

QFT and Invariants of Exotic Geometries

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N-PACT. 2020, Kristiansand

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DISCLAIMER: Officially, I am a mathematical physicist. My work is often not connected with reality in any meaningful sense.

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Can lead to some (valid) complaints:

- What is it good for, if not testable or even a description of reality?
- If not physics, then mathematics?
- How can methods of physics with no rigorous mathematical basis (e.g. path integral) be used to derive mathematically sound results?

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Tentative responses:

- The physics often leads one towards deep conjectures about geometry (e.g. Mirror Symmetry).
- Though the methods are mathematically iffy, the end results might still make sense (e.g. well-defined invariants).
- Personal hope: To connect the mathematics with phenomenological model building in particle physics and string phenomenology.

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String Theory is a higher dimensional UV-completion for the unification of the fundamental forces (including gravity). Playground for KK theories.

String Phenomenology: We postulate a space-time of the form

$$M_D = M_4 \times X_n, \quad ds^2 = g_{\mu\nu} dx^\mu dx^\nu + g_{ab} dx^a dx^b$$

where $D = 4 + n$. Here M_4 is the external space-time and X_n is usually taken as compact and small.

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Idea:

- Deformations of X_n geometry \leftrightarrow Fields in 4d geometry.
- Look for 4d theories containing the Standard Model spectrum.
- Need mechanism to lift additional massless fields out of theory (moduli stabilisation). **This is the difficult bit!**
- In particular, very hard to cook up models without runaway directions where X_n decompactifies.
- Geometric invariants can be useful in this regard.

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In string phenomenology, one often ignores (quantum) effects of fields propagating in the internal directions of X_n .

- Hard then to stabilize all moduli (runaway directions + extra massless fields).
- These quantum effects must eventually be included to have a consistent model.

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Threshold corrections:

- Such internal quantum effects are colloquially called threshold corrections.
- Often required to get the right RG-flow of effective field theory couplings, etc.
- Might help with moduli stabilisation: Generate non-perturbative couplings in effective physics.
- Drawback: Often hard to compute.

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In special cases (often with supersymmetry, etc) the computation of threshold corrections simplify: They derive from **topological QFTs**.

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The gist:

- The internal space X_n is equipped with some geometric structure Ψ (gauge-field, metric, complex structure, etc).
- The geometric structure Ψ is governed by some topological action $S(\Psi)$ deriving from the (stringy) higher-dimensional supergravity.
- Moduli space: Space of solutions of Ψ to equations of motion $\delta S(\Psi) = 0$.

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Quantum corrections: We quantize $S(\Psi)$

$$Z(X_n) = \int \mathcal{D}\Psi e^{iS(\Psi)}$$

Note: This is a **topological invariant** associated to the geometry X_n (integrate over geometric structures). Threshold corrections are associated to such topological invariants of internal geometries.

Example: Chern-Simons Theory

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Chern-Simons (CS) theory appears in string compactifications when higher dimensional objects called branes are wrapped on sub-manifolds M_3 within X_n .

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Chern-Simons (CS) theory appears in string compactifications when higher dimensional objects called branes are wrapped on sub-manifolds M_3 within X_n .

CS theory can be defined for a (curved) three-dimensional manifold M_3 , equipped with a gauge theory and associated gauge connection A_a :

$$S_{CS}(A) = \int_{M_3} \text{tr} \left(A_a \partial_b A_c + \frac{2}{3} A_a A_b A_c \right) \epsilon^{abc} .$$

Note: CS-theory is defined without a metric \Rightarrow topological theory.

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Equations of motion (suppressed gauge indices):

$$F_{ab} = \partial_{[a} A_{b]} + A_{[a} A_{b]} = 0 .$$

These correspond to flat gauge bundles.

The moduli space of flat gauge bundles on a compact space is finite-dimensional. Gives rise to a finite number of fields in the effective theory.

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We define the partition function for a gauge group G (e.g. $G = SU(N)$)

$$Z(M_3, G) = \int \mathcal{D}A e^{iS(A)} .$$

No metric dependence $\Rightarrow Z(M_3, G)$ should be an **invariant** [Witten '89, Reshetikhin-Turaev '91], depending only on the topological data of M_3 and the gauge group G .

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One can use $Z(M_3, G)$ to extract other interesting invariants of M_3 . E.g. the one-loop partition function [Witten '89], and:

- A perturbative computation of $Z(M_3, G)$ leads to the definition of **universal perturbative invariants**, independent of the gauge group [Axelrod-Singer '91].
- CS theory is related to the topological open string [Witten '92] \Rightarrow relations to Gromov-Witten invariants by open-closed string duality [Vafa '01, ..].
- Expectation values of Wilson loops: **Knot invariants** proportional to the **Jones Polynomial** of the knot [Witten '89, ..].

The study of quantum CS theory is a fruitful endeavour to this date.

There are many other generalisation and (quasi) topological theories which deserve mentioning:

- Holomorphic Chern-Simons theory (on six-dimensional Calabi-Yau manifolds): Give rise to Donaldson-Thomas invariants.
- Different types of volume (Hitchin) functionals in ordinary, generalised and exceptional geometry.
- Kodaira-Spencer gravity and Wittens topological strings.
- Generalisations combining gauge and gravitational degrees of freedom (my own work in heterotic string theory).

These all appear in different parts of string geometry, sometimes with equivalences (string dualities) between them.

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Explicit computation (resulting in actual numbers) for invariants can be very tricky.

Exact methods in QFT such as localisation can be useful. Usually requires extended supersymmetry implying a more rigid geometric structure.

Localisation: Reduces formal expression (path integrals) to integrals over finite dimensional moduli space of classical solutions which can in principle be performed (and sometimes is).

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Localisation: Reduces formal expression (path integrals) to integrals over finite dimensional moduli space of classical solutions which can in principle be performed (and sometimes is).

First step: One-loop computation. Compute the partition function of the quadratic approximation to the action. This action is usually of the form

$$S(\Psi) = \int_{X_n} \delta\Psi^* \Delta \delta\Psi + \dots,$$

where Δ is generically some elliptic linear operator.

Result: Ray–Singer torsion of Δ . This is often a topological invariant, but can have gravitational anomalies interesting in their own right.

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Let's recall the main points of the talk:

- There is a vibrant mathematical community studying geometric invariants.
- These invariants can be defined, and often computed using methods of QFT.
- They can be relevant for phenomenology, and string model building in particular.
- Can be hard to compute explicitly (numerically??).
- Often necessary to resort to special cases with extended supersymmetry and rigid geometric structures where exact methods of QFT can be employed.

Thank you!

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Thank you for your attention!