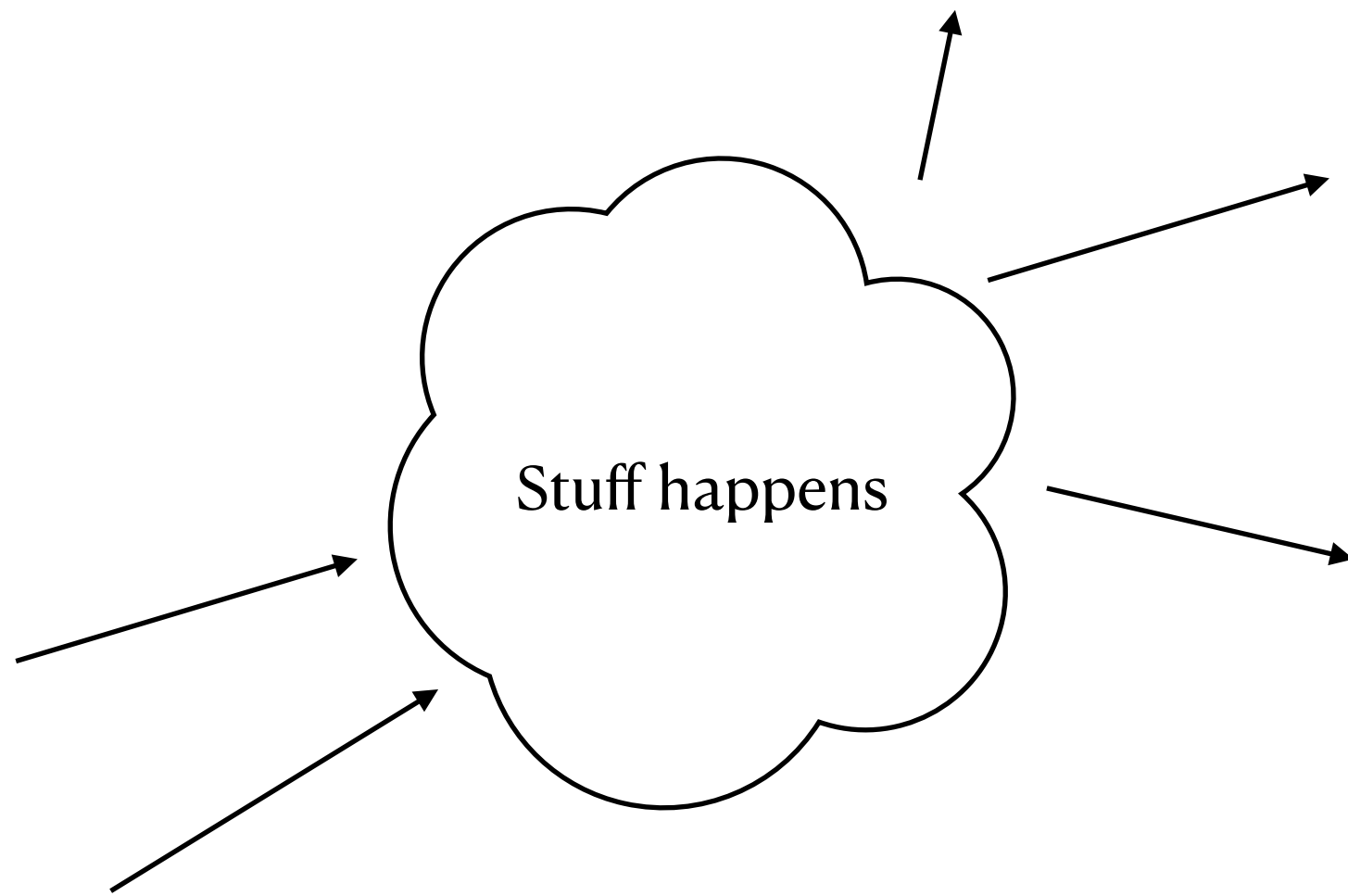


# Geometry of Scattering Amplitudes

Andreas Helset, CERN

# Scattering Amplitudes



Scattering matrix, "S-matrix"

$$\hat{S} : |\psi_{\text{in}}\rangle \rightarrow |\psi_{\text{out}}\rangle = \hat{S}|\psi_{\text{in}}\rangle$$

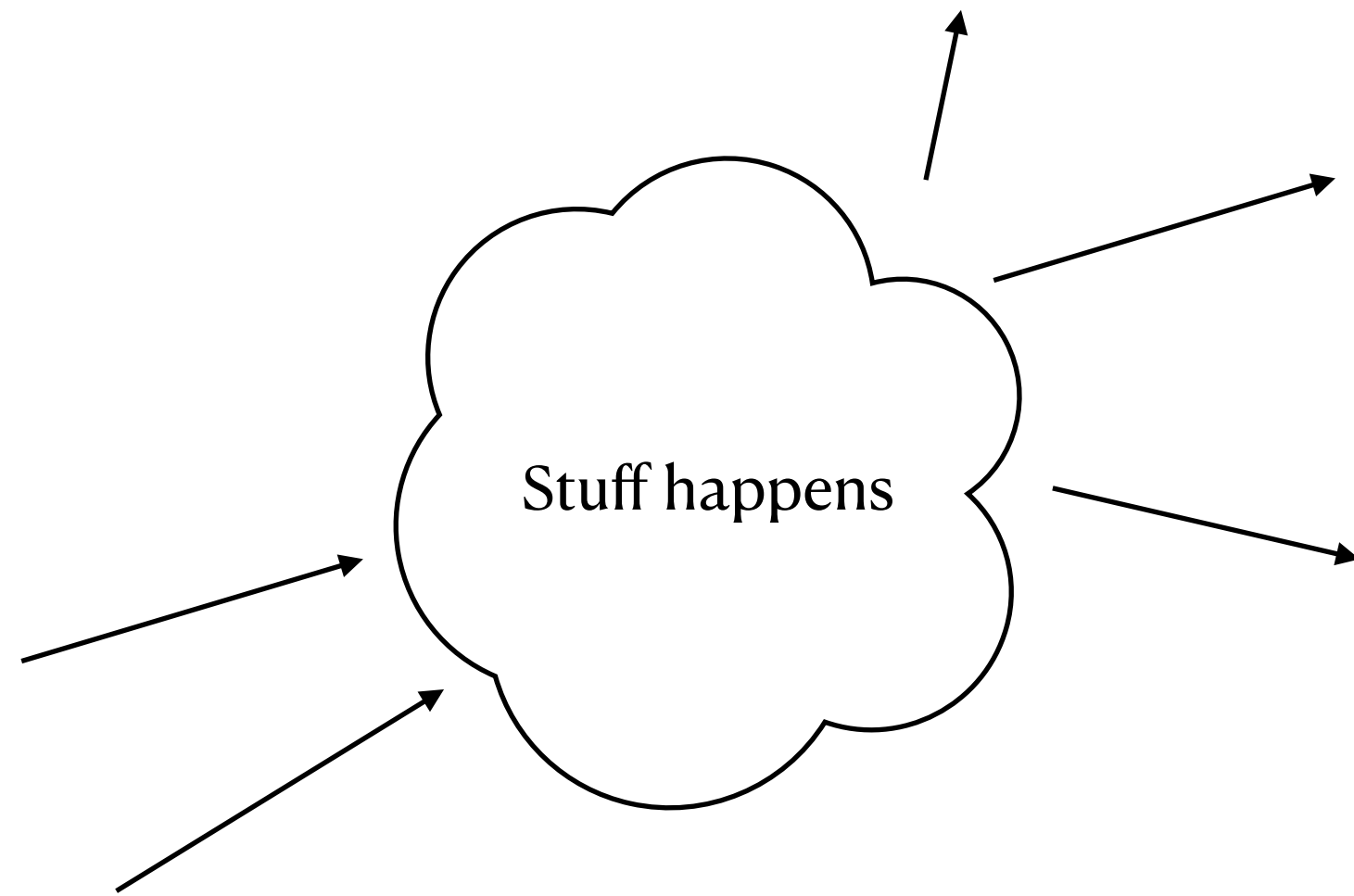
Key object for collider experiments

$$\frac{d\sigma}{d\Omega} \sim |\mathcal{A}|^2$$

Unitarity

$$\hat{S}\hat{S}^\dagger = 1$$

# Scattering Amplitudes



Scattering matrix, “S-matrix”

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Key object for collider experiments

$$\frac{d\sigma}{d\Omega} \sim |\mathcal{A}|^2$$

QFT 1:

Lagrangian

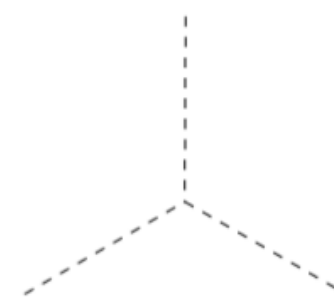
→

Feynman Rules

→

Scattering Amplitude

$$\mathcal{L}_{\text{int}} = \frac{\lambda}{3!} \phi^3$$



$$\mathcal{A} = \lambda$$

# Scattering Amplitudes

Surprising simplicity!

Gluon amplitude from Feynman diagrams

-8-      FERMILAB-Pub-85/118-T

The diagrams  $D_2^G$  are listed below:

$$D_2^G(1) = \frac{\delta_2}{s_{14} s_{25} s_{36}} \left\{ \begin{aligned} & [(p_2 - p_5)(p_3 - p_6)] \cdot [(p_1 - p_4)(p_3 + p_6)] \\ & - [(p_2 - p_5)(p_3 + p_6)] \cdot [(p_1 - p_4)(p_3 - p_6)] \\ & + [(p_2 + p_5)(p_3 - p_6)] \cdot [(p_1 - p_4)(p_2 - p_5)] \end{aligned} \right\},$$

$$D_2^G(2) = \frac{1}{s_{25} s_{36}} \left\{ \begin{aligned} & 2E(p_2 - p_5, p_3 - p_6) - 2E(p_3 - p_6, p_2 - p_5) \\ & + \delta_2 \cdot [(p_2 - p_5)(p_3 - p_6)] \end{aligned} \right\},$$

$$D_2^G(3) = \frac{4}{s_{25} s_{36} t_{125}} \left\{ \begin{aligned} & [(p_1 + p_2 - p_5)(p_4 + p_3 - p_6)] \cdot E(p_2, p_3) \\ & - [(p_1 + p_2 - p_5)(p_4 - p_3 + p_6)] \cdot E(p_2, p_6) \\ & - [(p_1 - p_2 + p_5)(p_4 + p_3 - p_6)] \cdot E(p_5, p_3) \\ & + [(p_1 - p_2 + p_5)(p_4 - p_3 + p_6)] \cdot E(p_5, p_6) \\ & - [p_1(p_2 - p_5)] \cdot E(p_3 - p_6, p_3 + p_6) \\ & - [p_4(p_3 - p_6)] \cdot E(p_2 + p_5, p_2 - p_5) \\ & + \delta_2 \cdot [p_1(p_2 - p_5)] \cdot [p_4(p_3 - p_6)] \end{aligned} \right\},$$

... + 8 more pages

# Scattering Amplitudes

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Gluon amplitude from Feynman diagrams

-8- FERMILAB-Pub-85/118-T

The diagrams  $D_2^G$  are listed below:

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Parke-Taylor amplitude

$$A_n = \frac{\langle 12 \rangle^4}{\langle 12 \rangle \langle 23 \rangle \cdots \langle n1 \rangle}$$

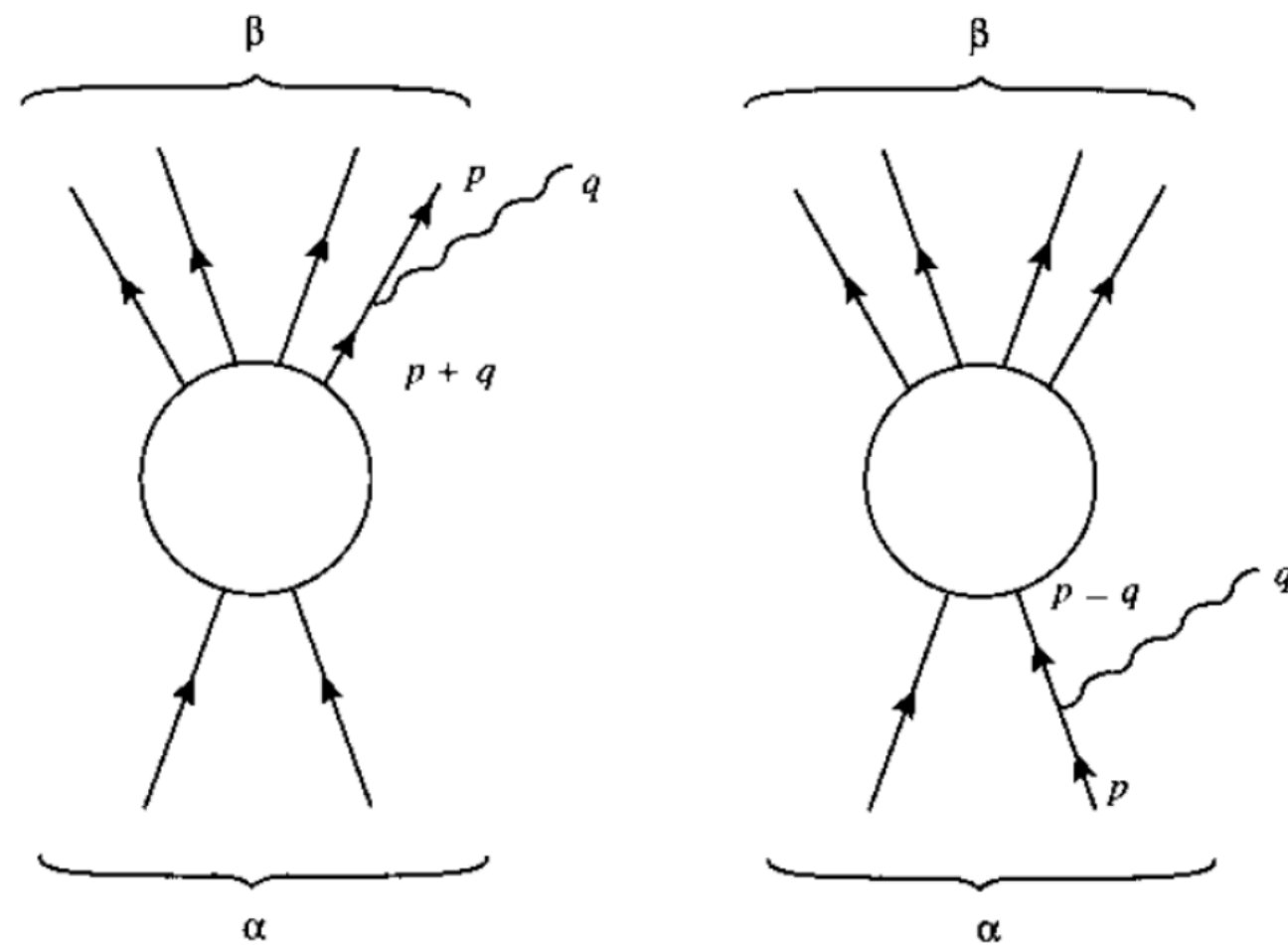
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# Scattering Amplitudes

## Universal relations!

Soft photon theorem

$$\lim_{q \rightarrow 0} \mathcal{A}_{n+1} = \sum_{i=1}^n Q_i \frac{\epsilon \cdot p_i}{q \cdot p_i} \mathcal{A}_n$$



Gauge invariance implies charge conservation

$$\mathcal{A}_{n+1}|_{\epsilon \rightarrow q} = 0 \Rightarrow \sum_{i=1}^n Q_i = 0$$

# Scattering Amplitudes

## Hidden structure!

Gauge theory

$$\mathcal{A} = \sum_{i \in \Gamma} \frac{c_i n_i}{d_i}$$

$$c_i + c_j = c_k$$

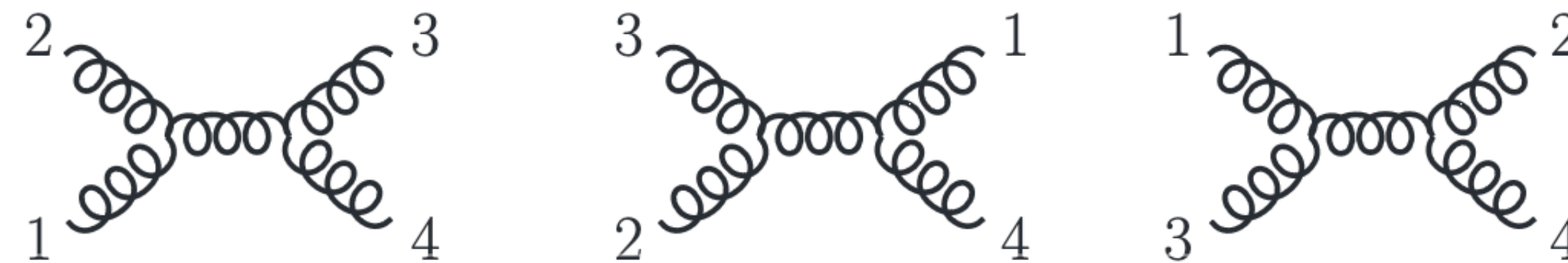


Figure 3: The three Feynman diagrams corresponding to the  $s$ ,  $t$  and  $u$  channels.

# Scattering Amplitudes

## Hidden structure!

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$$\mathcal{A} = \sum_{i \in \Gamma} \frac{c_i n_i}{d_i}$$

Color-kinematics duality

Bern, Carrasco, Johansson '08

$$c_i + c_j = c_k$$

$$n_i + n_j = n_k$$

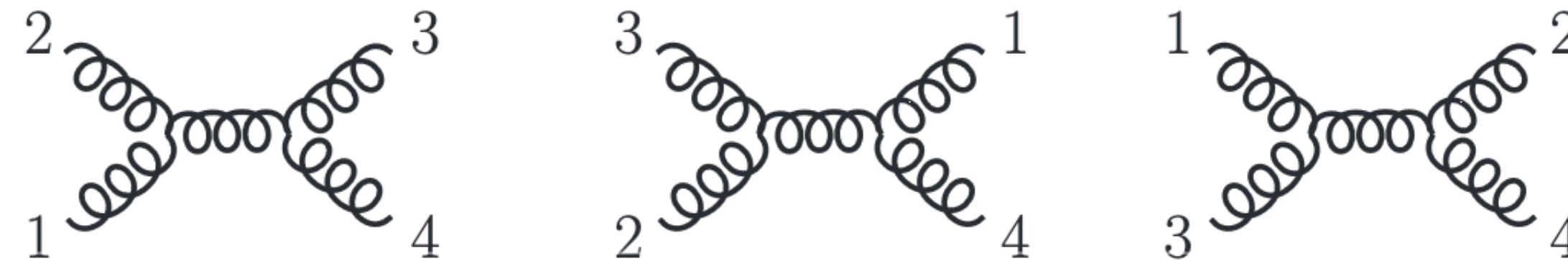


Figure 3: The three Feynman diagrams corresponding to the  $s$ ,  $t$  and  $u$  channels.



# Scattering Amplitudes

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Color-kinematics duality

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$$c_i + c_j = c_k$$

$$n_i + n_j = n_k$$

Gravity

$$\mathcal{M} = \sum_{i \in \Gamma} \frac{n_i n_i}{d_i}$$

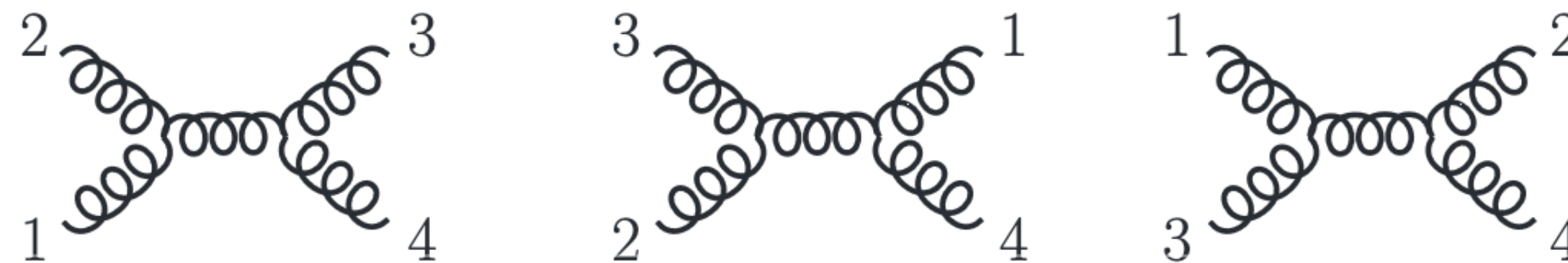


Figure 3: The three Feynman diagrams corresponding to the  $s$ ,  $t$  and  $u$  channels.

# Scattering Amplitudes

**Surprising simplicity!**

$$\begin{aligned}\mathcal{L} &= \frac{1}{2}(\partial_\mu\phi^i\partial^\mu\phi^i) + \lambda_{3,ijk}\phi^k(\partial_\mu\phi^i\partial^\mu\phi^j) + \lambda_{4,ijkl}\phi^k\phi^l(\partial_\mu\phi^i\partial^\mu\phi^j) + \dots \\ &= \frac{1}{2}g_{ij}(\phi)(\partial_\mu\phi^i\partial^\mu\phi^j)\end{aligned}$$

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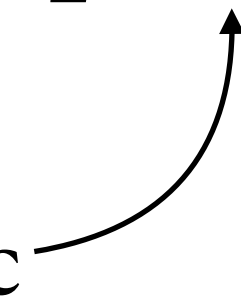
Field redefinition  $\phi^i \rightarrow \varphi^i(\phi)$

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Metric



Field redefinition

$$\phi^i \rightarrow \varphi^i(\phi)$$

Christoffel symbol

$$\Gamma_{jk}^i = \frac{1}{2} g^{il} (g_{lk,j} + g_{jl,k} - g_{jk,l})$$

Riemann curvature

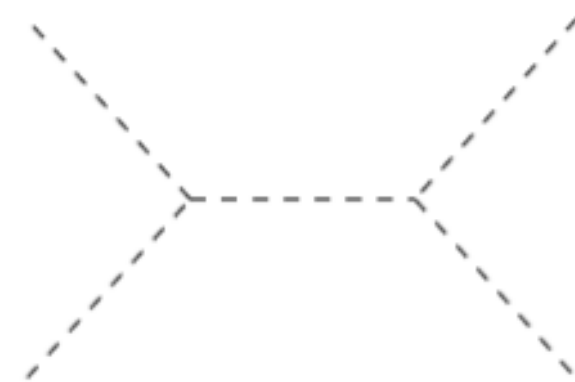
$$R^i_{jkl} = \Gamma^i_{lj,k} + \Gamma^i_{kn} \Gamma^n_{lj} - (k \leftrightarrow l)$$

# Scattering Amplitudes

Surprising simplicity!

$$\begin{aligned}\mathcal{L} &= \frac{1}{2}(\partial_\mu \phi^i \partial^\mu \phi^i) + \lambda_{3,ijk} \phi^k (\partial_\mu \phi^i \partial^\mu \phi^j) + \lambda_{4,ijkl} \phi^k \phi^l (\partial_\mu \phi^i \partial^\mu \phi^j) + \dots \\ &= \frac{1}{2} g_{ij}(\phi) (\partial_\mu \phi^i \partial^\mu \phi^j)\end{aligned}$$

Feynman diagrams



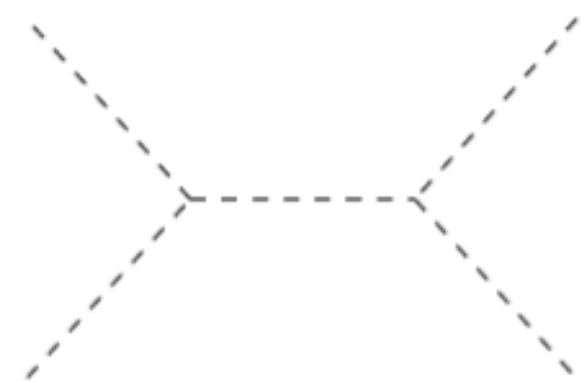
$$\mathcal{A}_4 = (\partial g)(p \cdot p') \frac{1}{s_{ij}} (\partial g)(p \cdot p') + \dots + (\partial^2 g)(p_i \cdot p_j)$$

# Scattering Amplitudes

Surprising simplicity!

$$\begin{aligned}\mathcal{L} &= \frac{1}{2}(\partial_\mu \phi^i \partial^\mu \phi^i) + \lambda_{3,ijk} \phi^k (\partial_\mu \phi^i \partial^\mu \phi^j) + \lambda_{4,ijkl} \phi^k \phi^l (\partial_\mu \phi^i \partial^\mu \phi^j) + \dots \\ &= \frac{1}{2} g_{ij}(\phi) (\partial_\mu \phi^i \partial^\mu \phi^j)\end{aligned}$$

Feynman diagrams



Geometric amplitude

$$\mathcal{A}_4 = R_{ikjl} s_{34} + R_{ijkl} s_{24}$$

$$\mathcal{A}_4 = (\partial g)(p \cdot p') \frac{1}{s_{ij}} (\partial g)(p \cdot p') + \dots + (\partial^2 g)(p_i \cdot p_j)$$

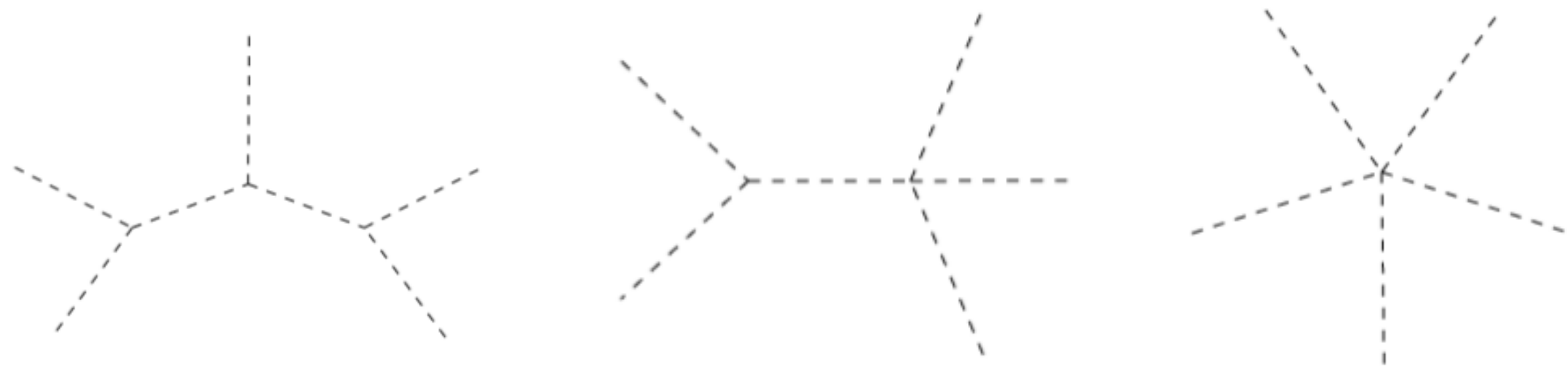
$$s_{ij} = (p_i + p_j)^2$$

# Scattering Amplitudes

Surprising simplicity!

$$\begin{aligned}\mathcal{L} &= \frac{1}{2}(\partial_\mu\phi^i\partial^\mu\phi^i) + \lambda_{3,ijk}\phi^k(\partial_\mu\phi^i\partial^\mu\phi^j) + \lambda_{4,ijkl}\phi^k\phi^l(\partial_\mu\phi^i\partial^\mu\phi^j) + \dots \\ &= \frac{1}{2}g_{ij}(\phi)(\partial_\mu\phi^i\partial^\mu\phi^j)\end{aligned}$$

Feynman diagrams



Geometric amplitude

$$\begin{aligned}\mathcal{A}_5 &= \nabla_k R_{iljm} s_{45} + \nabla_l R_{ikjm} s_{35} + \nabla_l R_{ijkm} s_{25} \\ &\quad + \nabla_m R_{ikjl} s_{34} + \nabla_m R_{ijkl} (s_{24} + s_{45})\end{aligned}$$

# Scattering Amplitudes

## Universal relations!

Geometric soft theorem

$$\lim_{q \rightarrow 0} \mathcal{A}_{n+1} = \nabla_i \mathcal{A}_n$$

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$$\mathcal{A}_4 = R_{ikjl} s_{34} + R_{ijkl} s_{24}$$



# Scattering Amplitudes

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$$\mathcal{A}_4 = R_{ikjl} s_{34} + R_{ijkl} s_{24}$$

# Scattering Amplitudes

## Hidden structure!

Nonlinear sigma model

$$A = \sum_{i \in \Gamma} \frac{c_i n_i}{d_i}$$

$$c_i + c_j = c_k$$



FIG. 1. Trivalent graph for four particles.

$$c_i = R_{abcd} \sim f_{abx} f^x_{cd}$$

# Scattering Amplitudes

## Hidden structure!

Nonlinear sigma model

$$\mathcal{A} = \sum_{i \in \Gamma} \frac{c_i n_i}{d_i}$$

Curvature-kinematics duality

$$c_i + c_j = c_k$$



$$n_i + n_j = n_k$$

FIG. 1. Trivalent graph for four particles.

$$c_i = R_{abcd} \sim f_{abx} f^x_{cd}$$

$$n_i = s_{12}(s_{12} + 2s_{23})$$

# Scattering Amplitudes

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FIG. 1. Trivalent graph for four particles.

$$c_i = R_{abcd} \sim f_{abx} f^x_{cd}$$

Double copy: Galileon theory = (NLSM)<sup>2</sup>

$$n_i = s_{12}(s_{12} + 2s_{23})$$

Galileon Theory

Curvature-kinematics duality

$$c_i + c_j = c_k$$

$$n_i + n_j = n_k$$

$$\mathcal{M} = \sum_{i \in \Gamma} \frac{n_i n_i}{d_i}$$

# Scattering Amplitudes

## Hidden structure!

Nonlinear sigma model

$$\mathcal{A} = \sum_{i \in \Gamma} \frac{c_i n_i}{d_i}$$

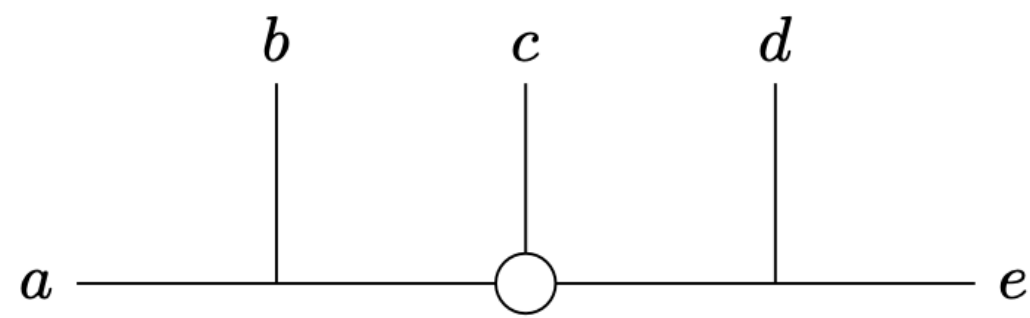


FIG. 2. Trivalent graph for five particles.

$$c_i = \nabla_c R_{abde} \neq f_{abx} f^x_{cy} f^y_{de}$$

Double copy: Galileon theory = (NLSM)<sup>2</sup>

Galileon Theory

Curvature-kinematics duality

$$c_i + c_j = c_k$$

$$n_i + n_j = n_k$$

$$\mathcal{M} = \sum_{i \in \Gamma} \frac{n_i n_i}{d_i}$$

# Scattering Amplitudes

## Geometry with fermions

$$\mathcal{L} = \frac{1}{2} g_{ij}(\phi) (\partial_\mu \phi^i \partial^\mu \phi^j) + i k_{\bar{p}r}(\phi) \bar{\psi}^{\bar{p}} \overleftrightarrow{\partial}_\mu \gamma^\mu \psi^r + i \omega_{\bar{p}ri}(\phi) (\partial_\mu \phi^i) \bar{\psi}^{\bar{p}} \gamma^\mu \psi^r$$

# Scattering Amplitudes

## Geometry with fermions

$$\mathcal{L} = \frac{1}{2}g_{ij}(\phi)(\partial_\mu\phi^i\partial^\mu\phi^j) + ik_{\bar{p}r}(\phi)\bar{\psi}^{\bar{p}}\overleftrightarrow{\partial}_\mu\gamma^\mu\psi^r + i\omega_{\bar{p}ri}(\phi)(\partial_\mu\phi^i)\bar{\psi}^{\bar{p}}\gamma^\mu\psi^r$$

Surprising structure!

$$\mathcal{A}_4 = (\bar{u}\not{p}_j u)\bar{R}_{\bar{p}rij}$$

$$\mathcal{A}_5 = (\bar{u}\not{p}_j u)\nabla_k\bar{R}_{\bar{p}rij} + (\bar{u}\not{p}_k u)\nabla_j\bar{R}_{\bar{p}rik}$$

# Scattering Amplitudes

## Geometry with fermions

$$\mathcal{L} = \frac{1}{2}g_{ij}(\phi)(\partial_\mu\phi^i\partial^\mu\phi^j) + ik_{\bar{p}r}(\phi)\bar{\psi}^{\bar{p}}\overset{\leftrightarrow}{\partial}_\mu\gamma^\mu\psi^r + i\omega_{\bar{p}ri}(\phi)(\partial_\mu\phi^i)\bar{\psi}^{\bar{p}}\gamma^\mu\psi^r$$

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Geometric soft theorem

$$\lim_{q\rightarrow 0}\mathcal{A}_{n+1} = \nabla_i\mathcal{A}_n$$



# Standard Model Effective Field Theory

Encode heavy new physics in effective operators

$$\mathcal{L}_{\text{SMEFT}} = \mathcal{L}_{\text{SM}} + \sum_i \frac{C_i}{\Lambda^2} \mathcal{O}_i + \dots$$

$X^3$			$H^6$		
$Q_G$	${}^6Q_{G^3}$	$f^{\mathcal{A}\mathcal{B}\mathcal{C}} G_{\mu}^{\mathcal{A}\nu} G_{\nu}^{\mathcal{B}\rho} G_{\rho}^{\mathcal{C}\mu}$	$Q_H$	${}^6Q_{H^6}$	$(H^\dagger H)^3$
$Q_W$	${}^6Q_{W^3}$	$\epsilon^{abc} W_{\mu}^{a\nu} W_{\nu}^{b\rho} W_{\rho}^{c\mu}$	$X^2 H^2$		
$H^4 D^2$			$Q_{HG}$	${}^6Q_{G^2 H^2}$	$(H^\dagger H) G_{\mu\nu}^{\mathcal{A}} G^{\mathcal{A}\mu\nu}$
$Q_{H\Box}$	${}^6Q_{H^4\Box}$	$(H^\dagger H)\Box(H^\dagger H)$	$Q_{HW}$	${}^6Q_{W^2 H^2}$	$(H^\dagger H) W_{\mu\nu}^a W^{a\mu\nu}$
$Q_{HD}$	${}^6Q_{H^4 D^2}$	$(D^\mu H^\dagger H) (H^\dagger D_\mu H)$	$Q_{HB}$	${}^6Q_{B^2 H^2}$	$(H^\dagger H) B_{\mu\nu} B^{\mu\nu}$
			$Q_{HWB}$	${}^6Q_{WBH^2}$	$(H^\dagger \tau^a H) W_{\mu\nu}^a B^{\mu\nu}$

**Table 3:** Bosonic even-parity dimension-six operators in the SMEFT. The first column is the notation of Ref. [26], and the second column is the notation used in this paper.

# Standard Model Effective Field Theory

## Renormalization Group Equations

$$\mathcal{L}_{\text{SMEFT}} = \mathcal{L}_{\text{SM}} + \sum_i \frac{C_i}{\Lambda^2} \mathcal{O}_i + \dots$$

1-loop UV divergence

$$\Delta S = \frac{1}{32\pi\epsilon} \int d^4x \left\{ \frac{1}{12} \text{Tr}[Y_{\mu\nu} Y^{\mu\nu}] + \frac{1}{2} \text{Tr}[X^2] \right\}$$

Curvature shows up

$$Y_{\mu\nu} = R^i{}_{jkl} (D_\mu Z)^k (D_\nu Z)^l + F_{\mu\nu}^a \nabla_j t_a^i$$

# Standard Model Effective Field Theory

## Renormalization Group Equations

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AH, Jenkins, Manohar '22

Assi, AH, Manohar, Pages, Shen '23

SMEFT RGE to dimension 8

$$\begin{aligned} \dot{C}_{H^6 D^2}^{(1)} = & -96 {}^6C_{H^6} {}^6C_{H^4 \square} - 12 {}^6C_{H^6} {}^6C_{H^4 D^2} + \left(352\lambda + 20g_1^2 + \frac{20}{3}g_2^2\right) \left({}^6C_{H^4 \square}\right)^2 \\ & + \left(-23\lambda + \frac{1}{8}g_1^2 + \frac{161}{24}g_2^2\right) \left({}^6C_{H^4 D^2}\right)^2 + (-64\lambda - 2g_1^2 + 12g_2^2) {}^6C_{H^4 \square} {}^6C_{H^4 D^2} \\ & - 22g_2^2 {}^6C_{H^4 \square} {}^6C_{W^2 H^2} + 6g_1^2 {}^6C_{H^4 \square} {}^6C_{B^2 H^2} - \frac{32}{3}g_1 g_2 {}^6C_{H^4 \square} {}^6C_{WBH^2} \\ & + 8g_2^2 {}^6C_{H^4 D^2} {}^6C_{W^2 H^2} + 6g_1^2 {}^6C_{H^4 D^2} {}^6C_{B^2 H^2} + \frac{43}{3}g_1 g_2 {}^6C_{H^4 D^2} {}^6C_{WBH^2} \\ & + 512\lambda \left({}^6C_{G^2 H^2}\right)^2 + (192\lambda + 4g_2^2) \left({}^6C_{W^2 H^2}\right)^2 + (64\lambda + 12g_1^2) \left({}^6C_{B^2 H^2}\right)^2 \\ & + (-3g_1^2 - 3g_2^2) \left({}^6C_{WBH^2}\right)^2 + \frac{80}{3}g_1 g_2 {}^6C_{W^2 H^2} {}^6C_{WBH^2} + \frac{8}{3}g_1 g_2 {}^6C_{B^2 H^2} {}^6C_{WBH^2} \\ & + \left(68\lambda + \frac{1}{2}g_1^2 - \frac{31}{6}g_2^2\right) {}^8C_{H^6 D^2}^{(1)} + \left(-8\lambda + 7g_1^2 + \frac{17}{3}g_2^2\right) {}^8C_{H^6 D^2}^{(2)}, \end{aligned} \quad (\text{C.23})$$

# Scattering Amplitudes

Surprising simplicity!

$$\mathcal{A}_4 = R_{ikjl} s_{34} + R_{ijkl} s_{24}$$

Universal relations!

$$\lim_{q \rightarrow 0} \mathcal{A}_{n+1} = \nabla_i \mathcal{A}_n$$

Hidden structure!

Double copy: Galileon theory = (NLSM)<sup>2</sup>

Practical calculations

RGE for Standard Model Effective Field Theory

**Thank you!**